Io's **Heat** Flow from Infrared Radiometry: 1983-1993

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Abstract. We report the following results from a decade of infrared radiometry of Io: (1) The average global heat flow is more than -2.5 Wm², (2) large warm (<200 K) volcanic regions dominate the global heat flow, (3) small high-temperature (> 300 K) "hotspots" contribute little to the average heat flow, (4) thermal anomalies on the leading hemisphere contribute about half of the heat flow, (5) a substantial amount of heat is radiated during lo's night, (6) high-temperature ($\ge 600 \text{ K}$) "outbursts" occurred during -4% of the nights we observed, (7) "Loki" is the brightest, persistent, infrared emission feature, and (8) some excess emission is always present at the longitude of Loki, but its intensity and other characteristics change between apparitions. Observations of 10 at M (4.8 pm), 8.7 µm, N (10 pm), and Q (20 μm) with the Infrared Telescope Facility presented here were collected during nine apparitions between 1983 and 1993. These measurements provide full longitudinal coverage as well as an eclipse observation and the detection of two outbursts, Reflected sunlight, passive thermal emission, and radiation from thermal anomalies all contribute to the observed flux densities. We find that a new thermophysical model is required to match all the data. Two key elements of this model are (1) a "thermal reservoir" unit which lowers daytime temperatures, and (2) the "thermal pedestal effect" which shifts to shorter wavelengths the spectral emission due to the reradiation of solar energy absorbed by the thermal anomalies. The thermal anomalies are modeled with a total of 10 source components at five locations. Io's heat flow is the sum of the power from these components.

1. Introduction

Thermal anomalies on the Jovian satellite Io were first recognized unambiguously in data from the Voyager 1 infrared spectrometer (IRIS) experiment [Hanel et al., 1979]. Since that time, a number of ground-based techniques have been used to study Io's volcanic activity. A review of pre-Voyager infrared radiometry (including eclipse observations) showed that the signature of the volcanically heated areas was present in the older data. Furthermore, reanalysis of these data provided the first estimate of Io's large heat flow, namely, $2 \pm 1 \text{ W m}^{-2}$ [Matson et al., 1981]. Subsequent observations of lo's infrared spectrum and further analysis of the Voyager IRIS data have confirmed these conclusions [Morrison and Telesco, 1980; Sinton, 1981; Sinton et d., 1983; Tittemore and Sinton, 1987; McEwen et al., 1992]. Recent reviews are given by Pearl and Sinton [1982], Nash et al. [1986], and McEwen et al. [1989].

lo's heat flow results primarily from tidal interactions with Jupiter. Energy is extracted from **Io's** orbit and deposited as heat in the interior of **Io[Peale et al., 1979; Yoder, 1979]**. Thus, the level of 10's heat flow is deter mined by the tidal dissipation properties of 10, Jupiter, and the complicated orbital mechanics of the Jovian system. These interactions make **Io's** heat flow a key parameter which constrains models of 10, Jupiter, and the interactions between them (see review by Greenberg [1989]).

Several observers have reported outbursts, or large temporary increases in 10's thermal emission [Witteborn et al., 1979; Sinton, 1980; Sinton et al., 1983; Johnson et al., 1988]. These

relatively rare events, associated with temperatures greater than 600 K, are quite impressive at 3.4 or 4.8 μ m. However, when integrated over time, they do not contribute significantly to the average heat flow because of their relatively short durations. An additional discussion of the variability of Io is found in the. work by *Howell and Sinton* [1989].

Several Earth-breed telescopic techniques have been used to identify the locations of hotspots. These include near-infrared imaging, speckle interferometry, occultations, and polarimetry [Howell and McGinn, 1985; Goguen and Sinton, 1985, 1988; Sinton et al., 1988; Goguen et al., 1988; Spencer et al., 1990; McLeod et al., 1991]. The locations of (positive) thermal anomalies have generally been referred to in the literature as "hotshots," even though the range of temperatures associated with these areas includes values less than 250 K. We use the terms "thermal anomaly" and "hotspot" interchangeably in this paper.

The Voyager IRIS data and the ground-based eclipse observations provide information about hotspots for only part of lo's surface. In order to determine the distribution and variability of thermal anomalies, we initiated a program of multiwavelength radiometry of Io's thermal emission as a function of orbital phase angle, or west longitude of the sub-Earth point (see appendix note 1). These observations take advantage of the fact that hotspots of the type originally detected by IRIS will dominate the disk-integrated thermal emission from Io in the 7 to 14 µm region even though they cover only a small amount of the total surface area (see Figure 1). In contrast, the emission from 10 at 20 µm comes mainly from regions which are heated by the sun and from lower temperature thermal anomalies. At wavelengths shortward of -6 μ m, an increasing fraction of the observed flux density comes from reflected sunlight. High-temperature hotspots still can make an identifiable contribution down to wavelengths as short as ~2 µm during an eclipse [Sinton et al.,

Some observations in our 10 program have been discussed by Johnson et al. [1983, 1984, 1988], Matson et al. [1993], and Veeder et al. [1993], The supporting details of those measurements are all included here. Early results showed ' (1) the 8.7 µm emittance varies considerably with longitude, (2) the longitude of the maxima in the infrared light curves is very close to the largest anomaly observed by the Voyager IRIS (Loki Patera at 309"W and + 13°N), and (3) the 8.7 μ m emittance at all longitudes is larger than expected from the reradiation of absorbed solar energy alone. We now have a relatively uniform data set for four band-passes centered at 4.8. 8.7, 10, and 20 µm obtained over the period from 1983 to 1993. In this paper we will (1) report the data acquisition and reduction methods, (2) describe a model which includes the "thermal pedestal effect" for calculating the contributions to the observed flux densities from both absorbed insolation and geothermal heat, and (3) use the results of this model to characterize the total heat flow from Io and the variations in its volcanic activity and power levels from 1983 to 1993.

2. lo Observations

The radiometric observations of 10 reported here were collected between February 18, 1983, and March 29, 1993 (UT), at the 3-m Infrared Telescope Facility (IRTF) on Mauna Kea. Hawaii. A summary of the circumstances of our lo observations is given in Table 1 (see appendix note 2). In total, we had 175 hours of photometric condition sduring 55 nights. The IRTF bolometer and data acquisition system were used with the standard filters at M (4.8 μ m), N (10 μ m), and O (20 μm) as well as with the narrow-band-pass 8.7 μm filter $(\Delta \lambda \sim 1 \,\mu\text{m})$. The IRTF Cassegrain system utilizes a wobbling secondary mirror to difference two fields of view on the sky. An 8 arcsec diameter aperture, a beam separation of 10 arcsec. and a frequency of 13 Hz were selected as a practical compromise to obtain accurate radiometry with adequate signal to noise ratio (SNR), sky subtraction, and scattered light rejection. The 8 arcsec aperture is much larger than the convolution of typical infrared seeing with the disk of **lo**, which helps to ensure accurate radiometry, since it minimizes the effect of small guiding errors. For eclipse observations, Io was acquired with a precise offset from Jupiter or another satellite. Tracking rates were updated by means of the lRTF automatic centering capability. The two telescope beams were typically oriented north-south except for observations very near the disk of Jupiter, which required that they be rotated perpendicular to the direction of the center of Jupiter's disk. This arrangement suppressed the offset which otherwise might occur from imperfect cancellation of the sky gradient due to scattered light from Jupiter.

During each night, the data stream from the IRTF bolometer system was captured in real time as an ASCII file. This tile was converted to a condensed format by means of our own semi-automatic editor (written in FORTRAN). This standardized log was checked against our handwritten notes and corrected for typographic and other errors. Logs for a few nights had to be keyed from the IRTF printout and verified by hand.

A typical observational sequence consists of four pairs of 10-s integrations at each of N, M, Q, 8.-? pm, and N. This cycle was repeated quickly (within 20 rein) relative to the rotation of **Io** (8.5" per hour). Measurements of 10 were alternated with cycles on appropriate standard stars to derive the atmospheric extinction and system sensitivity for each night (see also section **3**).

The basic quantity determined from the observations is the magnitude of the whole disk of Io relative to Alpha Lyrae. Nearby standard stars for each night were selected from a subset of the IRTF list. These generally included at least one star before and after Jupiter in Right Ascension. Approximately one star per hour was observed throughout an observing session. Typically, three or four stars were observed two or three times each during a night at both higher and lower air mass than the range of Io. Almost all observations were made at an air mass of less than 2. A few observations of other Jovian satellites (in particular, Callisto) provided an additional

check on the stability of the IRTF system. These are especially useful for confirming the reality of infrared outbursts on 10.

The IRTF bolometer system sensitivity was monitored in real time and remained stable. throughout a night. Experience has shown that the residual, systematic change in responsivity is less than -5% between equipment setups for adjacent nights, All of our reported radiometry have relatively small statistical errors (i.e., SNR >> 10). Adjacent measurements of 10 reproduce approximately as well as or better than repeated measurements of, and the self-consistency between, the standard stars on each night (i.e., better than 5 %). The comparison of 10 data from one apparition to another is dependent on the accuracy of the whole standard star network, as the stars used slowly change with the motion of Jupiter around the sky (see also the appendix).

3. Data Reduction Procedure

The adopted standard star magnitudes are given in appendix Table A 1. This photometric system defines Alpha Lyrae as 0.0 magnitude (mag) for each wavelength. The absolute calibration of the photometric system is given in Table A2, and the monochromatic corrections are discussed in the appendix.

The system sensitivity in magnitudes was plotted for each night as a function of air mass, and a linear best tit was used to derive the extinction and average sensitivity at M, $8.7~\mu m$, N, and Q. The typical range of values for extinction among photometric nights on Mauna Kea is 0.15-0.30 mag per air mass at 4.8, 8.7, and $10~\mu m$ and 0.3-0.50 snag per air mass at $20~\mu m$. In addition, the system sensitivity in magnitudes was plotted for each night as a function of time at each wavelength. These plots were examined for obviously spurious events as well as for systematic trends correlated with time, seeing or humidity. Marginal weather conditions sometimes required the deletion of blocks of observations at the beginning and/or end of a given night. Useful standard star observations always bracketed those accepted for Io. In all, only a few percent of the raw Io data were rejected as nonphotometric.

The reduction of the observations to emittance is described below for each wavelength. The reduced data are given in Table 2 (4.8 pm), Table 3 (8.7 pm), Table 4 (10 pm), and Table 5 (20 μ m) and are plotted in Figure 2 (4.8 μ m), Figure 3 (8.7 μ m), Figure 4 (10 pm), and Figure 5 (20 pm).

The reduction of **Ioradiometric** data at 8.7 μ m follows the moat straightforward procedure. The observations are sorted by **universal** date and time. The orbital **phase angles** are given in degrees (i.e., the sub-Earth west **longitude** or central meridian of **Io**). The **observationally** determined quantities for 10, the 8.7 μ m mags and the corresponding tlux densities (F) at the Earth (above the atmosphere), are given both in **Janskys** and in units of 10^{17} W cm¹¹. The spectral emittance (E) at Io's surface in units of 10^{2} W m¹² μ m¹ is calculated from

$$E = \frac{F\Delta^2}{R_{lo}^2}$$

where A is the geocentric distance and R_{\star} . Is the radius of 10. The 8.7 μ m band-pass is relatively narrow, and therefore no monochromatic correction is necessary.

The reduction of 10 radiometric data at N (1 O μ m) and Q (20 μ m) is similar to the procedure for 8.7 μ m data, with the exception that the wide N and Q band-passes require monochromatic corrections (see the appendix).

The reduction of 10 radiometric data at M (4.8 μ m) includes the additional consideration that the M data have contributions from both reflected sunlight and thermal emission (see appendix note 3). The reflected component is about 70-80% of the total [Witteborn et al., 1979; Sin(on, 1980; Morrison and Telesco, 1980]. This component varies as a function of the geometric albedo distribution on the surface of lo, geocentric and heliocentric distance, and the solar phase angle. The calculated flux density (10⁻¹⁷ W cm⁻² μ m⁻¹) of the reflected and thermal components and the emittance (10 ² W m⁻² μ m⁻¹) of the thermal component at 10's surface are included in Table 2.

We use a simple model to calculate the reflected flux density component for each of the observations at 4.8 μ m at Io's particular geometry. An approximate value for the reflected sunlight can be determined directly by differencing the values in and out of an eclipse, We make an improved estimate of the reflected component at 4.8 μ m by means of our background model, which includes the *thermal pedestal effect* (see section 4.4). Data from the well-observed August 3, 1986, eclipse were used to calculate this effect. Our results imply a geometric albedo for 10 of -0,7 at 4.8 μ m in good agreement with $P_{\rm M}$ = 0.72 [Goguen and Sinton, 1985]. This value is similar to the albedo at 0.6 μ m and consistent with the spectrum of Jo over this interval [Clark and McCord, 1980; Cruikshank, 1980; McEwen et al., 1988].

In order to derive the reflected component as a function of lo's longitude, we assume that the relative amplitude and phase of the light curve of this component at 4.8 µm is the same as for visual wavelengths as compiled by Morrison and Morrison [1977]. This step is justified by the small rotational color variation observed for lo between the visual and near-infrared [Johnson, 1971]. Similarly, the behavior with solar phase angle of the reflected component at 4.8 µm is assumed to be the same as that observed at visual and shorter infrared wavelengths which is approximated by a phase coefficient of 0,022 mag per degree (see also Figure 1 in Sinton et al., [1983]). This parameter is not strongly dependent on color except in the ultraviolet [Morrison et al., 1974; Veverka, 1977]

The largest uncertainty in this procedure is probably introduced by the difficult eclipse observations which determine the reflected component. An error in this normalization would propagate equally "over all longitudes. Nonphysical negative values for the emittance of Jo at 4.8 μ m may result if too large a value is calculated for the reflected contribution. I. Uncertainties in the phase and orbital corrections will introduce scatter in the calculated emittances between runs and possibly a small false component to the longitude variation. The absence of negative values in our data set and the presence of some apparitions showing weak longitude variation and others

with strong variations in Figure 2 suggest that such errors are not significant in this context,

A major change in the 4.8 µm albedo (for instance, due to a hemispheric-scale volcanic resurfacing event) could show up in our reduced data as a spurious longitude variation which is not due to thermal emission. However, Io's visual light curve has been remarkably constant over decades [Morrison et al., 1974; Morrison and Morrison, 1977; Veverka et al., 1986] and current HST ultraviolet observations of Io are consistent with the overall albedo pattern seen by Voyager [Sartoretti et al., 1992].

The loradiometric data from Tables 2, 3, 4, and 5 are binned by apparition and plotted against decreasing west longitude in Figures 2, 3, 4, and 5. These data directly show spatial and temporal changes in the thermal emission from Io. The general shape of these infrared light curves does not vary greatly from year to year, but the amplitude does vary, It is obvious that the distribution of hotspots on **lo** is not uniform. The Loki source (see appendix note 4) dominates lo's thermal emission at short wavelengths in the trailing hemisphere (1 80°- 360°W). As a result of the changing viewing geometry, Loki produces a distinctive, approximately cosine, peak near 309"W in the plots. By contrast, the leading hemisphere (O" - 180°W) usually shows relatively little variation with longitude. Further discussion on the variability in the light curves is given in section 5.1, in the context of our model for thermal emission,

4. Thermal Emission Model

Thermal models of airless planetary bodies generally fall into one of two categories: (1) those dealing with thermal, emission from the body as an idealized whole, and (2) those dealing with the details of temperature distribution as a function of position on the surface. The latter distributions include variations in thermophysical properties and time of day (so-called diurnal models). Disk-averaged models similar to those used in asteroid radiometry [e.g., Matron et al., 1978; Morrison and Lebofsky, 1979; Lebofsky and Spencer, 1989; Veeder et al., 1989] have usually been used to analyze telescopic, disk-integrated lodata [e.g., Morrison et al., 1972; Motion and Cruikshank, 1973; Mornson and Telesco, 1980; Matson et al., 1981; Sinton, 1981; Johnson et al., 1984, 1988]. When dealing with Voyager data and combinations of spacecraft data and ground-based telescopic measurements, both approaches have been used [e. g., McEwenet al., 1988; Simonelli and Veverka, 1988]. To date, no model has been successful in explaining all aspects of Jo's thermal properties.

Our approach in this paper is to use versions of 'slow" and "fast" rotating whole-disk models (see appendix notes 5 and 6) modified to meet Voyager-derived hemispheric albedos (see also sections 4.1 and 4.2). In addition, volcanic thermal anomalies or hotspots are included. This has enabled a number of new insights into the type of model required to match 10 infrared data and has allowed us to interpret most of Jo's observed characteristics. **Our** model is not unique, however,

and much work remains to be done to achieve a full thermophysical, time dependent description for this satellite.

To understand the motivation for the characteristics of our model, it is important to note the problems with previous approaches (both ours and those of other workers). In the past, Iong-wavelength (e.g., 10 and 20 μ m) observations of 10 have been analyzed in terms of absorbed solar energy using a thermal model appropriate for airless, slowly rotating bodies with very low thermal inertia. Previous models assumed that the anomalies had such high temperatures that the contribution from absorbed insolation was negligible at short wavelengths (i.e., less than 4.8 μ m). Eclipse results were viewed as due to a combination of two components: (1) the rapidly cooling background regions whose heat was only due to absorbed insolation, and (2) the hotspots whose temperatures were controlled by volcanic processes and remained constant on the timescale of an eclipse.

Such models have been recognized as failing to satisfy completely either the details of eclipse observations [Sinton and Kanninski, 1988] or the absolute level of Io's 20 µm emittance [McEwen et al., 1988]. The large size of this discrepancy can be seen by direct comparison with the eclipse of August 3, 1986. The observations of Io at 4.8, 8.7, and 20 µm before and after immersion are plotted in Figure 6a. A first-order attempt to use a "standard" model (see appendix note 5) in instantaneous equilibrium with solar insolation for the background emission of Io fails (see also section 4.1). The expected emittance from such a sphere with a bolometric Bond albedo of 0.44 (Table 8 for longitude 340°W) is shown by the dashed line. It predicts too much emission at 20 µm and too little at 8.7 µm. In contrast, an analogous "standard" model fits similar data for Callisto quite well.

Since this "standard" model (see appendix note 5) only accounts for background emission, it is easy to understand why it yields too little emittance at 8.7 pm. The shortfall can be fixed easily by including some contribution from hotspots. However, this approach fails because the predicted model emission at 20 µmincreases even further above the observations. The key question is, Why is the 20 µm emittance of loso low?

Bolometric albedos for Io high enough to match the 20 µm data (e.g., 0.7) were deemed reasonable before the detailed analysis of Voyager photometry. Then, it became evident that Io's Bond albedo was significantly lower: an average of 0.5 from Simonelli and Veverka [1984]; hemispheric averages of 0.45-0.52 from McEwen [1988] and McEwen et al. [1988]; and 0.6-0.8 for the brightest regions from Simonelli and Veverka [1988]. These improved albedo estimates result in a factor of 2 more absorbed power compared to older values. One of the suggestions of *McEwen et al.* [1988] was that the 20pm emittance levels, combined with Voyager albedo maps, might require a wavelength-dependent emissivity, with the value near 20 µm being about 15% lower than elsewhere. However, continued analysis of the Voyager IRIS data has ruled out any significantly lower emissivity at 20 µm[McEwen et al., 1992; J.C. Pearl and A.S. McEwen, private communications, 1993].

Another possible wavelength-dependent effect results from polarization. Polarization is not normally considered when interpreting radiometry of airless, particulate surfaces. 10 is unusual, and thermal radiation from hotspots has been shown to be polarized [Goguen and Sinton, 1985]. Because both surface temperature and emissivity decrease with distance from the center of Io*s disk, the contribution of long-wavelength radiation will be reduced compared to that expected from a unit emissivity model. However, our calculation of this effect for a smooth surface with the index of refraction (1.5 at $4.8 \,\mu\text{m}$) determined by Goguen and Sinton [1985] shows that emissivity would change by less than 1 % between 4.8 and $20 \,\mu\text{m}$ for a slow-rotator temperature distribution. Furthermore, the smooth surface represents a worst case calculation in the sense that it would produce the largest polarization effect.

Sinton and Kaminski [1988] encountered a related problem in analyzing the heating and cooling of Io during eclipses from their data set. They consider a range of models which illustrate the need for a mixture of regions with very different thermal inertia values (cf. the equilibrium and reservoir units in sections 4.1 and 4.2). However, even their best model (i. e., number 5) did not match the Voyager bolometric Bond albedo for Io, and they assumed that the thermal anomalies remained at constant temperatures.

Possible evidence for failure of the "standard" model for Io was also noted by Simonelli and Veverka [1988]. They pointed out that the limited Voyager IRIS observations of the nightside of Jo suggested temperatures about 35 K higher than calculated from a "normal" thermophysical model. The Voyager nighttime observations unfortunately suffered from low spatial resolution and poor pointing knowledge, leading to an uncertainty in whether the elevated temperatures resulted from multiple unresolved hotspots or were a true measure of nighttime "passive" emission [Pearl and Sinton, 1982; McEwen et al., 1992; J.C. Pearl, persona} communication, 1993], Nevertheless, Simonelli and Veverka [1988] suggested that the discrepancy might be an indication that heat was being transported from the dayside to the nightside of **lo**, possibly by a solid-state greenhouse effect [Brown and Matson, 1987].

Our strategy is to construct a simple model consistent with lo's known properties and including hotspot thermal emission from both geothermal emission and the reradiation of absorbed sunlight. We then apply this model to a well-observed eclipse to constrain several key parameters. Our thermal emission model for 10 involves three types of surface units: (1) a low**albedo** unit which is in instantaneous equilibrium with sunlight to provide for the falloff in emission when 10 enters eclipse, (2) a high-albedo thermal reservoir unit to allow for more thermal emission in the antisolar direction and thus lower the observed 20 µm flux density, and (3) the thermal anomalies. An essential step is to account properly for the heating of the anomalies by sunlight (i.e., the thermal pedestal effect defined in section 4.4), These thermal units produce five spectral components: the emission due to radiation of absorbed sunlight for each of the three units, the emission due to heat flow, and reflected sunlight at 4.8 μ m (see appendix note 7). The equilibrium and reservoir thermal units are assumed to be interspaced uniformly over 10's disk (except for [he sites of thermal anomalies). The anomalies, which are at specific locations, occupy several percent of the surface.

4,1. Equilibrium Thermal Unit

The equilibrium thermal unit is a passive background unit which is heated only by sunlight. It has a low thermal inertia such that it is in instantaneous equilibrium. Its emissivity is taken as 0.9. Later, we will find that it has a bolometric Bond albedo of 0.29 and that its areal coverage is about 20% of the surface of 10. This is similar in extent to the darker *regions* seen in the Voyager pictures. For the purpose of comparing with disk-integrated radiometry, the equilibrium unit is treated as a nonrotating sphere (see appendix note 5) whose surface temperature is a function of the Sun's zenith angle z (i.e., the angular distance from the sub-solar point). There is a small (geometrically) projected defect of illumination (from which there is no emission) due to the solar phase angle. Thus, for the daytime hemisphere

$$T_z = T_{subsolar} \cos^{1/4} z$$

$$T_{\text{nubsolar}} = \left[\frac{S_{\odot}}{r^2} \frac{1 - A}{\epsilon o} \right]^{1/4}$$

where So is the solar constant (1.388 x 10^6 erg cm⁻² s⁻¹ AU"*), A is the bolometric Bond albedo, ϵ is the emissivity, σ is the Stefan-Boltzmann constant (5.669 x 10^{-5} erg cm⁻² s⁻¹ T"), and r is the heliocentric distance in AU.

4,2. Thermal Reservoir Unit

A heat storage unit is the key to solving many of the problems that limit the approach of the "standard" model as defined in section 4.1 (see also appendix note 5). This passive background unit absorbs heat during the daytime but reradiates it at night as well as during the day. The emission in the direction of the Earth is reduced compared to that of an analogous equilibrium unit. The emissivity of the reservoir unit is taken as 0.9. Later, we will find that its bolometric Bond albedo ranges from 0.58 to 0.70 and that its area! extent is -80% of the surface of 10. This unit may correspond to the. bright, higher albedos seen in the Voyager images. "Computationally, the reservoir unit is treated as a rapidly rotating sphere (see appendix note 6) with a large thermal inertia whose surface temperature is a function of latitude ϕ . In this **case** no correction for defect of illumination is required. The temperature distribution is given by

$$T_{\phi} = T_{equatorial} \cos^{1/4} \varphi$$

$$T_{igoplus quatOddl} \xrightarrow{S_{\bigodot}} \frac{1}{\pi \epsilon_0} \frac{1}{\epsilon_0}$$

The nature of the reservoir unit is not further constrained by our present analysis. There are several processes that could operate on lo to provide the needed heat storage. The thermal inertia of the surface, alone, is an obvious candidate for carrying heat to the nightside. The fast rotator (see appendix note 6) may appear to be an unrealistic assumption. However, models developed by *Spencer et al.* [1989] indicate that, for lo's rotation period and heliocentric distance and with a thermal inertia similar to other Galilean satellites, the calculated diurnal temperature variations are relatively small and thus closer to a fast rotator than an equilibrium model.

If significant sublimation occurs, then the latent heat of volatiles could help provide the characteristics needed for the mode.]. Another possibility is a solid-state greenhouse in which sunlight is absorbed within the upper few centimeters *or more* of the surface. Such deep deposition increases the timescale for conduction to the surface. and heat is carried onto the nightside of Io[Brown and Matson, 1987].

4.3. Thermal Anomaly Unit

This unit consists of all of the thermal anomalies or hotspots. Each anomaly is assigned a particular location (latitude and longitude). Later their temperatures and areas are adjusted as needed to tit the orbital phase data. We find that this unit typically occupies -2 % of the surface. The anomalies are modeled as gray-body radiators with an emissivity of 0.9 (for all of our band-passes). In particular, their emissivity is not constrained by the 4.8 μm geometric albedo of lo which we determined in section 3. We assume that the anomalies may have significantly different properties from the disk average values.

Hotspot albedos were assigned following the general correlation between measured temperatures and albedos by McEwen et al. [1985]. For hotspots identified by IRIS, higher temperature anomalies have lower albedos (see appendix note 8), For the anomalies in our model which correspond to the Voyager hotspots, we have adopted the McEwen et al. [1985], albedos as listed in Table 6. These are normal albedos in the Voyager orange filter (-0.6 μ m) but given the uncertainties in the scattering laws from these anomalous surfaces, we have used them as reasonable estimates of the biometric Bond albedos for the purpose of the calculation of the thermal pedestal effect (see section 4.4). For the non-Voyager locations, we have assigned small warm components an albedo of **0.15.** typical of observed spots with temperatures greater than 300 K, and the larger, cooler components an albedo of 0.5, typical of observed areas with temperatures of less than 200 K [see McEwen et al., 1985]. The only exception was the third component required for Loki, where we used 0.3, close to the value for the warm unit (Loki 2) observed by McEwen et al. [1985]. The albedos and locations are tabulated in Table 6. They are used later in sections 5.1 and 6.1, where we adjust their temperatures and areas to fit the infrared light curves.

4.4. Thermal Pedestal Effect

The fourth spectral component is one which is due entirely to sunlight absorbed by the thermal anomalies themselves. Previously, this background component had been assumed to be negligible, since the anomalies occupy a small percentage of the surface and because their temperature was viewed as being regulated only by 10's heat flow. In terms of total power, this absorbed insolation is a small percentage. What has been overlooked, however, is that the anomalies provide a way to spectrally redistribute emission due to absorbed sunlight, resulting, in a significant effect at the wavelengths of our observations. We find that a key element in fitting the data is recognizing that the *solar heating and cooling of even volcanically heated regions* must be taken into account, especially for the lower temperature anomalies.

The *thermal pedestal effect* is the spectral blue shift which occurs in the background spectrum when sunlight is absorbed on a thermal anomaly caused by heat flow. Recognition of this effect leads to the separate concepts of *active* and *passive* components of the background spectrum. Passive components (i.e., the equilibrium and *reservoir* units) can be calculated a priori, while the active components cannot be computed until after the geothermal temperatures of the anomalies are specified. In practice, the active background is calculated simultaneously with the fit of thermal anomalies to the data. See Figure 7 for illustrations of the pedestal effect (see appendix note 9).

The solar energy absorbed by the anomalies elevates their temperatures during the day. If ϵ o T_{geo}^4 is the geothermal heat flow at an anomaly and ϵ o T_{abs}^4 is the absorbed solar power which is governed by the albedo, then the observed daytime temperature is obtained from

$$T_{obs} = \left(T_{gso}^4 + T_{abs}^4\right)^{1/4}$$

In our model, T_{geo} is specified for each anomaly and then T_{obs} is calculated. As our nomenclature indicates, the warming due to sunlight is superimposed upon a geothermally supported temperature pedestal. The amount of warming and the spectral distribution of the reradiated solar power are both functions of the temperature of the anomaly, Another complexity is that this effect is a function of the zenith angle of the sun at each anomaly. Thus, latitude and time of day must be considered (in addition to heliocentric distance).

The pedestal effect has a temporal signature which should be evident during **Io's** eclipses, when the sun is eclipsed, the **solar** energy absorbed on **hotspots** is radiated away much more rapidly than for the cooler, passive background elements. In eclipse, the hotspot temperature drops quickly to $T_{\rm geo}$, the level supported by **geothermal** heat flow alone, Previous analyses of eclipse data have not taken the thermal pedestal effect into account, This omission yields an apparent thermal inertia for **Io** which is too low.

4.5. **Determination** of Background Parameters

An accounting of the number of model parameters quickly shows that they exceed the constraints from daytime whole-disk observations alone. Observations both in and out of Jovian eclipse are necessary to solve for all of the parameters. For this purpose, we have chosen from our data set a well-observed eclipse disappearance on August 3, 1986 (see Tables 2, 3, and 5). The observed daytime and eclipse emittances at 10 (W rr' $^2\mu$ m 1) for 4.8, 8.7, and 20 μ m are summarized at the bottom of Table 7.

Since **lo** rotates -10° during an eclipse disappearance, we assume that we are seeing the same (or equivalent) *passive* surface units both before and after entry into eclipse. For the thermal anomalies in general, and for the strong, discrete Loki source in particular, we allow for the change in geometric foreshortening over this time interval.

From our eclipse radiometry, we have two data points at each of three wavelengths (Figure 6a) for a total of six constraints on the model as well as information on the area of lo's disk, the bolometric Bond albedo estimated from the Voyager photometry, and the relationship between the bolometric Bond albedo and the total background power radiated from Io. In addition to the general properties of our model previously discussed, there are nine model parameters to be determined. The passive background includes five of them, namely, the bolometric Bond al bedos and areas for the equilibrium and reservoir units; and the Bond albedo of Io at 4.8 pm. A reasonable fit to the eclipse data requires the use of two hotspot components at Loki, each defined by a temperature and an area. Since, at this stage, we are solving for the passive background only, we are not now concerned about the distribution of the hotspots. Later, when the longitude data are taken into account, we shall see that a further adjustment in the thermal anomalies is required.

The particular values and components of the model spectral emittances (W m' $^2\mu$ m') which are consistent with both the daytime and the eclipse data are given in Table 7. The model emittance in eclipse is the simplest. The 20 μ m emission has two components. Slightly more than half is due to the thermal reservoir unit. The remainder is from the thermal anomalies. At 8.7 μ m almost all of the signal is from the anomalies. At 4.8 μ m, only a trace of the emittance is from the passive background,

The daytime spectral distribution (i.e., before eclipse) has more components. The chief surprise is the size of the thermal pedestal effect at 8.7 μ m, which accounts for almost 30% of the **observed** radiation. At 4.8 μ m the thermal pedestal effect is -13%! The recognition of this component is important because previously it has been generally assumed that there was **no** significant background radiation at such short wavelengths. The sum of all the background units is more than 45% of the **emittance** at 8.7 μ m. At 20 μ m, the background is -75% of the total. Here, the 25% due to emission from the thermal anomalies is significant because it corrects the earlier assumption that hotspot emission was negligible at this wavelength. The percentages due to the pedestal effect depend on

the model albedos assigned to the hotspots.

Data at particular wavelengths constrain some model parameters more than others. The equilibrium unit is constrained by the daytime 8.7 and 20 μ m emission. 'I'he reservoir unit is constrained by the eclipse 8.7 μ m data and even more by the eclipse 20 μ m data. The pedestal effect is determined from daytime 4.8 and 8.7 μ m emission. In general, cooler regions are more constrained by the longer wavelength data, and hotter regions more by the 4.8 μ m data,

'I'he areal combination of the smaller, dark equilibrium unit with the larger, bright reservoir unit results in lowering the physical temperature of much of 10's daytime surface to well below the values that would be expected based on rapid equilibrium with absorbed insolation. This effect is what allows the low observed values of 20 μm emittance to be modeled successfully. The subsolar temperature of the equilibrium unit is -167 K whereas the equatorial temperature of the reservoir unit is -109 K (see Table 7). The latter does vary slightly with longitude as a function of its albedo (see discussion of background parameters below and also Table 8). Both temperatures also vary slightly with the heliocentric distance of Jupiter.

Figure 6b shows the different background models discussed for 10. The "standard" model (see appendix note 5) has a bolometric Bond albedo of 0.44. The eclipse model contains emission solely from the reservoir unit. The daytime model includes emission from the equilibrium unit as well. These passive backgrounds were then combined with the thermal anomalies to fit the daytime and eclipse data.

4.6. **Thermal** Model as a Function of Longitude

The previous discussion yields an average bolometric Bond albedo of -0.53 for Io which only applies for a narrow range of orbital longitude near 340" W (see Table 7). To interpret our infrared data, it is necessary to extend the model of the background to other longitudes. We have done this by utilizing the form of the albedo versus longitude curve labeled "solar" from Figure 8 of *McEwen et al.* [1988]. Table 8 lists the longitudinal values of the bolometric Bond albedo averaged over 10"W longitude bins. These albedos have been scaled to match the albedo of our eclipse model at longitude 340"W. This scaling will be discussed again later in section 4.7.

The **albedo** of the equilibrium unit is taken to be 0.29 at all **longitudes.** This unit occupies -20% of the surface. We then derive the **albedo** for the reservoir unit as a function of longitude as given in Table 8. The reservoir unit occupies -80% of the surface of **lo**.

At least five locations with a total of 10 temperature components are needed to match the shape of the corresponding parts of the infrared light curves (at all three wavelengths simultaneously). The model hotspot designations, albedos, longitudes (°W), and latitudes are listed in Table 6 (see also Table 7 and section 6.1). Since the Voyager IRIS experiment identified several major hotspots on the trailing hemisphere, we began with the locations of the most prominent ones, in the spirit of Occam's Razor. On the leading hemisphere, our

model includes additional thermal anomalies at two locations, each with two components. These are needed to match the shape of the corresponding parts of the infrared light curves (again, at all three wavelengths simultaneously). We locate these new hotspots at 170° and 80" W longitude, but at 0° latitude. This choice of latitude minimizes their heat flow contribution, If any of the actual hotspots on this part of 10 are. at higher latitudes, they would have larger surface areas (due to foreshortening) and thus produce even more heat flow than our model predicts. We find for all the spots that although their characteristics change from apparition to apparition, the data are well fit, assuming that the locations remain unchanged.

For the Loki region, it is necessary to consider three active components to match the peak in emittance near 300-320"W at all three wavelengths simultaneously. We have approximated tire distribution of temperature and area in this region with (1) a small hotspot (~450 K and radius -22 km), (2) a warm area (-245 K and radius ~140 km), and (3) a large cool area (-180 K and radius -300 km). The Pele region contains two components which are hotter but much smaller (-654 K and radius -6 km; -350 K and radius -180 km). A more subtle feature which appears in our data is near the longitude of the Voyager source Colchis. A model component at this position is necessary to match the asymmetry and broadness of the (composite) trailing side emittance peak (see also section 6. 1).

4.7. Caveats

The model presented in this paper should be viewed as the first in an evolutionary sequence. It is the first thermal emission model which is consistent with the known properties of Io, resolving large discrepancies between previous models and the Earth-based observations. It is a simple model which provides an accurate way to calculate the background emission and determine the basic characteristics of the thermal anomalies which are responsible for Io's heat flow, but it is not yet a complete description of Io's thermophysical properties. We arrived at the adopted model parameters by manual iteration after making appropriate range checks. A discussion of limitations of our model follows.

- 1. Our model conserves energy but it does not attempt to model the transient response of the surface either diurnally or in eclipse. Additional information remains to be determined about the **thermal** inertia of the equilibrium and reservoir units, characteristics of a possible solid-state greenhouse, and/or latent heat as well as possible effects due to atmospheric transport.
- 2. The passive background units in our model are distributed uniformly over the surface of 10 in a manner analogous to a tine "checkerboard." Since the poles a.. well as the midlatitudes of Io are darker than the equator, an obvious improvement would be to include the albedo distribution seen by Voyager at an appropriate resolution.
- 3. There is an offset of 0.09 between the Voyager bolometric Bond albedo of 0.44 at 340"W [McEwen et al., 1988] and the value of 0.53 in our model. The Voyager

albedo weighs all parts of the projected disk equally. In contrast, our model is most sensitive to the infrared emittance from warm low latitudes. The difference in albedo between the equatorial regions and higher latitudes is sufficient to explain the direction and size of this offset. Considering uncertainties in the calculation of Io's phase integral, we regard the level of agreement as satisfactory for this stage of analysis.

4. The approximation of the thermal anomalies as flat disks has obvious restrictions with respect to the large, cool regions. An improvement would be to represent them as elements on a spherical surface and to associate them, where possible, with identifiable surface albedo features in the Voyager (and eventually Galileo) images.

5. Thermal Emission Variability

Our infrared data set provides the first relatively uniform, long-term overview of 10's thermal characteristics. The basic nature of Io's thermal emission has remained relatively stable throughout 1983-1993. Nevertheless, there are significant variations between apparitions. Dramatic outbursts at 4.8 μ m from relatively hot sources (greater than 600 K) can vary rapidly on timescales of hours to days and occasionally reach very high levels of thermal emission.

5.1. Orbital Infrared Light Curves

The thermal emission model discussed in the previous section was used to match the infrared data as a function of Longitude for each apparition. First, the passive background is adjusted to account for the relatively small changes in absorbed sunlight due to heliocentric distance variation. Then, each thermal anomaly is assigned a "geothermal temperature," radius, and location. The active background, due to the pedestal effect, is calculated from the physical temperature of the hotspot in sunlight. We attempt to make a minimum number of changes, from year to year, in anomaly temperatures and sizes, consistent with producing a good match to the emittance versus longitude data. It does not appear to be necessary to change the model locations of any of. the thermal anomalies to produce good fits, suggesting that at least the average longitude distribution of thermally active areas on Io has remained approximately constant throughout the last decade.

The radius and temperature for each of the 10 model thermal anomalies (at five locations) are listed in Table 9 for each apparition. The model tits to the data at 4.8, 8.7, and 20pm during each apparition are shown plotted as the curves in Figures 2, 3, and 5. The 10 μ m data are not modeled because of the wide bandwidth of this filter which causes the effective wavelength to change, depending on the temperatures of the thermal anomalies. The model radii and temperatures vary from year to year, with the largest changes generally in the higher temperature components required to match the 4.8 μ m emission. One way of visualizing both the variation in Io's emittance by year and the models required to match the

data is to overlay the model emittance curves at each wavelength for all apparitions as in Figure 8 (see also section 6.1).

As seen in previous studies, there is more variation at 4.8 μ m than at longer wavelengths (8.7 and 20 μ m). For most years the thermal emission at all wavelengths is very similar in intensity and shape from apparition to apparition. The thermal emission from the trailing hemisphere is roughly consistent with that expected from the hotspots detected by Voyager [Johnson et al., 1984].

However, there are several exceptions to this broad generalization. Most obvious are the outbursts which will be discussed below in section 5.2. Other exceptions result from general increases in emission levels. On the leading hemisphere (O"- 1800), the 1987 apparition had more than twice the ernittance at 4.8 µm than the other apparitions. Emittances at 8.7 and 20 µm were also elevated. Emittance at 4.8 µm from the trailing hemisphere (dominated by Loki) appears to have two modes. One mode peaks at less than 0,05 W m² μ m and is representative of most apparitions. The other mode peaks at greater than 0.07 W m² µm¹ and occurred during the 1983, 1984, and 1989-1990 apparitions. This bimodal activity is much less distinct at 8.7 μ m, but the 1989-1990 and the 1984 model fits do have higher emittances than other apparitions. At 20 μ m the variation in emittance at Loki forms a continuum, with 1984 and 1989-1990 data showing the highest emittances.

5.2. Outburst Activity

Outbursts on **Io** are characterized **bv** a large increase in the $4.8 \,\mu\text{m}$ emission over a short period of time (i.e., hours to days). The first report of such an event was the **pre-Voyager** $5 \,\mu\text{m}$ enhancement reported by *Witteborn et al.* [1979]. Other events were subsequently observed, including one during the interval between **Voyager** encounters in 1979, which was apparently **connected** with a change in the appearance of the feature Surt [Sinton, 1980].

We have observed two well-characterized events of this sort, one on August 7, 1986, the second on January 9, 1990 (see appendix note 10). Both events were observed for several hours during the course of a single night. Our current data set includes observations from 175 hours over 55 nights. We conclude that such **events** are occurring -4% of the time during our program, This appears to be consistent with the number of other reports, but we do not have enough information concerning the total time other observers could have detected outbursts to calculate a more precise frequency.

These two outbursts have similar emission characteristics. The August 7, 1986 (leading hemisphere), outburst is quite large. The emittance at 4.8 μ m increased by over a factor of 10 from -0.02 to greater than 0.3 W m⁻⁻² μ m⁻¹! At 8.7 μ m it almost doubled. On January 9, 1990, the. emission from Io increased by a factor of 2.5 at 4.8 μ m and also 20% at 8.7 μ m. This event occurred on the trailing hemisphere near the same longitude as Loki(309°W longitude), Since we have data at each of our wavelengths during the outbursts, it is possible to characterize the sizes and temperatures of the sources responsible.

Preliminary analysis of the January 9, 1990, outburst on the trailing hemisphere is illustrated in Figure 9. This suggests a small (-5.6 km radius) high-temperature (-1200 K) source at the beginning of the event, and also indicates that the model source required at the end of the event is significantly cooler (-700 K) and larger (-13 km radius). While the latter size is large by terrestrial standards, it is only -0.2% of the area of Loki Patera. These characteristics are consistent with a high eruption rate silicate lava flow and the suggestion that 10's heat flow is coming from an ensemble of volcanic features with differing emplacement histories and temperatures [Carr, 1986; McEwen et of., 1992]. These temperatures are consistent with liquidus temperatures of basaltic magmas (1550-1600 K [cf. Lofgren et al., 1981]).

The August 7, 1986, outburst was discussed by *Johnson et al.* [1988], who reported that a model temperature of 900 K and a radius of 15 km were required to account for the flux density increases at each of the wavelengths. They further concluded that such high temperatures were strong evidence for the source being a silicate lava rather than a sulfur melt (whose STP boiling point is 715 K).

The earlier analysis of the 1986 event did not attempt to match the characteristics of the event in detail. Rather, the maximum emittance excesses at 4.8 and 8.7 µm were used to produce an estimate of the required high-temperature source. Since the 4.8 µm flux density decreased rapidly over the course of about 2.5 hours, it was clear that more detailed modeling would either require a different hotspot location, a changing source emittance, or both. Using our model for the anomaly characteristics required to match the rest of the apparition's data, we find that an even higher temperature source (- 1550 K and ~8 km radius) at a longitude of -35" W best fits the observations at both wavelengths during the outburst (see Figure 9). This is consistent with the location of a persistent source imaged by Spencer et al. [1990]. The fit is not perfect, and it is probable that this source changed both temperature and size during the observing session. However, there are not enough 4.8 µm data during the early hours of the night, nor enough $8.7 \mu m$ data late in the night, to obtain simultaneous solutions at both wavelengths and constrain a time variable model of the type used for the 1990 event.

6. Heat Flow

The heat flow from thermal anomalies on 10 is calculated from the \mathbf{sum} of $\mathbf{nonbackground}$ power radiated by each of the $\mathbf{hotspots}$ needed to match the infrared light curves. The heat flow power P for each model thermal anomaly is

$$P = \pi \epsilon \sigma T_{geo}^4 R^2$$

where ϵ is the emissivity, σ is the Stefan-Boltzmann constant, $T_{\rm geo}$ is the geothermal temperature, and R is the radius of the

anomaly. The hotspot radii and temperatures are given in Table 9, and the resulting heat flow is listed in Table 10.

6.1. Active Regions

Figure 10 shows the variations in power for the different temperature components in the five model regions: Loki, Pele, Colchis, Leading Hemisphere (170°W), and Leading Hemisphere (80° W). We will discuss the behavior of these regions and then turn to a discussion of the global heat flow from 10.

Loki. Loki is one of the best characterized volcanic features on Io. It was easily recognized in the Voyager IRIS data because of its temperature and size and the favorable geometry during the Voyager flyby [Hanel et al., 1979; Pearl and Sinton, 1982]. Loki has also proved to be the most easily detected thermal feature from Earth by a number of techniques, including our initial fill-disk radiometry [e. g., Johnson et al., 1984]. Other known or suspected anomalies, such as Pele and Babbar Patera, may contribute to the emittance near the longitude of Loki (see appendix note 4). The signature of the Loki region in our data is usually a prominent peak near 3090W longitude which falls off on either side more or less as expected from the change in projected area with rotation. In detail, the shape of the Loki feature is asymmetric and requires smaller sources to the east of 309"W to match the infrared data (see discussion of Pele, Colchis, and the leading hemisphere below),

The Io thermal model for Loki includes three temperature components which are probably closely grouped and whose variation with apparition is shown in Figure 10a. Pele is modeled separately (see below and Figure 10h). Our model values of Loki's characteristics are not constrained strictly by the physical size of the albedo feature at Loki or the IRIS derived temperature and area components, but rather are a composite of temperatures and areas which could in fact be distributed among-a number of sites near the same longitude. Based on the lines of evidence cited earlier, Loki itself has been the major contributor to the infrared emittance in this region throughout the decade.

The data in Tables 2, 3, 4, and 5 show that the volcanic activity in the **Loki** region has been continuous during the course of our study. The intensity of this source at 4.8 μ m has varied considerably due to changes in the moderate- and high-temperature components. Thus, in the 1988-1989 apparition, sources with temperatures greater than 200 K contributed much less emission than in 1983, 1984, or 1989-1990. In particular, the high-temperature component was near its lowest level at the time of the Ulysses flyby in 1992 as well as during the 1993 apparition.

Pele. Pele is about 50" east of **Loki** and **was** the site of one of the largest eruptive plumes seen by Voyager [Snuth et al., 1979]. It had one of the highest temperature anomalies detected by IRIS [Hanel et al., 1979]. The 10 thermal model includes a small high-temperature region similar to the Voyager-derived values [Pearl and Sinton, 1982] and, in addition, a larger but lower temperature component to match our data at longer wavelengths (see Figure 10b). The time history for **Pele** is qualitatively similar to **Loki**, with the bigh-

temperature component model area changing by an order of magnitude (essentially disappearing during three apparitions) and the larger low-temperature region area varying by a factor of 4 at most. This region and Colchis, discussed next, do not produce a clear separate "signature" in our data, in contrast to Loki, but together provide infrared emission at longitudes east of Loki required to match the shape of the emittance versus longitude data.

Colchis. Colchis was also identified as a thermal anomaly in IRIS data [Hanel et al., 1979]. It is approximately 50° east of Pele. This is a transition region in the infrared data between the peak due to Loki and the relatively constant emittance levels seen on the leading hemisphere at most wavelengths. The 10 thermal model includes a source with a single temperature component greater than 350 K which affects primarily the 8,7 and 4.8 μmemittance (see Figure 10c). The area and temperature of the Colchis model are relatively constant for most apparitions. The major exceptions are the 1989-1990 apparition, when the required area increased about a factor of 10, and 1993, when the area is similar to other apparitions, but a model temperature of 700 K is required to match the data.

170°W and 80°W. The 170°W and 80°W. regions on the leading hemisphere are not tied directly to Voyager observations. Voyager did not acquire high-resolution imaging or IRIS data on this hemisphere. The relatively constant emittance level with longitude over this hemisphere requires that sources of volcanic activity here not be concentrated in a single relatively small region, such as Loki, but rather be more uniformly distributed over Io's surface. Qualitatively this is in agreement with lower resolution imaging data from Voyager which show no large albedo feature or caldera structures analogous to Loki. Occultations reveal no evidence for any strong concentrated source on the leading hemisphere [Goguen et al., 1988].

The **lo** thermal model includes two relatively large sources in the leading hemisphere with two temperature components each, separated by 90" (see Figures 10d and 10e). These are sufficient to provide good matches to the observed longitude variations in the data. They should not be interpreted literally as specific volcanic localities but rather as parameterized "stand-ins" for what may be an ensemble of surface features having a similar combined area and temperature distribution. Partial confirmation of the. validity of the general characteristics of the 170°W source comes from recent analyses of IRIS data by McEwen et al. [1992], who find that a whole-disk spectrum from Voyager centered on 1550 W is matched by a thermal anomaly with a temperature of 169 K and a radius of 769 km. The lower temperature source at 170°W in our model varies from 140 K to 180 K and from 420 km to 760 km in radius.

The characteristics of all these model sources are relatively constant during the decade, the only major change being observed in 1987, where emittance from the leading hemisphere at all wavelengths was higher than normal. This required increased areas for the lower temperature components and higher temperatures for the high-temperature components. It may be significant that this higher emission level occurs the

year following the high-temperature outburst event in 1986 [Johnson et al., 1988]. If that event represented lava flow activity for some significant duration, elevated temperatures from cooling flows could still be releasing significant heat a year later [Carr, 1986].

6.2. Global Heat Flow

The model fits to the data at each apparition can be used to estimate the heat flow coming from the thermal anomalies, and thus the variation of the geothermal power with time and location on **lo**. Total heat flow power for a given apparition is calculated simply by summing the contributions from each of the anomalies.

Table 11 compares the total power (10" W) and heat flow (Wm⁻²) for Io during each apparition with the model contributions from the larger, cooler regions and the smaller, hotter regions. Clearly, the bulk of the power is coming from regions with a surface temperature of less than 200 K. Figure 1 la shows the history of power output from the thermal anomalies in our model. As expected from the larger variability of the 4.8 µm emittance compared with other wavelengths, the small amount of power from the hotter components tends to vary considerably from apparition to apparition. Total power tends to be less variable, however, since most of the power comes from the larger, cooler, and less variable anomalies which affect the 20 µm emittance the most, This implies that the activity in such regions is long lived.

In a similar manner, Table 12 compares the total power and heat flow with the model sources located on the leading and trailing (i.e., Loki) hemispheres. Approximately *equal* amounts of power come from each hemisphere. Figure 1 lb shows these powers and the total as a function of apparition. Both hemispheric and total power values are relatively constant, with an average total power of 1.05 x 10¹⁴ W (2.5 W m⁻²). This value is considerably larger than most previous estimates, but consistent with the even higher value of 3.2 ± 1.0 W m⁻² [McEwen et al., 1992; A.S. McEwen, personal communication, 1993].

Note that conduction through the crust of 10 for the more than -95 % of the surface that has no detectable thermal anomalies is not considered in the present model [see also Stevenson and McNamara, 1988]. As noted in previous discussions of volcanic power from Io, the power radiated by hotspot anomalies is a lower limit to the total, globally averaged heat flow [e.g., Johnson et al., 1984]. Figure 12a illustrates the key characteristics of the model spots for each apparition; it plots the area of each model component versus the temperature. Lines of constant power output are also shown, This plot clearly shows the dominant contribution of the large area, low-temperature anomalies to the total heat flow, as well as the larger variation of modeled area and temperature for the hotter components of the model.

The characteristics of our model anomalies for the past decade can be compared with what was observed by the IRIS instrument during the Voyager flyby. *McEwenet al.* [1992] have recently refined the pointing knowledge for Voyager and

reanalyzed the IRIS data. This permits the recognition and modeling of lower temperature anomalies than were previously possible [e.g., *Pearl and Sinton*,1982]. In Figure 12b, the IRIS thermal anomalies identified by *McEwenet al*. [1992] are compared with our data binned in 100 K increments. We find very good agreement between these and our model thermal anomalies. The trend between area and temperature from our results and those of *McEwenet al*. [1992] is qualitatively what is expected for a surface with multiple, cooling lava flows of different ages, as pointed out for sulfur models by *Sinton* [1982] and for silicate models by *Carr* [1986] and *McEwenet al*. [1992].

The consistency of our results with the IRIS analyses is particularly important for evaluating the new class of anomalies with temperatures less than 200 K. In our thermal model, these anomalies are a consequence of the new background model required to match our disk-integrated radiometry including eclipse observations. These anomalies are identified directly in the IRIS analysis of small Voyager regions [McEwen et al., 1992; A.S. McEwen, personal communication, 1993]. Lower daytime and higher nighttime background temperatures in agreement with our conclusions are also implied by the interpretation of IRIS data by McEwen et al. [1992]. The IRIS observations are an independent data set with respect to instrumentation (i. e., Fourier transform spectroscopy versus filter radiometry), observational resolution (i.e., disk-resolved versus disk-integrated), and calibration techniques. agreement of these two independent approaches lends considerable support both to the conclusion that higher heat flows result from the lower temperature anomalies and to the necessity of improved thermophysical models for Io's passive background.

Figure 13 compares various estimates of observed and calculated heat flow with constraints from tidal heating and orbital evolution. Both our present estimate and that of McEwen et al. [1992] increase the previously noted disc repancy between the geothermal output from lo and theoretical limits on the total amount of power dissipated in Io due to tides. A number of possible explanations have been proposed which include time variations in 10's internal properties, variations in the dynamical state of the system and variations in Jupiter's dissipation over the last 4.5 Gyr. It is not the purpose of this paper to explore this problem in detail (see Greenberg [1989] for a recent review). We note only that models of the dynamics of the system including tidal dissipation must now cope with even higher values of lo heat flow observed over the past decade. Accurate measurement of the change in Io's orbit during the current epoch from spacecraft observations and/or improved astrometry or else a better understanding of Jupiter's interior may be necessary to resolve this issue (see also Peale et al. [1979], Yoder [1979], Yoder and Peale [198 1], Cassen et al. [1982], Ojakangas and Stevenson [1986], Segatz et al. [1988], Fischer and Spohn [1990], and Ross et al. [1990]).

7. Summary and Conclusions

A new view of 10 can be synthesized from our multiwavelength observations made over more than a decade (nine apparitions). This data set defines the characteristic states of 10's thermal emission as a function of orbital phase angle. It allows anomalous conditions to be recognized and compared with other times, The probability of observing a large variation, or outburst, on any given observing night is about 4%. We observed two such outbursts and both had high temperatures typical of silicate lavas.

10's flux density at 20 μ m is curiously low compared to 8.7 μ m. From this, we have inferred that 10 is radiating a large part of its thermal emission at nighttime. We suggest three possible mechanisms which may be involved: (1) high thermal inertia (e.g., rock), (2) solid-state greenhouse, and (3) atmospheric transport (e.g., latent heat).

The thermal pedestal effect is an important aspect of physics which must be included in modeling Jo's thermal anomalies. This blue shift of the background spectrum for the reradiation of absorbed sunlight can lead to emittance components as large as 30% at 8.7 μ m and 13% (of the thermal emission component) at 4.8 μ m. The thermal pedestal effect gives rise to the concepts of active and passive background thermal emission. The spectrum for an active background cannot be calculated until the geothermally supported temperature of the underlying anomaly is specified, The thermal pedestal effect also has a temporal signature when 10 goes into eclipse, The solar energy absorbed on hotspots is radiated away more rapidly than for analogous, passive surfaces.

We have developed a new thermal emission model for 10 which has three types of surface units: a dark albedo equilibrium thermophysical unit, a high albedo thermal reservoir, and a set of thermal anomalies. The equilibrium unit covers about 20% of the surface. The thermal reservoir unit greatly increases thermal emission in the antisolar hemisphere by carrying significant heat into Io's nightside. This unit occupies about 80% of the surface. The areal combination of the smaller, dark equilibrium unit with the larger, bright reservoir unit results in lowering the 'physical temperature of much of, 10's daytime surface to well below the values that would be expected based on rapid thermal equilibrium with absorbed sunlight. The thermal anomalies are located at five positions, and the data for each apparition are modeled by varying the areas and temperature of 10 components. The area] coverage of the thermal anomalies is ~2%.

The thermal emission model allows us to compute accurately the daytime emission and thereby isolate the reflected component at 4.8 pm. The geometric albedo of 10 at 4.8 µm is about 0.7. A model for the orbital phase/longitude dependence of the passive background and the 4.8 µm albedo was synthesized by comparison with Io's visual light curve. The contributions due to individual hotspots were identified and tracked. Thus, individual histories for Loki, Pele, Colchis, and the behavior of anomalies on the leading side at longitudes of 170' and 80°W can be studied. The temperature versus area plot for our hotspots is similar to that predicted for a surface with multiple,

cooling lava flows of different ages. The anomalies may include overturning of crusts on lava lakes as well as lava flows. At least for the high-temperature outbursts, the lava must involve silicates.

The total heat flow is calculated by summing the radiated power from all of the anomalies. We find that 10's heat flow is greater than 2.5 W m⁻². Most of the beat flow comes from the large, persistent, low-temperature (less than 200 K) anomalies. Short-lived outbursts, despite their high temperatures., contribute very little to heat tlow due to their small sizes and short durations. Since we are not sensitive to conductive heat flow through the crust of the passive areas (most of Io's surface), our heat flow values are formal lower bounds on the actual heat flow. Nevertheless, this value alone is so high that tidal dissipation in 10 remains the only viable mechanism. The heat flow is large enough to rule out existing steady state models for dissipation in Jo. All of this suggests there may be further complexities in the Jupiter-10 tidal interactions and in the volcanic activity of Io.

Appendix

Magnitude

"For each observing night, standard stars were selected from a subset of the IRTF list. These generally included at least one star before and after Jupiter in RA. The Bright Star catalog number (BS) and the magnitudes are listed in Table A 1. These magnitudes have been adopted from *Hanner et al.* [1984], *Tokunaga* [1984], *Hanner and Tokunaga* [1991], and *Hanner et al.* [1992]. This network is expected to be self-consistent to within 0.05 msg. The value for p Gem at 8.7 μmwas developed as a secondary standard from comparisons with other stars within our own data set.

Absolute Calibration

The IRTF photometric system defines a Lyrae as 0.0 mag at each wavelength. The absolute calibrations at M (4.8 μ m), 8.7 μ m, N (10 pm), and Q (20 μ m) as listed in Table A2 have been adopted from *Beckwith et al.* [1976] and *Hanner et al.* [1984]. These values are expected to be more accurate than *10%. Calibration issues at particular wavelengths are discussed by *Campins et al.* [1985], *Rieke et al.* [1985], and *Tokunaga et al.* [1986]. other photometric systems that have been used for Jo are discussed by *Morrison and Lebofsky* [1979], *Sinton and Tittemore* [1984], *Campins et al.* [1985], *Rieke et al.* [1985], and *Tokunaga et al.* [1986].

Monochromatic Correction

A monochromatic **correction factor** is required for measurements of 10 with the N (10 μ m) and Q (20 μ m) filters. Over the width of these band-passes, the spectral flux density distribution of the **Galilean** satellites is significantly different from that for the standard stars used. There are a number of contributions which affect the correction which must be made to the observed (instrumental) magnitude (or flux density) to

obtain the equivalent monochromatic value which would have been measured with a very narrow band-pass at the effective wavelength. The factor Y is defined by

$$\frac{B(\lambda_{\star}, T_{\circ})}{B(\lambda_{\star}, T_{\star})} = Y \frac{\int B(\lambda, T_{\circ}) r_{\lambda} d\lambda}{\int B(\lambda, T_{\star}) r_{\lambda} d\lambda}$$

where the integrals are taken over the band-pass of interest, $B(\lambda, T)$ is the Planck function, T, is the effective temperature of the object, T_{λ} is the effective temperature of the reference star, and λ_{λ} is the reference wavelength, The system response function, r, as a function of wavelength is defined by

$$r = aft$$

where a is the atmospheric transmission function, f is the filter transmission function, and t is the combined detector and telescope response function. For the purposes here, we have taken t as constant over the filter band-pass, which is a reasonable approximation for a bolometer used with a reflecting telescope at N and Q, The relevant standard star temperature was taken as 10,000 K (i.e., approximating a Lyrae). The results are not very sensitive to the exact value adopted for T_{\bullet} above effective temperatures of -3500 K.

Atmospheric absorption from water vapor within the N (10 μ m) band-pass is approximately constant with wavelength. At N the measured flux density of 10 includes a significant hot volcanic component with an inferred color temperature of approximately 600 K and a cool "background" surface component with an estimated color temperature of perhaps 120 K at noon on the equator (depending upon the local albedo and also the heliocentric distance) as well as a range of various warm components with interrnediate color temperatures. Over the N filter we use an effective temperature of -170 K for Jo due to the average volcanic flux contribution. This yields the value used in Table 4 for Y (10 μ m) of 0.95 (i.e., -0.05 mag).

The column density of water in the atmosphere affects the calculation for the Q (20 µm) band-pass since the long-wavelength portion of the band-paw is effectively controlled by H₀0 The theoretical atmospheric absorption as a function of wavelength and H₂O column density was calculated by V, **Kunde** (personal communication, 1993). However, we found in practice that assuming a constant factor produced less scatter over many nights than calculating the effective wavelength of Q based on our estimate of the H₂O column density for each night. Thus, a 3-mm column water abundance and a $T_{\rm eff}$ at Q of -170 K yields the value used in Table 5 for Y (20 μ m) of 1,08 (i.e., 0.08 mag). Again, the effective temperature is higher than the expected physical temperature of Io of ~120 K due to the contribution from volcanic emission. Since this value for Y (20 µm) may vary by up to 20% between a very dry and a marginal night at the IRTF, our reported 20 µm measurements in principle could have a systematic error of up to this amount, However, the consistency of our observations from night to night and from apparition to apparition suggests that relative errors between nights are typically less than 10% (i.e., less than O. 1 mag).

Notes

- 1. Numerically, this longitude, has the same value as the orbital phase angle. Plots with west longitude decreasing to the right can be compared directly with U.S. Geological Survey maps of 10.
- $\overline{2}$. Data at 8.7 μ m are not available for April 12, 1980, and March 19, 1982, and so these dates are not considered further.
- 3. One approach that has been taken to this problem is to define a "pseudo-geometric albedo," p', from the total 4.8 μ m flux density which is thus larger than the physical albedo [Sinton, 1980]. This quantity can show changes in the hotspot activity of Io over a few days; but comparisons at different longitudes, and especially at different epochs, are confused by the different angular and heliocentric distance dependencies of the reflected and emitted components,
- 4. The association with Loki Patera is established by comparison of the Voyager observations from IRIS [Hanel et al., 1979; McEwen et al., 1985] with occultation, speckle, imaging and polarization observations which have resolved Loki[Goguen et al., 1988; Howell and McGinn, 1985; Spencer et al., 1990; McLeod et al., 1991; Goguen and Sinton, 1985, 1988]. Strictly speaking, our data only require that there be a source at the same longitude as Loki.
- 5.' The treatment in section 4.1 is the same as for the "standard" asteroid thermal emission model but with the empirical beaming factor (which is not used in our **Io** model) set to unity [cf. *Matson et al., 1978; Morrison and Lebofsky*, 1979; **Lebofsky** and Spencer, 1989; **Veeder** et al., 1989].
- 6. The treatment in section 4.2 is the same as for one of the "nonstandard" asteroid thermal emission models with an empirical beaming factor of unity [cf. Matson et al., 1978; Morrison and Lebofsky, 1979; Lebofsky and Spencer, 1989; Veeder et al., 1989].
- 7. For an emissivity of 0.9, the reflected component at 8.7 μ m is -1 % of the observed emission. The reflected component is even less at 20 μ m.
- 8. *Goguen et al.* [1988] occultation results confirm that the lowest albedo regions at Loki had the highest temperatures, which strengthens this relationship.
- 9, **Calculations** of the intrinsic luminosity of Jovian planets involve the same physics as the thermal pedestal effect. For example, see Figures 1-4 of *Hubbard* [1980].
- 10. Data for the night March 11, 1983, plot well above the leading hemisphere **measurements** collected on July 19 and 21 at all wavelengths, including 20 μ m. Rechecking the photometry indicates that while this is not an outstanding night, it is good enough and has no obvious problems that could immediately call the data into question. Given the highly variable nature of Io, we are including the data for completeness. However, observational concerns such as north/south differences in the sky could have caused the discrepancy. Addition-

ally, the intensity of the emission level at $4.8 \mu \text{m}$ is consistent with other nonoutburst measurements during other apparitions. We therefore do not consider this measurement an outburst.

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Figure 1. Io's infrared spectral components. The "hotspot" and "background" curves are from our model for August 3, 1986 (UT), with Loki on the central meridian (309°W). The background curve contains only the passive background component. The reflected solar component (which dominates at short wavelengths) corresponds to an albedo of 0.44. Detailed spectral features, such as S0₂ frost absorption, are not shown. The thermal anomaly emission (including the active background) peaks near 13 pm. Passive background emission dominates at long wavelengths, but emission from thermal anomalies remains significant.

Figure 2. Io's 4.8 μm (M) emittance (W m¹² μm¹) versus west longitude (of the sub-Earth point) for apparitions from 1983 to 1993, Reflected sunlight has been subtracted. The longitude scale spans 1.5 revolutions. Consequently, the observations from Table 2 between O" and 180°W are plotted twice. Two large outbursta are evident: the first on August 7, 1986 (UT) (seen near 90°W), and the second on January 9, 1990 (UT) (seen near 330°W). The solid curve is our model fit. Temperatures and areas for the model hotspots are listed in Table 9. Model curves calculated for the outbursts are shown in Figure 9.

Figure 3. 10's 8.7 μ m emittance (W m⁻² μ m⁻¹) versus west longitude. The outbursts during August 7, 1986 (UT) (seen near 90" W), and during January 9, 1990 (UT) (seen near 330 °W), are not as strong as at 4.8 μ m (Figure 2). The upper solid curve is our model fit. The lower solid curve (just above the bottom of the frame) is the passive background caused by the heating the surface by insolation. (See also caption for Figure 2).

Figure 4. lo's 10 μ m (N) emittance (W m⁻² μ m⁻¹) versus west longitude. The outbursts during August 7, 1986 (UT) (seen near 90" W), and during January 9, 1990 (UT) (seen near 3300 W), are only barely noticeable. (See also caption for Figure 2).

Figure S. 10's 20 μ m (Q) emittance (W m⁻² μ m⁻¹) versus west longitude. The upper solid curve is our model fit. The lower solid curve is the model passive background which is caused by the heating of the surface by insolation. 10's albedo distribution controls the variation with longitude of the passive background which is the largest component of the 20 μ m emission. (See also caption for Figure 2).

Figure 6. Jo's emittance spectra for August 3, 1986 (UT). (a) The data are shown as solid circles (daytime) and open triangles (in eclipse). The solid and dotted curves are our model fits. The components of the model at each wavelength are given in Table 7. The dashed curve is a "standard" thermophysical model (see appendix note 5) based on the bolometric Bond albedo of 0.44 from Voyager data at 340°W longitude [*McEwen et al.*, 1988] (see also section 4. 1). This curve illustrates the problem of attempting to match the 20 µm data with conventional models. (b) Comparison of passive background model spectra with the standard model from Figure 6a.

Figure 7. Emittance spectra illustrating the thermal pedestal effect. (a) An example of a hotspot with a nighttime temperature of 32S K and a (higher) daytime temperature of 331 K due to heating by sunlight. The spectrum of the active background is the difference between these two curves. (b) The active background components for a hotspot with a bolometric Bond albedo of 0.29 (and emissivity of 0.9) as a function of its nighttime geothermal and daytime temperatures. Note that the power is the same for each curve, even though the spectral distribution is different, The lowest spectrum is the degenerate case without any thermal anomaly (in this case, the geothermal temperature is identically O K and the temperature at equilibrium with insolation is 140 K). Thus, the spectrum of the absorbed solar power is blue shifted from the longer wavelengths expected for a passive reradiation spectrum.

Figure 8. Overlay of all the 10 models for the 1983-1993 apparitions. The curves include both the anomaly and the background emittance for 4.8, 8.7, and 20 μ m. These plots show both the *general* stability of the shape of the light curves, the locations of the thermal anomalies, and their variation in intensity. The variation at 4.8 μ m is significantly larger than at longer wavelengths.

Figure 9. Emittance versus longitude of the sub-Earth point for the outbursts of August 7, 1986, and January 9, 1990 (UT), compared with model calculations at 4.8 and 8.7 µm. (Note that for a given event time. increases towards the left). Available "nonoutburst" data for these apparitions are also included. The thermal emission at $4.8 \, \mu m$ is plotted after removal of the reflected component. The curves labeled "model" correspond to our fits for the entire apparition. The curves for the passive background at 8.7 µm are labeled as "background." The differences between the model and background curves are due to the emission from the hotspots including the pedestal effect. The curves labeled with values for T and r show the result of adding the thermal emittance from a hotspot of the indicated temperature and radius to the model for the apparition. For the 1986 outburst, the model hotspot was placed on the equator. For 1990, the event was assumed to be located at Loki.

Figure 10. Heat flow histories of the thermal anomalies for the 1983-1993 apparitions. In addition to the total, power levels are also binned by temperature as indicated (see also Table 10).

Figure 11. **Io** heat flow during the 1983-1993 apparitions. (a) In addition to the total, power levels are also binned by temperature as indicated (see also Table 11). (b) Leading and trailing (i.e., **Loki**) hemispheres compared to the total (see also Table 11).

Figure 12. Surface area versus temperature log-log plot for **Io's** hotspots. The diagonal lines are contours of constant emitted power (10° to 10¹⁴ **W**). (a) All of the thermal anomalies for the 1983-1993 apparitions. (b) **Hotspots** for each apparition binned by 100 K increments compared with results from **McEwenet** al. [1992].

Figure 13. Comparison of Io's heat flow determinations with calculations of tidal dissipation and other power values from the literature. The heat flow for two values of Q_{lc} (50 and 100) and the upper limit from the orbital evolution are from Cassen et al. [1982] following the method of Goldreich and Sorer [1966]. The upper limit from the rate of orbital acceleration for the last 300 years is from Greenberg's [1987] analysis of data from Lieske [1987]. These estimates all are significantly lower than the observed heat flow values determined from thermal emission. Heating by the decay of radiogenic elements and heating by electrical induction are 2 orders of magnitude smaller and plot off scale.

- Figure 1. 10's infrared spectral components. The "hotspot" and "background" curves are from our model for August 3, 1986 (UT), with Loki cm the central meridian (309" W). The background curve contains only the passive background component. The reflected solar component (which dominates at short wavelengths) corresponds to an albedo of 0.44. Detailed spectral features, such as S0₂ frost absorption, are not shown. The thermal anomaly emission (including the active background) peaks near 13 μm. Passive background emission dominates at long wavelengths, but emission from thermal anomalies remains significant.
- **Figure** 2. 10's 4.8 μ m (M) emittance (W m⁻² μ m⁻¹) versus west longitude (of the sub-Earth point) for apparitions from 1983 to 1993. Reflected sunlight has been subtracted. The longitude scale spans 1.5 revolutions. Consequently, the observations from Table 2 between 0° and 180°W are plotted twice. Two large outbursts are evident: the. first on August 7, 1986 (UT) (seen near 90°W), and the second on January 9, 1990 (UT) (seen near 330°W). The solid curve is our model fit. Temperatures and areas for the model hotspots are listed in Table 9. Model curves calculated for the outbursts are shown in Figure 9.
- Figure 3. Io's 8.7 μ m emittance (W m''² μ m') versus west longitude. The outbursts during August 7, 1986 (UT) (seen near 90°W), and during January 9, 1990 (UT) (seen near 330°W), are not as strong as at 4.8 μ m (Figure 2). The upper solid curve is our model fit. The lower solid curve (just above the bottom of the frame) is the passive background caused by the heating the surface by insolation. (See also caption for Figure 2).
- Figure 4. Io's 10pm (N) emittance (W m⁻² μ m⁻¹) versus west longitude. The outbursts during August 7, 1986 (UT) (seen near 90°W), and during January 9, 1990 (UT) (seen near 330°W), are only barely noticeable. (See also caption for Figure 2).
- **Figure 5. lo's 20pm** (Q) emittance (W m⁻² pm-¹⁾ versus west longitude. The upper solid curve is our model tit. The lower solid curve is the model passive background which is caused by the heating of the surface by insolation. Io's **albedo** distribution controls the variation with longitude of the passive background which is the largest component of the 20 μm emission. (See also caption for Figure 2).
- **Figure 6. Io's emittance** spectra for August 3, 1986 (UT). (a) The data are shown as solid circles (daytime) and open triangles (in eclipse). The solid and dotted curves are our model fits. The components of the model at each wavelength are given in Table 7. The dashed curve is a "standard" thermophysical model (see appendix note 5) breed on tJre bolometric Bond albedo of **0.44** from Voyager data at **340**° W longitude [*McEwen et al., 1988*] (see also section **4.1**). This curve illustrates the problem of attempting to match the 20 μm data with conventional models. (b) Comparison of passive background model spectra with the standard model from Figure 6a.'
- **Figure 7.** Emittance spectra illustrating the thermal pedestal effect. (a) An example of a hotspot with a nighttime temperature of 325 K and a (higher) daytime temperature of 331 K due to heating by sunlight. The spectrum of the active background is the difference between these two curves. (b) The active background components for a hotspot with a **bolometric** Bond **albedo** of 0.29 (and **emissivity** of 0.9) as a function of its nighttime geothermal and daytime temperatures. Note that the power is the same for each curve, even though the spectral distribution is different. The lowest spectrum is the degenerate case without any thermal anomaly (in this case, the geothermal temperature is identically O K and the temperature at equilibrium with in solation is 140 K). Thus, the spectrum of the absorbed solar power is blue shifted from the longer wavelengths expected for a passive reradiation spectrum.
- Figure 8. Overlay of all the 10 models for the 1983-1993 apparitions. The curves include both the anomaly and the background emittance for 4.8, 8.7, and 20 pm. These plots show both the general stability of the shape of the light curves, the locations of the thermal anomalies, and their variation in intensity. The variation at 4.8 μ m is significantly larger than at longer wavelengths.

Figure 9. Emittance versus longitude of the sub-Earth point for the outbursts of August 7, 1986, and January 9, 1990 (UT), compared with model calculations at 4.8 and 8.7 μ m. (Note that for a given event time increases towards the left). Available "nonoutburst" data for these apparitions are also included. The thermal emission at 4.8 μ m is plotted after removal of the reflected component. The curves labeled "model" correspond to our fits for the entire apparition. The curves for the passive background at 8.7 μ m are labeled as "background." The differences between the model and background curves are due to the emission from the hotspots including the pedestal effect. The curves labeled with values for T and r show the result of adding the thermal emittance from a hotspot of the indicated temperature and radius to the model for the apparition. For the 1986 outburst, the model hotspot was placed on the equator. For 1990, the event was assumed to be located at Loki.

Figure 10. Heat flow histories of the thermal anomalies for the 1983-1993 apparitions. In addition to the total, power levels are also binned by temperature as indicated (see also Table 10).

Figure 11. Io heat flow during the 1983-1993 apparitions. (a) In addition to the total, power levels are also binned by temperature as indicated (see also Table 11). (b) Leading and trailing (i.e., Loki) hemispheres compared to the total (see also Table 11).

Figure 12. Surface area versus temperature log-log plot for Io's hotspots. The diagonal lines are contours of constant emitted power (10° to 10¹⁴ W). (a) All of the thermal anomalies for the 1983-1993 apparitions. (b) Hotspots for each apparition binned by 100 K increments compared with results from *McEwen et al.* [1992].

Figure 13. Comparison of **lo's** heat flow determinations with calculations of tidal dissipation and other power values from the literature. The heat flow for two values of Q_{lo} (50 and 100) and the upper limit from the orbital evolution are from *Cassen et al.* [1982] following the method of *Goldreich and Soter* [1966]. The upper limit from the rate of orbital acceleration for the last 300 years is from *Greenberg's* [1987] analysis of data from *Lieske* [1987]. These estimates all are significantly lower than the observed heat flow values determined from thermal emission. Heating by the decay of radiogenic elements and heating by electrical induction are 2 orders of magnitude smaller and plot off scale.

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Table 1, Page i

The labulated columns are the UT date, solar phase angle in degrees, heliocentric distance in AU, geocentric distance in AU the range of orbital phase covered in °W longitude, and the observers and telescope operators. The observers and operators are designated by their initials and are presented below.

Telescope Observers		2. Telescope Operators:
AL = Arthur Lane	HH = Heidi Hammel	CK = Charlie Kamınski
AT=Alan Tokunaga	LL=Larry Lobofsky	DG=Dave Gnep
DB=Diana Blancy	MH=Martha Haancr	JH = John Hamikon
DM = Dave Morrison	RB=Robert Brown	M=Miller
DLM = Denau Matson	RD - Richard Disley	WG=William Golish
PT = Edward Tedesco	RN = Robert Nelson	
GV = Glenn Veeder	Tj = Torreace Johnson	
JG≃Jay Goguen	WS = William Sinton	•

Table 2. M (4.8 μ m) Photometry of Io (J1).

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557 234 352 600 782 16 2.65 519 235 342 600 782 516 2.65 610 237 348 623 8.11 516 2.95 757 252 347 6.29 8.19 510 3.09 813 254 346 6.34 8.26 510 3.16 841 259 341 6.64 8.65 505 3.60 921 265 340 6.70 8.77 5.05 3.48 930 266 343 652 8.49 5.05 3.44 95 269 3.8 6.83 8.89 5.05 3.44 95 269 3.4 5.89 7.67 5.40 2.28 521 73 3.4 5.89 7.67 5.40 2.28 531 74 3.4 5.89 7.67 5.40 2.28	13/07/19	948	63				90	1.88	2.79
555 235 352 600 782 316 265 610 237 348 623 811 516 295 757 252 347 629 819 510 309 811 254 346 634 825 510 316 841 259 341 664 865 510 356 855 261 341 664 865 505 368 921 265 340 670 877 505 344 950 269 343 682 849 505 344 950 269 343 682 849 505 344 950 269 34 589 767 540 228 521 73 34 589 767 540 228 521 73 34 589 767 540 228 521 78	83/07/20	557	134	3 S2	8	7 82		2.65	3.97
610 237 348 623 811 716 295 757 252 347 6.29 8.19 510 309 813 254 346 634 8.26 510 3.16 844 259 341 664 8.65 515 3.63 952 265 340 6.70 8.77 5.05 3.68 930 266 343 6.82 8.89 5.05 3.84 1009 271 3.2 722 940 4.99 4.40 522 73 3.4 5.89 7.67 5.40 2.28 551 78 3.4 5.89 7.67 5.40 2.28 551 78 3.4 5.89 7.67 5.40 2.28 551 78 3.4 5.89 7.67 5.40 2.28 551 78 3.4 5.89 7.67 5.40 2.28 <td< td=""><td>83/07/20</td><td>559</td><td>235</td><td></td><td>6 00</td><td>7.82</td><td></td><td>2.65</td><td>3.97</td></td<>	83/07/20	559	235		6 00	7.82		2.65	3.97
757 252 347 6.29 8.19 \$ 10 3 09 811 254 346 6.34 8.26 \$ 10 3.16 844 259 341 6.64 8.65 \$ 10 3.15 854 226 341 6.64 8.65 \$ 10 3.55 852 265 340 670 8.77 \$ 505 3.84 930 266 343 682 8.49 \$ 505 3.44 950 260 343 682 8.49 \$ 505 3.44 950 260 343 682 8.49 \$ 505 3.44 950 271 3-2 722 940 4.99 4.40 522 71 3-4 5.89 7.67 5.40 2.28 531 78 3-4 5.89 7.67 5.40 2.28 531 78 3.4 5.89 7.67 5.40 2.28	3/07/20	610	237			S 1.1		2.95	14 4
813 254 346 634 826 510 3.16 844 259 341 664 8.65 510 3.55 854 259 341 664 8.65 505 3.60 921 265 340 672 8.49 505 3.44 950 266 343 682 8.49 505 3.44 950 260 343 682 8.49 505 3.44 950 260 343 682 8.49 505 3.44 950 260 343 683 8.89 505 3.84 1009 271 3.2 7.22 9.40 4.99 4.40 522 71 3.4 <t>5.89 7.67 5.40 2.28 531 74 3.4 5.89 7.67 5.40 2.28 534 5.9 5.63 7.33 5.47 2.50 704 82</t>	83/07/20	757	252			8.19		3.09	
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704 88 350 611 796 5.47 2.50 774 92 354 589 767 5.49 2.19 775 95 351 606 789 5.49 2.40 973 109 355 584 7.60 5.46 2.15 941 111 354 5.9 767 5.39 2.29 814 317 340 6.7 804 4.19 3.84 830 319 343 6.2 849 4.19 3.84 831 321 350 611 7.96 4.17 3.50 851 324 354 589 767 4.17 3.50	3/07/21	615	8				5.47	.	2.80
734 92 354 589 767 549 2.19 755 95 351 606 789 549 2.40 933 109 355 584 760 5.46 2.15 941 111 354 5.9 767 5.39 2.29 814 317 349 6.7 804 4.19 384 830 319 343 6.2 849 4.19 3.84 831 321 350 6.11 796 4.17 3.50 882 321 354 589 767 4.17 3.50	3/07/21	704	86			7%	5.47	2.50	3.76
755 95 351 606 789 549 240 933 109 355 584 7.60 5.46 2.15 941 111 354 5.9 767 539 229 814 317 349 6.7 804 4.19 384 830 319 343 6.2 8.49 4.19 3.89 842 321 350 611 796 4.17 3.50 851 324 589 7.67 4.17 3.50	(3/07/2)	734	%			7.67	5.49	2.19	3.29
933 109 355 584 7.60 5.46 2.15 941 111 354 5.9 767 5.39 2.29 814 317 349 6.7 8.04 4.19 3.84 830 319 343 6.2 849 4.19 4.30 842 321 350 611 7.96 4.17 3.79 884 524 354 589 767 4.17 3.50	13:07/21	755	ž			7 89	5.49	2,40	J.62
941 111 354 5.9 767 539 2.29 814 317 349 6.7 8.04 4.19 3.84 830 319 343 6.2 849 4.19 4.30 842 321 350 611 7.96 4.17 3.79 854 324 354 589 767 4.17 3.50	12/20-53	933	8			7.60	5,46	2.15	3.23
814 317 349 6.7 8.04 4.19 3.84 810 319 343 6.2 849 4.19 4.30 842 321 350 611 7.96 4.17 3.79 854 374 354 589 7.67 4.17 3.50	13/07/21	91	Ξ				5.39	2.29	ر 4
830 319 343 62 849 4.19 4.30 842 321 350 611 7.96 4.17 3.79 854 324 354 589 7.67 4.17 3.50	83/08/07	8 2 4	317				4.19	3.84	6.37
842 321 350 641 7.96 4.17 3.79 854 324 354 589 767 417 3.50	83/08/07	830	319	3 43		8 49	4.19	4.30	7.13
854 324 354 589 767 417 3.50	83/08/07	842	321			7.96	4.17	3.79	6.29
	20.27	85:	324	3 54	5 89	7 67	4.17	3.50	5.81

Table 2, Page 1

Table 2. M (4.8 μ m) Photometry of Io (J1) (continued).

85/07/10	85/07/10	60//0/20	01,000,10	20000	24/26/10	92/00/10	01/00/10	26/06/06	85/06/10	85/06/10	85/06/10	93/00/07	95/00/09	85/06/00	85/06/00	85/06/09		85/06/09	85/06/08	85/06/08	85/06/08	85/06/08	85/06/07	27/00/07	85/06/06	85/06/06	85/06/06		85/06/06	85/06/06	84/08/20	84/08/20	84/08/19	84/08/19	84/08/19	84/08/19	84/08/19	84/07/19	04/0//19	84/0//19	84/07/19	84/06/27	83/08/07			6	Date		
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10.2	9.84	¥.		3	•	8.97	8	2	-	# 57	8.49	8	2 3		9			73		3	8 26	1.26			2	-	40		6 21	6 56	8 97	2 86	64 60 27	8.65	8.73	1.89	2 11	12.9	C.C	12.7	12.9	11.5	# 49			(10-77	Plux		-
7.74	704	2	3 6	3	, ,	0.33		3	6 1		6.33	٥ ٧	6	2 2	8	6.10		6 0	5 97	5 97		5 89	ş				. 3			4	5.62	8	5.95	3.95	5.95	5.97	5 97		6.72	6.71	6.71	2.0	4.17		1 1	(10.5	Sundicted	-	Cabulated Flore
2 47	2.20	1.00) i	2 6					10.		2.16	1.8.7	8	2 5		2.71	1	2.63	۰ 4	2.37	2.37	2.37		204			3 05		1.37		3 29	12 88		2.70	2.78	2.93	2.14	6.13	0.50	5.95	619	3.38	4.32		W	(10.17	Emusion	,	The Penedica
1.98	2 58	2 38			,	3 61	3 /2		17.		2 95	4 08	4 10				;	3 62	3 28	3.27	3.27	3 27	3.78	307	3 62	4 15	4 26		1.92	1 58	101	4.01	3 98	3 75	3 87	4 08	1.99	7 63	7.73	7.31	7.59	4.06	7.17		Lur tua	10-	Епицалис	Spectral	

Table Z, M(4.8 µm) Photometry of Io(II) (continued).

						001	***	C7(T1()
96 L	09.8	199	13.5	85.9	3.04	991	749	87/11/75
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	≯€ ′\$		1.21	00.6	\$0 E	651	5211	87/11/23
67 L		LL'9				P\$1	1001	52/11/28
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99.9	*L *	6L9	\$11.5	8.86	01.E	WI		EZ/111/L8
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£0.8	09.0	LC S	οCT	% 6	L6 Z	302	61*1	82/60/78
05.4	TT'>	s?*	971	S96	00.E	767	LOS	82/60/LB
£\$ >	*! **	85 B	Ţ. <u>\$1</u>	9L"6	66Z	98Z	1221	82/60/18
85 P	B1.5	82.8	8.51	08"6	66.2	283		82/60/28
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to ç	16.2	42 L	t 01	928	LIC	7337	2141	90/80/91
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89.5	2.28	29 L	066	09L	9ZC	62	OFEI	50/to/91
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5.63	> 2	Z9 L	986	85 L	UC	Lζ		\$0/80/91
2.43	LO.2	79 L	69'6	** T	6Z.E	zz		50/90/91
8£ 2	2 03	79 L	s96	14.7	6Z.£	02		1 50/80/9
Lti	7:10	ES L	S96	6£ L	62.E	81	\$121	50/80/9
€9 €	80.€	68 L	0-81	Z*'W	\$1 c	+\$1	8s6	*0/80/9 !
€£`€	91 €	68 L	1.11	68'S	+1 ′€	75 1	\$\$6	w/80/9
66 E	8E.E	76 L	£.11	39.8	Z1 c	ost	S26	+0/80/9
94.1	TT.E	76.T	4.11	26.3	90 E	891	116	10/80/9
		000	2.05	LS 1	86'?	9+€	1408	0/20/9
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6£ 8	70.T	00.0	70.T	LF 3	LF L	.76	.50:	20, 30, 7
	ee: 3	co I	001	89L	3.25	6CC	21.0	£0/80/9
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	(101)	21.01)	4.017		Magnitude	AgaA.	(tru)	(TU)
1003	noissima	Mgiland	Xuf-			Phase	Jun.T	Date
Spiritage	Themsal	Reflected		-	1	Orbital	1	ļ.
Specital	MX Densilies	Calculated F		Observed		l	1	
							<u> </u>	

01 c W! LZ9 SZ/11/L8

Table 2. M (4.8 µm) Photometry of Io (II) (continued).

The C		CAMBINAL VIII	Galculated Fi)	Observed		l	l	1
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S9	F01)		Suntight		•	abuángaM	>y∄tr∨		
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LFC 26Z EOL %6 59L 97E 97E 9E11 E0/80/				4.04			.16	C=01	(A)eA
		26 Z	£0 L	% 6	59 L	. 92 €	PLE	9611	£0/80/

Table 2. M (4.8 μ m) Photometry of Io (J1) (continued)

Reflected Sunlight	-		<u>}</u>		Observed	d	Calculated Flux	lux Deasities	
(OL) Angle Magnitude (IQ) (Date	Tune	Phase			Flux	Reflected	Thermal	Spectral Emittance
1037 199 324 775 101 6.56 352 1025 206 322 773 101 6.56 352 352 1025 206 322 773 101 6.56 373 373 375 376 376 376 376 377 377 377 378 379 378 379 378 379	(0.1)	(04)	Angle (deg.)	Magnitude	(Jy)	(10 ⁻¹⁷	Wcm²,m²)	(10 ¹⁷ Wcm ² µm ⁴)	(10 ⁻² Wm ³ µm ⁻⁶)
1051 201 324 778 101 6.54 3.59 1122 206 3.22 793 103 6.24 3.78 104 8810 222 3.22 793 103 6.24 3.98 103 6.24 3.98 103 6.24 3.98 103 6.24 3.98 103 6.25 3.98 3.98 103 6.25 3.98 3.9	87/11/25	1037	199		7.75	101	6.56	3.52	5.01
1122 206 322 793 103 6.54 378 378 376 320 322 793 103 6.28 338 339 322 7185 102 6.28 339 339 322 7185 102 6.28 339 339 322 7185 102 6.28 339 339 322 7185 102 6.28 339 339 322 7185 102 6.28 339 339 322 7185 102 6.28 339 329 321 321 322 9.38 2.28 323	87/11/25	1051	201	3 24	7.78	0	654	3.59	
306 20 32 76 989 628 39 980 32 32 78 989 628 39 990 32 32 776 101 6.36 34 1100 46 317 831 102 6.36 374 810 51 304 31 103 645 437 807 51 304 934 122 938 278 837 55 303 942 123 938 285 837 55 303 942 123 938 285 837 66 304 936 122 938 282 851 57 302 947 123 938 285 852 57 302 945 122 938 282 851 57 302 945 123 930 243 852 57 302	87/11/25	122	206	3.22	7 93	10.3	6 54	3.78	
948 36 321 7,89 103 628 3.99 1100 46 311 322 7,85 102 6.36 3.86 1100 46 317 831 108 645 437 437 886 51 3104 934 112 938 2,76 887 55 300 927 112 12 938 2,85 887 57 302 947 123 938 2,85 887 57 302 947 123 938 2,85 1025 70 300 945 1125 950 2,42 950 1025 70 300 945 122 956 2,93 1012 79 300 946 122 956 2,93 1012 79 300 947 123 956 2,42 950 1025 70 300 946 122 956 2,93 1012 79 300 947 123 956 2,93 1012 79 300 947 123 956 2,42 956 1025 70 300 947 123 956 2,42 956 2,93 1012 79 300 948 125 956 2,93 1012 79 300 948 125 956 2,93 1012 79 300 948 125 956 2,93 1012 79 300 958 125 956 2,93 1012 79 300 958 125 956 2,93 1012 70 921 2,93 10 327 950 302 331 70 921 2,93 10 327 950 302 331 70 90 118 100 5,79 42 12 100 5,79 42 12 100 5,79 42 12 100 5,79 42 12 100 100 5,79 42 12 100 100 100 100 100 100 100 100 100	87/11/26	756	8	3 26	8	9 89	6.28	3.61	
948 36 324 776 101 6.36 3.86 948 36 324 776 101 6.36 3.74 8807 51 305 927 121 9.38 2.85 881 54 303 942 122 9.38 2.85 881 55 300 9.42 123 9.38 2.85 881 57 302 947 123 9.38 2.85 881 57 302 947 123 9.38 2.85 881 68 300 9.61 122 9.38 2.90 1002 68 300 9.61 12.5 9.50 3.02 1002 68 300 9.64 12.5 9.50 3.02 1002 70 300 9.64 12.5 9.66 2.91 1127 79 300 9.64 12.5 9.66 2.91 1127 79 300 9.65 12.7 9.66 2.91 1127 79 300 9.65 12.7 9.66 2.91 1127 79 301 9.58 12.7 9.66 2.81 1141 81 2.99 9.74 12.7 9.76 2.96 809 2.92 3.1 7.96 9.60 2.97 809 2.92 3.1 7.96 9.60 2.97 1015 9.0 3.0 3.0 9.6 10.4 5.79 9.1 3.0 3.0 9.6 10.4 5.79 9.2 3.1 7.9 3.0 9.0 9.6 12.1 9.6 2.9 9.6 2.2 2.9 9.6 2.2 2.0 9.0 9.0 9.0 9.0 9.0 9.0 9.0 9.0 9.0 9	8//11/20	30	:	i id	7.89	103	6.28	3 98	
948 36 324 776 101 6.36 3.74 1100 46 36 1100 46 317 8311 108 6.45 4.37 120 1100 6.36 317 8311 108 6.45 4.37 120 120 938 2.78 120 938 2.78 120 938 2.78 120 938 2.78 120 938 2.78 120 938 2.78 120 938 2.78 120 938 2.78 120 938 2.78 120 938 2.78 120 938 2.78 120 938 2.78 120 938 2.78 120 938 2.78 120 938 2.78 120 938 2.70 120 936 122 938 2.78 120 936 122 938 2.78 120 120 120 936 122 938 2.70 120 120 936 122 938 2.70 120 120 120 946 122 950 2.42 120 120 120 120 120 120 120 120 120 12	87/11/26	930	¥	3.23	7.85	10.2	6.36	3.86	
1100	87/11/26	948	36		7.76			3.74	36
866 51 304 934 122 938 278 837 51 305 927 1121 938 278 837 55 303 932 122 938 285 837 55 304 936 122 938 285 851 57 302 947 123 938 285 851 57 302 947 123 938 296 851 57 302 947 123 938 296 851 57 302 947 123 938 296 851 66 304 936 122 950 302 1002 68 301 961 123 950 301 1012 70 301 964 125 950 301 1022 70 301 965 126 966 291 1123 79 300 965 126 966 291 1124 81 298 98 127 97 30 29 1124 81 299 72 92 966 291 1124 <t< td=""><td>87/11/26</td><td>100</td><td>\$</td><td>3 17</td><td>3.3</td><td>10.</td><td></td><td>4.37</td><td>6.26</td></t<>	87/11/26	100	\$	3 17	3 .3	10.		4.37	6.26
807 \$1 303 927 121 938 270 834 \$4 303 932 122 938 245 831 \$5 300 942 123 938 285 851 \$7 302 947 123 938 298 851 \$7 302 947 123 938 298 851 \$7 302 947 123 938 298 1012 68 301 961 122 938 303 1025 70 301 961 122 950 242 1025 70 301 964 122 966 291 1121 81 290 940 122 966 291 1122 79 300 965 126 966 291 1121 81 299 974 127 978 290 1122 79 300 965 126 966 291 1121 81 299 974 127 978 290 1122 79 300 302 732 966 281 123 79	88/11/16	80	51	₩ 2	9 74			2.78	3.08
834 54 303 939 172 938 285 836 55 303 942 123 938 289 837 55 304 936 122 938 289 852 57 302 947 123 938 296 852 57 302 947 123 938 296 957 66 304 936 122 950 242 1022 59 301 941 125 950 302 1025 70 301 946 125 950 301 1025 70 301 946 126 946 291 1142 81 299 946 122 966 291 1142 81 299 974 127 978 309 1142 81 299 974 921 597 366 281 1142	88/11/16	807	~	3 05	9.27			2.70	2.99
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Table 2. M (4.8 μ m) Photometry of Io (J1) (continued).

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Table 2. M (4.8 µm) Photometry of Io (11) (continued).

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66			557	, 0				- ••	EU
66									
66									
66									
66									
66			_						
66	381	2 83	10 L	£8.6	\$5.T	UС	96	€€8	82/60/66
66		\$6 Z							
66°C CZ Z 089 CO 6 [69 9C C 6ET LEET \$7/60/66 EL C C 02 089 CR 8 \$L9 6C 9ET 9T C/ \$7/60/66 O6 C ST Z 0R9 96°R BR'9 LC PET 60E T \$7/60/66 TO C PEZ SR9 606 869 9C C LET PT ZT \$7/60/66 T6C 06: SR9 SL6 6PL 8Z C 9ZT 10ZT \$7/60/66	90 E	LZZ		?.26	60°L		801	6S6	82/60/66
66°C CZ Z 089 C88 8L9 6CE 6ET LEET 8Z/E0/E6 EL C 02 089 C88 8L9 6CE 9ET 9T C/ 8Z/E0/E6 06°C 51 Z 08'9 96'8 88'9 LEC FET COE 1 8Z/E0/E6 10°C 7 Z 58°9 606 869 9C°C LET 1 T Z 1 8Z/E0/E6									
667	16 C	06.	CR o	C / 6	69 L	97. c	971	1071	R7/50/56
667	10 £	> ₹ Z	SB 9	606	869	9€ €	221	15 14	#2/E0/E6
667 77 089 706 [69 988 681 LEEI 87/60/66 61 CL 7 089 688 84.9 666 961 91 C/ 87/60/66									
66°T 77 2 089 706 [69 9EE 6ET LEET \$7/E0/E6									
OC 316 629 PP 8 889 PP 6 186 909 66/10/16			089						
04.7	7 60	512	629	PP 8	87 9	44.E	183	909	62/60/66

Table 2. M (4.8 µm) Photometry of Io (II) (continued).

			7. Page 7	əldaT				
₱6 C	262	11-9	££ 6	LIL	EE E	10	627	92/10/66
95€	Z+ Z	109	£88	86.9	5£ £	86	914	92/60/66
12 E	5 43	11-9	78 8	6L9	6€ €	96	20L	97/50/56
S1 c	EE z	110	SL 8	TL 9	9+6	₹.	244	92/60/66
,,,	112	1 . 7	>C 5	CL 9	1			,
0L 1	1.26	\$£ 9	19 L	₽8 \$	SS E	552	Litt	\$2/60/66
28.1	⊊ € 1	PE 9	69 L	16 S	₽\$ E	482	***	52/60/66
161	₹ 1	0t 9	ī8 L	00 9	3 25	845	FEET	52/60/66
90 T	£\$ 1	07 9	E6 L	60 9	05 €	992	8151	\$2/60/66
Z8 Z	60 Z	£\$ 9	28.8	†Ç Q	€₩ €	238	1333	\$2/60/66
942	so Z	c* 9	L+ 8	16.9	£\$ £	967		\$2/60/66
\$6 Z	61 Z	9t 9	59 🖁	199	1 ♥ €	550	1013	\$2/60/66
56 i.	61 z	91-9	598	t9 9	1 ₽ €	218	100!	57/60/66
co E	\$24	61-9	EL B	019	9 €	807	\$58	\$2,60/66
10 €	EZ z	6t 9	14.8	699	0+ E	907	T £8	52/69/66
10 €	EC Z	0, 3	, , , ,	033	٠, ٢	700		30700.00
87.2	902	15.3	45 8	86 9	3 45	661	127	\$2/60/60
9L Z	2.04	159	SS &	LS 9	₹ 45	161	T £7	22/60/61
99°T	Lb I	t 5 9	848	150	€⊅ €	\$61	773	\$2/60/66
g1 ●	LO £	2 Z 9	≯ € 6	81 L	EE E	911	6411	80/20/26
96 €	16 z	00 9	1€ 6	\$1 L	EE.E	SZI	916	80/20/26
102	ШI	6 5 5	9E'L	505	65 €	966	1151	10/20/26
787	3.06	65 5	59 L	B8 &	75 €	333	8151	10/20/26
80 C	\$7 z	65 S	*8 L	ZO 9	₹\$ €	331	EOST	10/20/26
		19 \$	16 4	80 9	15 €	416	1351	L0/20/20
ti c	0£ z							10/20/20
6E C	80 Z	99.8	*1 8	52.9	84.€	60€	927!	LOILOILE
76 €	LR Z	99'€	€\$:8	SS 9	Z+ E	300	1151	10/20/26
6Z E	177	08s	12 2	[£ 9	L+ €	78 <u>2</u>	£\$6	10/20/20
† € <i>c</i>	tt Z	08.₹	\$ 74	EE 9	94 (082	\$06	10/20/26
L+ E	£\$:	6L.2	ZE B	6€ 9	SP E	574	1321	50/20/26
≯ 5 €	85 z	98.₹	£7 B	81 9	₹ (1	392	ETET	50/20/26
5.86	802	L6 S	SO 8	81 9	60 €	902	4601	50/20/20
		919	178	0s9	LFE	69	11:51	10/20/20
2 83	\$0 z	9[9	91 2			69	2061	10/20/20
LLZ	10 z		-	LZ 9	LV E			-
99.2	£6 /	80.9	00 8	\$19	6 † £	85	STT 1 Wet	01/10/04
8£ C	28 z	T\$.T	L 01	128	81 €	122	wet	01(10/06
LO C	95Z	88 L	> 01	20 8	12.6	512	6171	01/10/0
Z6 Z	** z	88 T	€ 0!	£6 L	22 E	213	1535	01/10/0
91€	19 Z	16 (901	018	61 E	202	1511	01/10/0
2.44	7:04	16 L	% 6	\$9 L	92 E	Soz	9811	01/10 /00
26.2	96'T	₱6 L	886	6\$ L	72.5	661	2501	01/10(00
L8 1	951	₽6 L	156	Of L	te e	461	LE01	01/10/00
61 T	951 C8 /	00.8	£8 6	SS L	72.5	L91	017	01/10/00
85 Z	91 z	00.8	201	08L	≯2°€	591	159	01110106
₩S Z ● '*1	91 Z 0'Zl	66 T	¥ 61	<i>08L</i> 6 ≯ 1	7C E	87E	1001	60/10/06
	911	EE L	6 21	5 7 1	95 Z	9#E	846	(QI10I06
8 E I	911	LE L	0 81	371	7» L	782	370	(4110100
0 ?1	L IT	EE L	0 61	9 1	2.55	WC	256	60/10/06
5 ● 1	zζi	EE T	5 61	051	2 53	345	916	60/10/06
0.21	921	97 L	8 61	2.21	15.5	6EE	106	60/ 1 0/06
	1 met me	/ m# 1000	('mu'msW			· · · · · · · · · · · · · · · · · · ·	1	1
	Wcm²µm¹)	('mas'moW	(101)	(At)		(189p)		
(, w#, wM	(10							
(, w ₁ , w _M	4.01)	ั ^ก ับได้	11.0 47		ShukingsM	⊃ίΣα∧	(TU)	(un)
(01ع	noissim3	Mailnuc	Xuf-	,	Staningsto	Phase Angle	ነጥነ (ፓህ)	Date (T(J)
			xn _{{-}	Opserved	obusinseM			

Table 3. Photometry of Io (11) at 8.7 μm_{\odot}

Spectral Emittance (10 ⁻³	(10:17	બ લ	Magnitude	Staff Staff	miT (TU)	oneQ (TU)
(''m4 ⁵ 'mW	Wcm ² µm ¹)	(Åf)		(deg.)		
z●l	SLL	961	60.1	340		W/ewes
88 9	79.4	111	65.1	15	ZZB	61/20/68
17 L	73.4	7.21	95°t	ES.	658	61/L0/E8
12 L	5L'P	Z Z I 021	≯\$"! 9\$"!	6S LS	*16 006	61/L0/E¶ 61/Lo/es
8 OT LT'L	68.5	c 21	ES 7	55Z 19	£09 £ £6	02/L0/E8
111	09°L 61°L	2.81 7.81	80'T 11'T	9CZ	109	02/10/58
6.51	29.8	8 12	160	253	SW	
671	591	21.8	160	524	808	02/70/£8
13 ¢	£6.8	LTZ	LX O	258	678	02/10/08
961	90 6	6ZZ	920	092	+58	83/07/20
Z >1	\$1.6	53.9	18 0	99Z	976	02/10/68
9 ● 1	\$7.6	9 O Z	8L O	69Z L9Z		02/10/68
L > (+8 .6	8 +7	LL'O	69Z	1000	02/20/68
0\$1	o 0 1	c \$7	SL V	OLz	1001	02/1/0/€8
8 36	ιςς	6 Cl	0>1	PL.	925	12/10/28
● /8	1? \$	L c1	79 t	SL	SCS	1 2/10/28
62.8	15.5	6 Cl	01	6L	8SS	
LE 8	95 S	o † 1	6E T	28	619	1 2/40/68
97.6	\$19	5 51	\$ 2.1	68	SOL	12/10/68
106	865	15!	18.7	76		83/01/2 !
92.6	\$19	s \$!	1.28	96	6 \$ L	12/10/58
48 B	88.2	₹● 1	CC 1	601		1 5/F0\E8
10.6	86 S	151	131	111	576	17/10/68
031	· · · ·	3 ••	230			20,00,00
691	cii	s #2	290 290	815	028	TO/80/E8
691	£.11	s 82	290	350	558	TO/80/E8
8 9 î 1 L î	Z II ● 1!	C #4	[9 o 19 0	322 322	158 \$ 98	70\80\E8
99!	111	C 82	10.0	257	116	TO\80\E8
		0.07				
CL L	*! *	S 01	14.1	● z		72/80\E8
Z£ L	3.92	66	LL'I	LZ		LZ/80/C8
78.8	\$1.9	951	1.28	EII	1542	60/so/w
£16	95.8	19!	124	*!!	1522	
968	9 7.9	8.2£	9Z I	\$11	EDF 1	80/50/#8
10 6	0€ 9	6 \$ 1	1.26	130	PEE 1	80/50/98
606	EC 9	091	1 25	.71	1348	80/50/18
188	919	s 🕻 l	1.28	121	1452	80/50/#8
168	6.20	L S1	LZ"1	156	BEPT	84/02/08
Loc	f 0Z	€ Is	ZO: O-	66Z	+18	L1/90/#8
200	·	0 67	210	308	y>* 1	L 1/70/18
402	0 L!	o £≯	/1: 0	Los Sos		L1/90/#8
102	991	c 8E S It	050 Ozo	216	996 016	∠1/90/≯8 ∠1/90/≯8
€ £1 s®1	7•1 7:51	98.0	Lξ 0	91 c	\$101	L1/90/#8
÷ 91	5 c1	340	(† o	350	9101	L1/90/#8
951	671	35 s	89 0	354	L111	L1/90/#8

Table 2. M (4.8 µm) Photometry of Io (II) (continued).

(10. Emiterace Specitis	Incontrol Themat Emission	Calculated F Reflected Sunlight	¥nf∓	Observed		IstichO Searly	Len'	sieG (FU)
('m4'mW	(Long ² ma ³)	(101) (ใกณ ² msW	101) ('mu'moW	(kf)	Magnitude	(72p)		(10)
ES c	29 z	6 Z 9	16 8	• S •	35€	282	059	62/60/66
so E	9Z 🕻	6Z9	9s1	L\$ 9	Z* c	L8 Z	SOL	62/60/66
18 C	∠8 Z	029	LO 6	L69	9€.€	262	967	62/60/66
60 7	90 €	02.9	C: 6	60 L	• 1€	\$6Z	601	62/60/66
£\$'\$	LE C	+ 1'9	15 6	0E'L	3.31	10C	l w	67/09/66
15.	SE c	† 1 9	64 6	6ZL	t€ f	€0€	468	62/60/66
4 52	91 c	*! 9	0C6	9! L	cc £	016	226	62/60/66
LL'*	₱\$ c	60 9	£9 6	6£ L	6: C	116	* \$6	62/60/66
L1'S	₱8 c	609	76 6	79 L	97 C	350	ES01	62/03/56
89.4	8t c	so 9	6 23	7.E L	1€ c	355	2111	67 460/66
€6 ≯	99€	509	14.6	94 L	₽Z c	725	trii	62/60/66
L9 🕨	LP'E	\$0.9	55 6	157	16.6	3.59	4511	63/03/56
9Z Þ	91 €	90 9	2Z 6	80 L	₱£ C	338	2421	63/03/56
ተ ር ተ	22 E	90 9	82.6	CI L	€€ c	131	t571	62/03/26
90.4	ζ0 <i>c</i>	90'9	806	L6 9	9€ €	338	8061	62/ 60/66
98 €	78 :	21 9	66 8	069	LC €	146	1351	63/60/66
₹30	02 c	Z1 9	18 6	SIL	€€ €	343	Ott,	62/60/66
9 €	252	719	198	199	l⊅ C	342	PSE1	62/60/66
99C	ZL'Z	219	18 8	6L 9	6E E	746	30t1	62/60/66

I age 1, E aldaT

9.03	7.67	19.4	2	2	1423	85/07/10
8.86	7 53	0 61	. 8	67	1417	85/07/10
9.03	7.67		2	S	1359	85/07/10
8 16	6.93	175	1.15	4	134	85/07/10
7.79	6.62	16.7	3	¥	3	85/07/10
10.8	0.81	17.2	111	24	8	93/0//10
10.1	6.81	17.2	1.17	ß	858	85/07/10
10.0	8 48	21.4	0.93	\$	122	85/07/09
	6.87		1.16	\$	1259	\$5/06/10
9.13	6.62	6 9	1.19	2	1249	85/06/10
9.21	6.75	17.0	1.1	63	1239	85/06/10
9.21	6.75	17.0	- :	ŝ	1229	85/06/10
å	0.21	5.7	1.27	0,	ć	01/00/08
	5		3 5	: :	,	0.000.0
	?		5 5	<u>.</u>	3	85/06/10
. 0	7 47	je :	07	2 !	3	85/06/09
	8 21	3	0 97	31	305	85/06/09
-	%	8	0 92	225	1256	85/06/09
11.1	9	20.4	0.98	224	1247	83/06/09
			2 5	1		60/00/00
	7 80		2 5	3 5	3	84/26/09
• 3	7 2		5	3 5	,	60/00/09
8 9	2 :		5 8	, ,	1	30,70/33
786	577	4	- K	¥	1426	
8.12	0.88		ij	<u>.</u>	1402	80/00/08
	3.92		1.91	: נ	i	80/00/08
= -	\$	5	Ş	3		#3/uo/u/
11.2	# O#	20.3	9	0	101	85/06/07
13.5	9.86	24.4	0.79	335	1734	85/06/06
•	:	•			į	
13.5	9 66	÷	0.79	334	1249	85/06/06
13.5	9.66	24.4	9	333	1242	85/06/06
12.4	8 84	223	0.89	327	1202	85/06/06
11.9	8 52	21 5	0 93	326	1153	85/06/06
7.71	5.51	13.9	ŧ	8	101	84/08/20
3				,	,,,	07,00,40
ž ,	7 2		2 :	֓֞֞֞֜֞֜֞֜֓֓֓֓֓֓֓֓֓֓֓֓֓֓֓֓֓֓֓֡֓֞֓֓֡֓֞֓֓֡	3	84/04/20
9.54	6.03	3 6		8		41/00/10
2.0				Ē		41,000,40
3 2 2	9 9	3 3		ž \$	3	84/08/19
\$	X 93		<u>.</u>	ž	2	84/08/10
9.48	6.81	17.2	1.17	157	7	84/08/19
1.89	6.38	6.1	1.24	152	£	34/08/19
=	¥ #	37.3	0.33	33	3	84/07/19
.	7.			33	=	84/07/19
3 0	14.7	37.1	9	327	727	14/07/19
18.5	5		10.01	324	2	94/0//19
			2 6	;		14/07/10
	70		\$ 8	3 5	3	84/06/17
	-		\$ 1	3 8	•	84/06/17
<u> </u>	13 1	ب ب	ĵ	3	24.	34/06/17
	, ,					
Wm ',m')	Wcm ² ,m ³	(yc)		(deg.)		
(10-	(10.17		Magnitude	Angle	Ç	(10)
Emittance	Density	Flux		Phase	Time	Date
Spectral		Contrago		0		

Table 3. Photometry of io (J1) at $8.7 \mu m$

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Orbital	
Observed	ī
₹	
Spectra	
_	

22 25 25 25 25 25 25 25 25 25 25 25 25 2	382	314 063	314 0 65 315 0 62 319 0 62 323 0 66 304 0 69	314 0 63 28 3 3 3 3 3 3 3 3 3 3 3 3 3 3 3 3 3 3	0 65 274 255 268 268 268 268 268 268 268 268 268 268	314 0.53 23.6 315 0.62 23.6 315 0.62 27.4 302 0.66 27.4 303 0.69 26.8 316 0.75 26.3 316 0.75 26.3 316 0.77 26.3 316 0.77 26.3 317 0.41 34.7 318 0.42 34.3 318 0.42 34.3 319 0.43 31.1 340 0.53 31.1 340 0.53 31.1	314 0.65 22.6 315 0.62 23.6 315 0.62 23.6 315 0.62 27.4 302 0.66 27.4 302 0.69 26.8 304 0.69 26.8 316 0.70 26.5 316 0.77 26.5 316 0.77 26.1 317 0.41 34.7 318 0.42 34.3 319 0.43 31.1 340 0.53 31.1 340 0.53 31.1 344 0.66 29.1
(dy) Wo(1) 22.6 23.6 24.5 25.3 26.2 27.5 28.6 27.5 28.6 27.7 28.6 29.1 29.1 29.1 29.1 29.1 29.1 29.1 29.1	21 8 8 22 0 8 19 6 7 26 2 7 28 3 11			28 6 25 3 27 4 26 8 26 8 26 8 28 5 28 5 28 5 28 5 28 5	22 24 25 25 25 25 25 25 25 25 25 25 25 25 25	22.2 22.2 22.2 22.2 22.2 22.2 22.3 22.3	28. 11. 22. 10. 26. 26. 10. 26. 26. 26. 26. 26. 26. 26. 26. 26. 26

Table 3. Photometry of Io (11) ref 8,7 µm (continued).

15.21 9.41 002 Z **L**Z 00.1 611 SZZ 91/21/6S £6.7 128 SOL 91/21/68 8 01 911 L9 O 011 1 c1 6 L Z 26/ 15/121252 343 S90 01s1 \$1/21/68 o 1? €61 z 91 Szo *c 8 61 9.91 42.0 02'0 338 1591 51/21/68 9 **L**I 6 OZ. £.44 11.0 335 8011 51/21/68 89/12/12 1243 14 333 9 **2 i** 80.0 330 24.3 102 511 S9'6 6L'0 6111 • 1/21/68 ***11** LS'6 74.2 08.0 9701 • 1/ZI/6S 8.01 0s2 \$617714 1008 5.01 64.8 222 20.1 48 116 +1/21/68 680 001 78 158 +1/21/68 056 LL"L 418 9 02 **L6**'0 \$6/15/14 133 EL 6 οć 217 41/21/68 L"OZ 460 £.01 0'92 06+1 21/21/68 cZi 1.21 1.01 25.5 PL'0 89/12/12 1344 911 972 ш 1001 21/21/68 99.6 64.0 L'*Z L'11 **8**4.6 04 L*ZI 71/71/68 8L0 2211 21/21/68 901 n o **8** 6 Zoz ● 0Z 1.00 89/12/17 1020 79 6 60.8 860 4 1001 Z1/Z1/68 916 21/21/68 9SL 911 101 20 L of-s 611 ££.1 86 006 21/21/68 1901 21/10/68 971 9 01 t 92 690 60€ C.11 5 82 **29** o **∠**0€ 2701 21/10/62 t c1 671 8.01 CLZ **19** 0 106 C96 \$1/10/6\$ 1.21 z 0! LSZ CL O 662 876 81/10/68 9 61 • 11 8 82 19 o 167 OER 81/10/68 2.21 082 290 LIZ 906 81/10/68 0.11 \$6.6 1s2 8211 91/11/88 9s11 91/11/98 2.11 1.01 ?s2 *****LO 5.25 9101 91/11/88 1.11 SL 0 1.11 1.01 \$4.0 ● 101 91/TT/\$\$ 1.25 91.0 1001 91/11/88 0.11 79.Z 0.11 0001 91/11/88 19 PP6 91/11/88 90! S5"6 w 'o 901 85.6 z⊕Z MO 216 91/11/22 Uo 080 8.45 196 91/11/22 18.6 601 901 09.6 9s8 91/11/88 \$\$8 91/11/99 10.7 744 LS 13.1 28.0 9SL 91/11/81 1 01 **L16** 6Þ 6Þ ['01 \$1.6 73.1 SSL 91/11/81 SE 0 Ltt 61.8 L 07 £6.0 LP 6011 97/11/41 001 956 97/1 1/LE Ⅱ £6.7 0-02 LE 666 97/11/4 5.11 **E 0Z** 66.0 SE co 1 (,.w√,.wɔ,∧ (k_f) (-Zap) (, w_z.w_M Magnitude , O1) algaA (TU) (TU) Emittance Flux Density Phase Date Tume Distal Spectral Observed

Table 3. Photometry of Io (I I) at 8.7 µm (continued).

Spectral	p	Observe		D. Orbasel		
Sommen	x Density	пЫ		Phase	SauT	Date
(10°² Wea ¹ µm'¹)	(10 ⁻¹³) Wcm ² µm ³)	(A)	obuknysM	348nA (-83b)	(TU)	(TU)
	/ mm/ ma			,		
2.01	69.8	6'12	06.0	61	1552	\$0/80/9
z 01	co.s	\$ 12	160	tζ	1921	50/80/9
1 01	82.8	7 12	Z6 O	23	1526	\$0/80/9
001	• St	912	260	28	EEET	\$0/80/9
L 66	69 t	212	£6 O	15	BPET	\$0/80/9
88"6 62 6	01 1 11 t	Z:1Z	n o n o	K EE	TE+1	\$0/80/9
				.,		
\$6'6	67 \$	912	S60	0+	ES>1	W/m n
8.11 1.01	1 01 65 E	52 ¢	SL 0 Z6 0	Coz 1♥	1101 SOS!	90/80/9
8:11	101	72.5	*L'0	902	6701	90/80/9
8'11	1.01	52 2	* L 0	208	5901	90/80/9
1 c1	. !!	r sc	390	at.	8171	90/80/9
s c1	5:11 5:11	787 2813	09°0 S90	243 238	OS* I	90/80/9
o 21	s \$1	1 6E	82.0	0s	8001	L0/80/9
7 B1	9\$1	₱ 6E	LZ O	95	6601	70/80/8
z % [951	₹.6€	۲Z:0	9s	ZSOI	LO/80/9
∠ \$1	191	9.04	> 2 0	08	OPE T	70/80/2
7.81	091	\$ 07	9 ₹ o	28		LO/80/9
9 81	651	co?	Szo	48	87+1	L0/80/91
0 / 1 c \$ 1	5 51 £ \$₹	1 6E 1 6E	97 o	282 6t	1500	87/60/LS
	_					
69!	2 \$1	6 85	82.0	782		87/60/2
† <i>L!</i> 991	L \$1 z \$1	5.9E	LZ` 0 0s0	10E 56Z		87/60/L
* 9 !	0 \$1	31.6	16 0	80s	2051	82/60/L
s 91	151	o T E	16.0	309		8Z/60/L
10 L	173	271	a c 1	711	217	LUI I/L
8 ∲ ! 16 <i>L</i>	901 £9:\$	L 92 z • 1	69°0 2€ 7	∠11 511		62/11/ <i>L</i> 1 1/ <i>L</i> 1
0.51	L 01	OLZ	\$9.0	133		1/11/53
8+1	'901	L 92	690	128	60L E	Z/ 11 /L
6'91	901	R 9Z	690	PET	\$7R i	Z211 /L!
271	9 o!	L 9 Z	690	9E1	048	EZ/1 1/L
s +1	♦ 01	z 92	11.0	S¥ i	896	EZ/1 1/L
C†! 9†!	7 01 • 01	LSZ z 9Z	CL O IL O	091 551	TE11	EZ/11/L Ci/11 /L
L 101	9 01	76.2	IL O	121	EIS	\$7/11/1
0 71						
z S I 6 † I	L 01 S O I	1 LZ 9 9 Z	89 O OL O	193 128		\$7/11/L! 7/1 /L!
12 4	801	C LZ	L9 0	591		Z/ 1/L/
9 c1	SS 6	I ● Z	080	002	**01	52/1 1/21
9 E i	SS 6	1 92	080	707	s s01	\$7/11/15
z c1	82 6	23 👁	£8 o	`Zôz	6211	52/11/1/
101	69 L CO L	C 61	sol ∳1`t	15/		2/1 1 /L8 2/1 1/L8
601						

6S2

c 01

140

OCL 91/71/68

Table 3. Photometry of Io (11) at 8.7 µm (continued)

69.8	\$1 .8	c 91	1.23	15	998	92/20/26
07.8	9 12	c 91	1.23	61	££8	97/50/56
CI 8	10.9	c St	oc i	EÞ.		97/E0/E6
108	?6s	0 \$ 1	ZE 1	19		92/60/66
[L L	ILS	701	9s′1	60	+ZL	92/60/66
12.7	ÞE S	S.E1	EP 1	Ĺξ	OIL	
					· · -	, 4. 20. 20
♦ € <i>L</i>	99°S	L C!	1) 1	SE	959	97/50/56
45 6	60 L	6'11	£1.1	322	1791	52/60/66
\$2.6	28.6	£ 41	1.16	6*Z	1340	\$2/60/66
\$6.6	BC L	981	90°t	247	1327	\$Z/ED/E6
U 6	72 L	E.E.1	01.1	6 a		52/60/66
	,,,,		•••	o u	002	30 00/20
94 6	CCL	£.81	01.1	La	9121	\$2/60/66
97.6	9 t 9	£.71	91'!	122	1201	\$2/60/66
\$ 0.6	14.9	061	21 1	612	9001	\$2/60/66
L9 \$	2 + 9	16.2	1.23	506	206	\$2/60/66
72 2	\$\$.6	5.01	12.1	Loz	99t	
						20.00,00
99°B	099	ፒ 91	I 3¢	300	8 \$4	\$2/60/66
L?	£ • 9	791	£Z 1	661	SPL	\$2/60/66
£E.8	41.9	9 🕻 १	128	161	167	\$2/60/66
z 01	LV L	681	LOI	121	\$76	80/70/76
6/ 6	LIL	181	11'1	BEE	9551	LO/ZO/Z6
u_ 0	• • •		• • • •		,,,,	20,00,00
61 '6	21.9	041	2 1 /	755	9261	10/20/26
0.11	60.8	70.4	86.0	334	1254	10/20/26
sli	CP 8	EIZ	+6 '0	255	1151	10/20/26
2 C!	101	9 \$ 2	• L 0	615	4551	10/20/26
9 € 1	LOI	072	t90	310	123¢	10/20/26
				• • •		
v ?1	L 01	69Z	89.0	305	SET	L0/20/26
L Z1	Z£ 6	5 62	68.0	222	1001	10/20/26
611	7L'8	1.22	060	282	916	
021	\$£'\$	1.22	060	575	1329	\$0/70/76
6'11	69.8	2.12	[6 O	69Z	1251	\$0/70/76
	•, •		•			20.00.00
101	TE.T	9.81	\$0 T	L+2	1101	\$0/20/26
SO 6	9s"9	991	121	OL	ISEL	
L6 T	0s"9	P'9!	22.1	69	LPE	10/20/26
LE 8	£9 9	Z-91	1.23	w	60E1	10/20/26
52.8	21 9	9.61	1.28	6s	1533	
\$1 '9	•! '\$	0 £1	27 '?	223	**£1	01/10/06
011	61 '6	7 77	18'0	122	6261	01/10/06
8'01	106	\$.22	LT O	216	1525	01/10/06
1'01	90'6	6"ZZ	9to	213	SEZI	01/10/06
601	60.6	0 EZ	9to	300	1122	01/10/06
1'11	LZ.6	* EZ	£1 o	302	6611	01/10/06
8.01	90'6	6'22	98.0	661	SSOI	01/10/06
201	20.6	1.22	9to	L61	1901	01/10/06
E.11.	17.6	\$ 62	28.0	WI	EIL	01/10/06
011	9.20	2.02	9to	591	'*S9	01/10/06
					_	
4	('ma''ma)W	100		(.0)		
(1.m45.mW	(10.)	(46)	Magnitude	(- 8 >p)	,	
,.01)		1		Man	(TU)	(TU)
ววกงาหกรี	ykened z	MH.	L	Phase	>m1	Date
Spectral	P	Observe		IssetO	l	
	L			<u>. </u>		

Table 3. Photometry of lo (11) at 8.7 µm (continued).

902	CL!	9 64	91'0	187E		60/10/06
c 17	8 / 1	6 11	c! 0	LPE		50/ 10/06
i. 17	t LT	6 **	cl o	344	9C6	60/10/06
7.12	181	8 SP	110	345	126	60/10/06
1, 22, 1	s 81	L 99	too	0+€	\$06	50/ 10/06
s £ Z	961	96V	70 o	335	018	60/10/06
6CZ	Ooz	90s	00 0	LZC	EE4	60/10/06
0 \$ 2	1 02	8 02	10 W	SZE	614	60/10/06
1 11	CC 6	9CZ	E8 0	991		90/ 10/06
0.11	CZ6	c £7	78 0	ŵi		90/10/06
2.11	** 6	23 8	78 o	651	7* / 1	90/10/06
£:11	89.6	53 9	180	751		w/ 10/06
5'11	59.6	? • Z	6L0	LÞ1	1073	90/10/06
₽ 111	SS 6	I ● Z	Oto	Z* 1	L\$6	90/10/06
cff	94 6	6 a	18.0	LET	L06	80/10/06
2.11	£† 6	8.62	28.0	SE 1		80/10/06
zli	Z⋫ 6	8 CZ	78 o	121	008	80/10/06
2.11	ZÞ 6	t £Z	ZS o	521	S*L	90/10/06
601	9.20	23.2	98 0	II	629	W/ 10/06
011	€Z 6	CR	PB 0	211	919	80/10/06
Ooz	8 91	\$ 23	6! O	ccc	WC 1	40/10/06
20.3	141	1 69	41.0	CC		LO/ 10/06
012	941	5 99	€10	358		LO/ 10/06
	S 8!	L'99	60'0	● ZC		LO/ 10/06
33 0		2 22	90'0	416	* \$11	
0.62	€ 61	X IV	PO 0	416	7511	20/10/00
22.3	881	* L*	L00	916	LFIT	L0/10/06
* EZ	961	9 6?	ZO 0	115	•	20/10/06
Ocz	C 61	8.84	PO.0	60s	8501	10/10/06
23 o	961	96+	Z00	96Z	LZ6	LO/10/06
c £Z	9'61	● 6≯	Z00	≯6 Z	16	LO/10/06
c ZZ	2 21	• 4	LOO	682	1+8	40/10/06
\$12	1 11	LS*	11.0	182	928	10/06
122	9 21	6 99	30.0	SRZ	119 /	OI1OICX
z 1 Z	8.71	6'99	c1 o	612	728	40/10/06
LOZ	+41	6 EÞ	\$10	LLZ	217	40/10/06
£.61	z 91	o I*	Czo	692	419	40/10/06
	1.81	L'0+	\$7°0			
ī 61				292 120		0/10/06
9 ! !	U 6	L DZ	8L 0		1771	'w/lo(of
L 01	96.8	9ZZ	L \$ 0	* [1	8511	90/10/06
121	101	25.5	₽L O	12	\$28	90/ 10/06
s 1!	89.6	• • Z	64.0	6L		90/10/06
2.11	8£ 6	1 EZ	28 o	CL	217	90/10/06
2 11	∠ ₹ 6	6CZ	18 0	IL	829	90/10/06
\$°11	19 6	● \$2	6L0	9+1	4501	91/21/6s
z 11	L+ 6	6CZ	It o	SET	126	91/71/68
	Wcm ⁻² km ⁻¹)	(4()		(·8ap)	[
('m4'mW						
*01) (¹ -mң ² -mW	41.01)	""	Magnikude	31 7 6V	(to)	(10)
(10.5		L	ShuangaM	otan^ oign∧	(TU)	(TU)
	Consity	L	Magninude	issidiO əssiff əigaA	Time (TU)	asa (TU)

Table 3. Photometry of lo (11) at 8.7 µm (continued).

				بسميب	_	_	
1	\$6:11	88.8	**22	no	9+6	1405	62/60/66
	PP'11	02.8	5.12	66.0	** E	L PE1	62/60/66
	9 t h	158	512	C6 O	ZPE	0EE [67/60/66
	65.11	19\$	1.12	76'0	940	SIEI	62/60/66
	56 11	\$77.8	022	060	338	1001	62/60/66
ı	\$9"11	05 8	\$ 12	6.93	966	6921	62/60/56
ı	15 41	22'6	E.ES	18.0	Осс	1304	6Z/EO/E6
	9C Z1	\$1.6	23.2	58 o	328	1 125	63/03/56
	15.84	95'6	34:1	08.0	373	1 150	6Z/EO/E6
	0621	SS'6	2+2	0s"0	1 ZE	201 T	62/60/66
	12.80	15'6	0 ≯Z	18.0	312	7001	63/60/66
ı	L621	S96	24.3	64.0	311	696	62/50/56
	98 21	\$\$ 6	74.1	080	90E	106	62/00/06
ı	98.21	ss'6	74.1	0s"0	Zoc	m	63/03/56
ı	/ m=/ m	Mcm.'mm')	((4)	l	(1900)		
ı	⁽¹⁰⁻¹) (10-1)	(10-ري	(11)	ShukingaM	(Japy)	un	(44)
ı	Sommen	укизод и	NJ.		Staff	Тыте	Date
ı	Spectral	tz.	Observe		IslidiO		

Table 3. Photometry of lo (11) at 8.7 µm (continued).

Special Enitance (10°)	tx Densky		obuningald	District Office of the Control of th	anuT (TU)	эь.a.c (TU)
War's part)	Mcm.,hm.,)	(44)		(-2op)		·
6 8 t	65.9	991	12.1	?s	* E6	92/ E0 /E
L6 2	S99	191	1.20	99	9s01	97/50/5
S5'6	80.7 87. a	6 41	£1.1	9L		92/20/20
\$16	8 7. '9	1.71	41 1	98	csz i	97/60/6
926	989	CLI	91 T	88		92/60/6
616	189	Z 41	LII	26	1581	97/E0/E
02.8	0s9	6 \$ 1	52 1	\$6	7291	92/60/6
S6 B	19.9	8 91	02.1	100	SEP!	92/20/2
026	E8 9	z Lī	L I'T	462	150	LZ/E0/E
+ 0′6	OL9	691	61'1	239	6S9	LEIEDIE
6! 6	Z T 9	ZLI	L1'1	6S2	512	LZ/EO/E
71.6	U 9	161	81 1	747	6ZL	12/60/6
996	LIL	1 81	111	* Z	208	LZ/EO/E
66.01	14 L	S 61	1.04	248	628	LZ/E0/E
98 01	80.8	Coz	660	552	L06	72/60/6
11.03	\$1.8	L'07	L6.0	LSZ	196	LZ/EO/E
12 11	ZE 8	012	s6 O	Z9Z	cco	1/2/60/6
18.11	9L 1	77.1	060	OLz	9101	12/ED/E
5021	≯ 6 g	922	L\$ 0	142	1504	LZIEDIE
12 34	\$1.6	1.62	\$8.0	282	RIC !	72\E0\E
28.11	67.8	Zzz	69 0	167		LZ/ED/E
66'11	68 2	5 7 Z	88 0	562		LZ/EO/E
L T 11	18.8	77.7	610	L6Z		LZ/ED/E
c6 1 I	28.8	22.3	6 8 O	10£	£1+1	LZ/E0/E
co Z!	€6 t	22.5	88.0	303	967 I	LZ/EO/E
• T Z !	00 6	L 22	18 0	Coc	0791	
8 82	cc 9	091	1 32	18	509	82/60/6
8 16	ſšô	zL!	L! "!	83		82/60/6
11'6	9 L 9	1 L!	81 1	58	SIL	87/E0/E
190	B1 L	. •	•••	00	7*L	aL/LU/L
6L6 196	LZ'L B1 L	F 81	111	06	6SL 9†L	87/60/6 87/60/6
01 01	OCL	681 E81	W / 01 1	S6 16	7 <u>58</u>	87/60/6
LO 01	87 L	611	LOI	L6	078	82/60/6
10.24	09L	Z 61	so !	101		82/E0/E
	, C _	-				
≱6 6 t L 6	8£ L 9Z L	981 C#1	01 [111		82/E0/E
88.6	££ L	s % (981	60 ! 100 î	721	6021	82/E0/E
\$\$ 6	60L	6.71	113	621	1555	
L£ 6	% 9	9 L I	21.15	561		82/60/6
810	Cas	LU.	L. 1	LLI	ect 1	Bartort
88 8 81 6	6S 9 Z S 9	991 Z <i>L</i> I	17 T 11 T	a l LEI	SPEL	\$7/60/5 \$103\58
0s 21	82.6	73.4	£8 o	w:	119	67/60/6
17.17	10.6	22.8	98 0	286	859	67/€0/€
Z\$ 11	SS 8	9 12	76 O	7882	217	67/60/6
300.				-{		
88 SI	9S6	747	w o	. 667	EDL	62/60/64

Table 4. \times (10 μ m) Photometry of Io (J1).

												_													_																										
84/05/08	84/05/08	84/05/08	83/08/07	83/08/07	83/08/07	83/08/07	83/08/07	83/08/07	83/08/07	12/10/28	83/07/21	83/07/21	83/07/21	83/07/21		83/07/71	83/07/21	83/07/21	83/07/20	87/07/20	83/07/20	83/07/20	83/07/20	83/07/20	83/07/20	83/07/20	83/07/20	83/07/20	83/07/19	83/07/19	83/07/19	83/07/19	83/07/19	83/07/19	83/04/24	83/03/11	11/03/11	83/03/06	83/02/18	83/02/18	83/02/18	83/02/18	83/02/18	83/02/18	80/04/12			Ŝ	Date		
1331	1259	174	916	Ę	849	837	ž	816	\$	ş	2 2	3 2	2	203	:	4	553	529	8	8	934	923	857	8 \$	812	759	8	80	935	910	857	168	820	313		1506	1504											ŝ	Tune		
121 /	1191	Ξ	328	326	323	320	916	-	316	5	§ 5	5	2 /5	87	:	<u>.</u>	78	74	271	360	265	264	260	<u> 1</u> 58	254	252	237	226	సి	ş	8	5	<u>.</u>	æ	36	193	10	339	ŧ	339	336	332	321	326	ĸ	1	G S	Mase	0		
0.03	004	8	0 21	622		0.21			ii ii	9	\$ 5					0 1	0 5		6 6 6 6		-0.21		-0.20	0.20	6 %	0.16	9	÷0 0		0.24	0 26	0 27	0 21		023	0.24	0.25	-0.17	015	0.14	0.14	0.12	0.05	0 07	0 25		waganoc	N. I.			
37 5		265	8 8	\$	44.7	\$		8	473	9			. A	34.5			3 6	<u>ي</u> ن	50 8	^ 84	\$ \$		80	\$ 4	± 	44.7	9.1	<u>+1.5</u>	3).2	\$ 6	20	20	29.1	29.8	31.2	4.2	9	\$	33.6	33.9	33.9	34.5	32 .	X	8		Ş	[ū	Observed	
1.07	8	<u>.</u>	33	3	1.27	1 33	1.33	1.33	1.34	1.02	9 0	98	0.990	0.981	ì	0 946	0 954	0 946	44	- - -	1.33	1.38	1.32	1 32	: 26	1.27	1.19	1.18	0.886	0.878	0.862	0.854	0.847	0.847	0.886	1.37	1.39	1.23	0 954	9.963	0 %3	186.0	8	107	0.870	Acm, way	(101		Flux Density	<u> </u>	
1.53	1.52	49	1.20	3.20	2.11	2.20	2.20	2.20	123	Ę	1.52	.49	. 49	48		5	1	1.42	216	?	%	206	1.97	1.97	: :	1.96	1 78	177	1.32	1.31	1.28	1.27	1.26	1.26	1.26	2.#	2.48	2.35	1.93	1.95	1.95	.98	2 E	2 2	22		Wm², m¹)	(10·	Ematance	2	

Table 4. N (10 μm)

of lo (J1)

85/06/09 1 85/06/09 1 85/06/09 1 85/06/10 1 85/06/10 1		-			-		_		_		15/06/09	•	85/06/08 1	\$5/06/01	85/06/08	_	85/06/07 L	85/06/06	15/06/06	-	-			84/08/20			84/08/19			_	84/08/19	-	•	•	84/07/19			Ξ.	_		14/05/08	<u> </u>			7
1209 61			-			1241 223	1224 221	1217 220	1203 218				1352 30		1240 20		1057 162			1236 332			_	347		\$27 166	100				643 151		_		716 325	77	653 322			421 126	122	(deg.)	5 5	Orbital	-
		9 07			ь,		0.24	0.24			0.24		0.03		8		 b	0.32				6		6.14	. 6						607	-			6					9		C 7			-
-	37.8	<u> </u>	38.2	4		=	±	<u>+</u>	47.9	Ė	<u></u>	37.3	37.8	38.3	38.6	49.0	£	52.0	51.7	52.5	\$	43 6	33.9	\$5.2		3.4	2.	. 9	3.1	5 .1	Ê	67.0	67.8	8	2	67.7	68.7	2	2	25	37.2	(ענ)	Flux	Cosciver	²
,	<u>8</u>	3	.09	38	- : - :	2	1.37	1.37	1.36	1.25	1.37	ĩ. 8	1.07	-98	. 10	1.39	- 1	.4	1.47	1.49	1.32	1 24	0.963	5 5	1.27	1.23	1.22	2 2	1.22	1.22	1.17	1.91	1.93	. 3	1.93	ร์	1.95	*	<u>.</u> 04	10.1	.0	(10 ⁻¹⁶ Wcm ⁻² µm ⁻¹)	10		
8	- E	- B	1.48	8	7	1 89	- 8	1.	1.87	1.72	- 82	\$		1.50	i. 52	1.94	.ī. 28	2.07	2.06	8		74	25	2.70	1.77	ב	5	1.74	7	7	2	2.34		2 32	237	٠ ۲	2.40	. 75	- 49	1.5	34	Wm²μm¹)	(10°	Spectral	

Table 4. N (10 µm) Photometry of Io (11) (continued).

,,_	07 -	011	SL.0-	007	6024	82/60/18
i≯こ S≯z	07 z ≯7 z	9 LL L'8L	LL 0	98Z 283		\$5/60/L\$
15 2	67 Z	7.08	08.0-	182	5511	82/60/28
11.2	78 1	6S9	≯ \$`0-	65Z	\$501	80/80/98
612	88 I	5.33	86.0-	L 8		L0/80/98
		• //				
Z1Z	L8 !	9s9	72.0-	56	81 ● 1	LO/SO/96
71 z	28 1	0 19	SS.0-	12	SPCT	L0/80/98
11.2	18 /	\$ 69	15.0-	84	8251	70/80/88
21.2	ī 8 1	1'99	22.Q-	14	1352	LO/90/98
602	64.1	1.69	£\$.0-	L9	9021	70\80\88
902	u.i	62.2	15.0-	ss	Z#01	LO/80/98
902 90 Z	12.1	CZ9	28.0-	es .	6201	LO/80/98
90 Z	5/ T	C29	02.0	27		L0/80/98
5.04	*L 1	c 19	05.0	341		90/90/98
66'1	0/ 1	16s	L+ 0-	La		90/90/98
•	•••					,
06!	59!	€.72	E+ 0-	LOZ	1034	90/90/98
76'1	W I	6 4\$	PP 0-	*0 Z	\$101	90/80/98
68 T	Z9 I	0 45	24.0-	102	LS6	90/80/98
6S I	SC.1	5 L\$	22.0-	Z >	8001	\$0/90/98
8 5 T	≯£ :1	£.74	ZZ 0-	00	95* I	50/80/98
72.1	1.33	6 9+	12.0	B.C.		\$0/80/98
9s1	76.1	9 91	02.0	96		\$0/90/98
\$\$ I	76 1	\$ 99	02.0-	cc	90+1	\$0/80/98 \$0/90/98
\$\$ T ! 22	75 T 25 T	9 99 9 99	0Z 0- 0Z 0-	1E 62		\$0/20/92
33 :	CL i	7 77	0.0	oc.	yet.	50/50/91
SS T	Z€ I	\$ 90	0Z 0	Lζ	1350	\$0/90/98
9\$ T	EE'T	L 90	0Z 0-	٠Z	6521	\$0/90/98
\$\$ t	1.32	??9	0Z 0-	12	1543	50/90/98
9\$1	1.32	999	0Z 0-	61	1336	\$0/80/98
\$\$`T	\$C.1	Þ 9Þ	02 O-	<u> </u>	1313	\$0/90/98
8 5 T	≯C'T	c ()	ZZ:0-	91	LOZ	\$0/80/98
LO Z	9/.1	1 19	15.0-	9s1	6001	10/80/98
2.08	91'1	029	15'0-	124	ES6	10/80/98
2.06	9L'I	E 19	15'0-	751	626	10/90/98
01 -z	% L'1	929	25.0-	671	920	10/80/98
9CL O	129.0	51.9	29.0	9+€	90+1	E0/90/98
287.0	0990	73.2	55.0	344	ESET	£0/90/98
2.03	ደረግ ሂረግ	0- [9	69 0	340	1356	E0/80/98 E0/80/98
SI z	18 1	7.E3	\$9 0- 95 0-	SEE	EEZT	£0/90/98
>1 Z	18 1,	LLY	73 U	UUU	FECT	CO1 501 78
3° I*	12.1	L'59	95 10-	Occ	2121	E0/90/98
∠16 o	SLL 0	Z.T.Z	\$5.0	6ZC	1171	EO/80/98
9Z z	16'1	£.78	09:0-	BIE	1023	Co/win
6ZZ	76 1	9'L9	19.0	316	+601	EO/80/98
2.31	561	9.89	Z9 0 -	EIE	+101	E0/80/98
				ı	_	
(''ma,''mW	(f-ma,f-moW	(4g)		(-gap)		
(101)	(10.0	L	abuxingaM	Angle	(Ln)	(LL)
Emittance	I Density	H		Orbital Dead?	DOM.T	Date.
Spectral	 	Observe		14440	ł	
	P	-,useqD		l	1	1

Table 4. N (10 µm) Photometry of Io (11) (continued).

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5 Z 4	9s!	o SS	8C.0-	81€	1334	Z0/L0/98
5 2 2	LS 1	285	\$E.O-	∆ 1 €	P991	20/70/28
97 z	85 1	555	6E 0-	SIE	1455	20/70/28
5.29	091	2.98	T# 0-	60€	1334	20/L0/98
0s:	091	Þ 95	19:0-	LOS	1351	20/10/98
72.27	651	8 68	0+ 0-	€0€	1521	ZO/LO/98
If':	19 1	?9s	100	Zoc	8921	20/70/38
5 3 3	LS T	I ÇÇ	8E.O-	355	L161	\$2/90/98
<u>2</u> 38	65 1	8 22	O+ O-	318	1252	\$7/90/98
L♥ Z	W I	8 7 č	****	91€	וגגגו	\$2/90/98
6? Z	99!	Sec	SP 0-	CIE	1171	52/90/98
99 1	PE I	€ L?	22.0	28	850	21/60/58
69 1	LE 1	1 8>	12.0-	99	808	21/60/58
WI	cc i	8.64	12.0-	64	617	21/60/58
+9 °1	cc !	899	12.0-	и	cot.	21/60/58
06!	95 T	z PS	7£.0-	L#	1101	82/01/1
и!	75.1	• ES	SE.0-	78		
79 I	66.1	8.85	\$2.0-	87	656	21/70/28
69.1	6E 1	6.81	52.0-	u	616	21/20/58
c> <i>I</i>	1 33	o EP	11 0-	OL	006	21/20/58
LZ.2	C61	0.89	19'0-	784	9051	11/20/58
● Z z	lb f	1 49	09 0-	282		11 //0/58
82 z	P6 I	2.89	29 0 -	182	9251	11/20/58
17 <u>z</u>	88.1	€.66	8 5 0	9LZ		n/LO/W
52 Z	96"!	59?	62. ū -	275	5791	11/20/59
2.20	w ı	0 99	85.0-	L92	9+61	11/L0/\$8
c1 z	18.1	LC9	PS 0-	192		11/L0/58
ii z	6/ 1	2.69	£2.0-	258	1545	11/10/58
90 Z	9L I	8.13	15.0	255	1333	11/20/58
10 Z	● LT	c 19	08.0	222	1303	11/20/58
\$61	991	c 85	\$1.0-	142	2001	11/20/58
16 T	C9 I	z LS	Z+'0-	6CZ		11/10/58
16'1	C9 1	• 12	6+ 0	962		11/20/58
28.1	82.1	s \$\$	6E.O-	232	7+6	11/20/58
Z S 1	\$\$.1	9 0 S	7£.0-	922	728	11/20/58
65.1	1.35	9 L>	22 O-	w	1451	01/20/58
L5"1	+€ I	1 L\$	120	w	60+1	01/20/58
ΚI	SE'1	♦ L?	2Z 0-	w	1324	01/20/58
91 1	1.24	8 c?	E1 0-	54	LC11	01/10/51
1 • 1	1.20	2.24	60.0-	cc	*101	01/20/\$8
9917	22.1	1 69	71 0	53		01/ 40/51
77 1	021	c ZÞ	01.0-	77	028	01/20/51
(8	ES I	0 %	90'0	200		60/L0/\$\$
6 8 I	09!	z %	1+0-	\$61	1115	60/20/51
SB 1	LS T	5 55	6E.O-	861	1111	60/L0/51
	Мст³µт³)	(41)		(¹ 2 3p)	1	
(†mus'mW	(10. _{PP}	(71)	Magnitude	AgaA	(11)	(10)
				36864		(TU)
101)	х Оспаку	MJ.			DOM:	2)8/7
Spectral Smittance 1017	х Оспаку			IssichO	>miT	Date

Table 4. N (10 µm) Photometry of Io (11) (continued).

ì						
ts !	£\$:1	L'ES	0?0-	L* 1	S601	91/21/68
£ 8 1	SS 1	1.15	19.0	5+1	1028	91/21/68
0/, 1	EP 1	20.4	EE 0-	LET	OC6	91/71/68
L? (!≯ <i>t</i>	P 6P	10.0-	PE1	£16 9	86/15/ 19
65 1	≯€ /	E. Th	92.0-	m	6CL	91/21/68
w ı	6E I	L'89	62:0-	211	817	91/21/68
08.1	75 1	S.EZ	00-0-	\$11	959	91/21/68
7 SO	10 I	9.86	50.0	**E	1230	\$1/21/68
01 z	9L !	0.29	95.0-	345	LIST	\$1/21/68
3 33	L8. I	9 59	Z9 O-	340	1200	\$1/21/68
52.52	6 8 °T	£.99	£9.0-	466	£9+1	\$1/21/68
5 26	06'1	6'99	19 '0-	333	9111	\$1/21/68
727	161	Z:19	19 0-	1 cc	00+1	\$1/21/68
Z 3 Z	S6 /	9.89	∠9 ′0−	331	1521	\$1/21/68
LE Z	661	1.07	69'0-	61E	1232	\$1/21/68
7L 1	S# 1	6.08	PE 0-	901		86\ 15\I¢
0L T	Ep. 1	£.08	cc V	601	3011	\$1/21/6 8
17.1	W I	2.05	66.0	36	1034	\$1/21/68
991	65 /	6.84	06.0-	96	L101	\$1/Z1/68
w !	8 E /	2.84	62.0-	C6	∠\$6 €	ILZI168
84.1	₹ 7	9 (>	LI v	22	176	\$1/Z1/68
06 ⊤	41 T	€ 19	71 0	28	006	*1/21/68
LS T	135	£ 9+	≯ Z 0-	(8		1/21 /68
+5 T	62 1	+ 5+	22 O	S 1	9+ <u>/</u>	P1/21/68
151	LZ 1	• ••	00.0	* L	10/	● 1/zl/6t
(51	82.1	7.50	13.0	٤L	€:1	● IIZI16,
£\$ 1	42 !	(60	27.0	49	(0)	● IfzI/6t 11/21 /61
98 1	95 1	1 %	9(b	58	1291	21/21/61
64.1	05 1	32.8	• 0	(8		
12 1	60-1	£ 25	CC v	64	1320	21/21/61
94 t	SP 1	1 15	0C 0-	u	9661	21/21/69
L9 1	0+1	2 69	92 v	u		21/21/69
L9 1	66.1	1 69	92 v	17		21/21/61
99 1	6E T	6.81	62.0	69		21 /21/69
#\$ T	62.1	E.24	71.0-	68	BCIT	Z1/Z1/6 1
0s' I	1.26	€ **	51.0	LS	7111	21/21/69
L+ I	61.1 61.1	£.Eb	21.0-	0s	9701	21/21/61
EP /		0.54	600-	27		21/21/69
8E T	\$1.1	9 09	SO 0-	7		89/12/12
9E'1	91 11	1.05	10.0-	66	016	21/21/61
\$£:1	£1'1	9 6€	coo-	38	LSt	
SE 1	£1 1	6 60	€0.0-	9s	548	21/21/61
5 36	SL'I	19	OS O-	01 c	9101	81/10/68
2.31	891	1.65	9+ 0-	310	1044	21/10/61
7€ 2	IL-I	0.09	27°0-	306	1601	11/10/61
	Mcm.3km.)	(41)		(- 3 5p)	ĺ	
101) (1:ma, [©] mW	₉₁ .01)	(~[/	Magniude	aiga∧	un)	(Ln)
Emmance	и Бевыку	MH.		laidiO əsaff	>m.T	Date
Special	P	Observe				
					_	

Table 4. N (10 µm) Photometry of 10 (11) (continued).

cc z	OL I	66S	Lt v	,90€	9101	81/10/69
26 T	Z9 1	69S	ZÞ 0-	VZOE	286	81/10/69
6 8 (65 T	6 SS	0+ 0-	662	256	81/10/68
E6 I	Z9 I	1 45	74 v	862		81/10/61
<i>II</i> .1	8+ i	225	££ 0-	262	198	81/10/61
64.1	05"1	9 75	Cc 0	288	118	\$1/10/61
11.2	+5 t	242	LE O-	286	6SL	\$1/10/61
OL 1	• S1	06s	9s 0	64	0611	91/11/81
W I	15°1 75°1	z ES S ES	\$€ v \$€ o-	6 <u>/</u> 0L	1156	91/11/81
•••	٠,٠	3 63	300	OL.		7171176
891	75 7	S ES	SE O-	OL		91/[[/8
\$91	601	s ZS	EE v	89		91/11 /81
99.1	05.1	9 75	cc v	89		91/1 1/81
99:1	0s1	t 25	160	w		91/1 1/81
89.1	25.1	9 68	SE V	W	556	91/11/81
£9°1	27 1	615	2£ 0-	SS		91/t 1/81
WI	87 I	225	2E O-	ss	6E8	91/11/81
19 f	9+1	c !S	!E v	0s	708	91/11/81
79'1	9+1 9E1	s I S	1E 0-	0s 9+	208	91/11/81 92/11/29
\$6!	yn. i	ot6	w	47	\$501	76/1 1/ L 1
08.1	57.1	1.44	910 -	96		92/1 1/29
8/L 1	134	L EP	El v	74		7/1 1 //1
* L I	1.22	t 2 9	11.0	22	808	92/11/ <i>L</i> 1 92/1 1/ <i>L</i> 1
W /	21 I 21 I	C 6E	ΖΟ <i>ν</i> ΖΟ Ο ⁺	02 61	154	92/1 1/L1 92/11/L1
ז יכש	-					
ZE 2	1 63	£.72	EP 0	Soz		S2/1 1/L
ic z	29'1	1 15	Z+ O-	[02		\$7/11/2
EE <i>z</i> ZS <i>z</i>	W I	9 L\$ z Z 9	€# v Z S O-	661 991		52/1 1/L1 52/1 1/L1
2,43	16.1	Z 29	25 O-	991 991		57/11/22
[• z	٥٢. [L 6\$	L♥ V	791	010	57/1 1/41
\$5.5	89 1	685	9+0-	L51	965	57/11/2
	\$9°!	Z 85	99°O-	251		2/1 1 /4
5 32 5 46	8/ T	s 79	75 O-	651		57/11/13
6C'Z	ا کُو	8 6¢	L+ V	>\$ 1	6101	EZ/11/23
15.2	₱9°T	6 4 S	99°0-	● C t	058 (Z/1 1 /L1
Ocz	ώì	LLS	E+ 0-	135	918	EZ/11/L
₱€ z	E9 1	\$ 45	c?&	LZ		EZ/T T/L
162	£9 I	272	ZP 0-	150		17/1 1/53
2.28	79 †	I LS	Z+:0-	911	1051	EZ/11/L1
5'46	2.25	06L	U.0-	60s		12/60/L1
15.5	62"2	9.08	08.0-	LOL		82/60/11
2.46	2,25	Z 6L	8/L (O	106		8Z/60/L
LY Z	57.22	£ 67	\$7.0-	300		87/60/21
! ∲ Z	2.20	S.TT	SL V	767	1530	\$Z/60/LS
	(f.m.x ^c .m.)W	((-)		(·\$ap)		
	er.01)	(At)	Magnitude	AgaA		
(1.m4.mW						
,01)				Staff	(TU)	(TU)
	х Бензиу	MH		laischO Saad?	AmiT (TII)	Date (TU)

Table 4. N (10 µm) Photometry of lo (II) (continued).

Table 4. N (10 µm) Photometry of Io (11) (continued).

6L'T 8L'T 5C'Z 7+ Z	05°1 69°1 99°1 981 60°2	17.58 0.58 0.58 9.59 2.17	00 0 20 0 20 0 99 0	491 WI 1991 615 415	LOL 01/10/06 LS9 01/10/06 \$19 01/10/06 9001 60/10/06 \$66 60/10/06
99 Z 99 Z 15 Z 65 Z	10 Z 70 Z 10 Z 11 Z	6 1L 5 ZL 0 PL + 9L	27 Q- 27 Q- 26 Q- 74 Q- 74 Q-	600 600 600	276 60/10/06 116 60/10/06 958 60/10/06 518 60/10/06 cm 60/10/06
19 T 19 T 89 T 58 T	12 Z 92 Z 92 Z 951 15 T	0.E2 0.E2 7.22	9L V LL 0- LL 0- BC 0- Cc 0'	\$26 926 926 926 291 591	9ZL 60/ 10/06 5ZL 60/10/06 1TL 60/10/06 1PZ1 80/10/06 9ZZI 80/10/06 01ZI W/lo/w
LL'1	15.1	1.68	4E.0-	651 451	2911 80/10/00 9s11 80/10/00
64.1				, , , ,	De 11 BOZION
	6+1 6+1 8+1 8+1 891 6+1 8+1	975 \$75 175 075 175 575 075	26.0- 26.0- 26.0- 26.0- 26.0- 26.0-	951 651 151 891 991 691	7211 80/10/00 201 1 80/10/00 0501 20/10/00 8201 80/10/00 5101 80/10/00 256 80/10/00

9L 1	89°t	075	ZE 0-	1191	6C6	80/10/06
14.1	PP11	905	62 0-	LEI	216	80/10/06
OL I	E+ 1	● 05	67 v	961	658	20/10/06
77.1	57'1	60S	0€ 0-	≠£1	948 84	D/ I 0/06
691	77 1	1.05	8Z O-	821	108	80/10/06
89!	ZÞ: t	86?	72 O-	92 t	05L 8	30/ 10/06
69 T	79 1	8 67	L: O	151	964	20/10/06
99 1	6£ 1	160	97 v	511	● C9	80/ I 0/06
W I	01/1	161	92 0-	£11	179	80/10/06
L9 1	0+1	€ 6₽	920-	111	809	80/10/06
80 Z	2.05	זינג	89 0-	ccc	ESE 1	∠0/10/0€
546	60 Z	L CL	04.0-	335		∠0/10/0€
75 Ž	5.13	056	27 O-	of c		20/10/0€
SS z	51.2	954	CL&	328		20/10/0€
5 2	2.18	L 9L	PL 0-	338		LOI 10/06
19 7	22.1	287	9L O	333	157.1	LO/ 10/06
89 2	5.25	2 62	8/ V	∠!€		LO/10/06
OL"z	72.2	L6L	6L o	SIE		LO/10/06
69 Z	92.2	L6L	87.0	116		40/10/0€
01 Z	72.2	8 67	6Z 0-	60E		0/10/06
272	5.29	t08	6L O	80€	nent -	.o/ to/06
07.5	72.2	008	64.0	L6Z		O110KU
SLZ	2.31	218	08.0-	562	616	20/10/0€
•LZ	2.30	0.18	08.0-	562		OI 10/06
59 z	5.23	E 84	11.0	585		0/ 10/06
/* z	2.03	CIL	99'0-	8s:	1 (0	L0/10/0t
7 2 Z9 z	5.20	SIL	SL V	982	418	LO/10/06
19 z	02.2	žü	SL &	787		10/01
95.2	\$1°Z	L'SL	EL V	082		0/ 10IW
ES Z	5113	8 € L	27.0	8LZ		LO/10/01
15°Z	11.2	C \$L	14.0-	9/2	co.	20/ 10/04
79 z	5 Oc	L IL	L9 V	OLZ		LO/10/0x
6C Z	to.Z	8 07	99 0	w:		LO/ 10/00
6€ Z	10.5	90L	\$9°0-	99Z		L0/10/00
17.1	99'T	905	62.0-	171		90/ 10/04
14.7	99 11	1'0s	6Z '0-	611	● C7	90/10/04
09'1	•€1	CLP	22 v	511		90/10/04
69 I	LE'T	● 8?	PZ v	£11		90/10/00
	∠€.1 6Ε.1	06	92.0			90/10/00
\$9 <i>l</i> 99 l	01 1 01 1	Z 6?	92.0	E X	61 8	90/10/06
		8:04	17.0		401	0011045
65.1	66.4	8'99	12.0-	6L	181	90/10/00
85 T	1:33	9.94	Oz v	u	REL	W/10/W
19 f 25'f	SE.1	9 / *	77 O-	**		90/10/04
45 t	26.1 66.1	99*	02.0	и 0 С		90/10/00
L> 1	CE 1	7 70	OC V	VL.	059	70/ TU/UI
	Wcm ² µm ³)	(41)		(-Zop)	i	
(''m4 ² 'mW	₉₁ .01)	*1/	Megnaude	algaA	(TU)	(TU)
eonamim3 (101)	x Densky	uiA		Orbital	Mal Till	Date (Till)
Spectral	-	Observe		14540		
	, ,				1	

Table 4. N (10 µm) Photometry of Io (11) (continued).

(1.mx, mW	(10-m / 10-m	(<i>L</i> _f)	opnazgaliń	(-gap)		
66.1	W I 90'1	c"% Z " LE	>0.0 100	9 8 St	1548	9Z/E0/E6 9Z/E0/E6
59"1	80°F	6.7£	20 0	76	1340	92/20/66
941	1.06	ľ 'K	Z00	L6	1410	92/60/66
99°1	LOT	9 LC	coo	66	5771	92/50/66
29 1	1.10	9.86	00.0	101	6611	92/60/66
09.1	61′1	LIP	80 0-	9£Z	159	LZ/E0/E6
LS 1	91′1	019	90 0-	238	1 9 9	LZ/E0/E6
85.1	41.1	1.14	Lo&	240	50Z	LZ/E0/E6
6 5 i	21'1	5.14	90.0-	Z > Z	026	LZ/E0/E6
191	zz i	6'2>	11:0-	LVZ	ESL	LZ/ED/E6
<u> </u>	1.24	T.E.	ET O	124	112	LZ/ED/E6
<u> </u>	9Z 1	Z.EÞ	c1 0'	952	958	LZ/£0/£6
62.1	6Z 1	c s?	L10-	192	126	LZ/E0/E6
♦८ °1	62'1	● S*	41 0 -	897	1201	<i>LT/</i> CO / E 6
IL ŧ	LZ'I	9 **	\$1.0-	0LZ	SEOI	75/60/69
u	15.1	Z:9¥	61 o -	7.1.7		LZ/CO/E6
64.1	. CC. I	L'9+	02 V	182	+511	L74E0/E6
üΊ	IC.I	1'9+	61 0-	283	1508	LZ/E0/E6
641	cc'!	9:9+	o: 0 -	062	9521	LZ/60/66
84 1	1.32	● 9?	02.0-	C6Z	1321	LZ/E0/E6
8 7. I	ZE 1	€.34	Oz o-	\$62	1332	LZ/E0/E6
94 1	ie i	0.54	61 0 -	66Z	1403	42/60/66
† L'1	1.29	* S*	L1 0-	10€	91+1	LZ/E0/E6
64.1	££ 1	L'9+	02.0	Cos	06+1	LZ/E0/E6
St 1	100	o 8 E	Z00	64	PE9	\$Z/EO/E6
1 28	156.0	S CC	91 0	18	L99	82/50/56
\$\$ I	LO 1	8 7.E	Z00	ES	LS9	82/60/66
1 15	1 02	O LC	SO 0	18	50C	87/60/66
St I	LO 1	8 TE	60.03	28	LEL	82/60/66
171	70 1	r Lt	+V V	vo	372	au turto
8+1 E+1	901	6 7 E	90.0	96 06	694	\$Z/E0/E6 \$Z/E0/E6
8t 1	01 t 011	98C 9 8C	00.0	96 16	0E2	8Z/E0/E6 8Z/E0/E6
61 1	011	6 8 E	000 000	001 04	328	82/60/66
911	601	185	20 0	801	LS6	82/60/66
LPT	601	€.8€	100	0!	6001	87/60/66
0s1	21.1	7.6€	70 0	152	6s11	8Z/ED/E6
91-1	60 1	Cuts	100	LZ1	1315	87/60/66
ít i	60°1	• ts	100	≯ET	1001	82/60/66
91:1	W I	1.85	ZO:0	961	* 151	82/50/56
6* I	11 /	6 ₹€	10.0	6E1	SECI	82/60/66
19 1	0:1	7 24	60.0	283	CC9	63/03/56
691	97 /	7.44	+1 0 ⁻	285	49	63/03/56
0L T	L: I	5.44	510-	782	707	62/60/66
8/ t	££ 1	9.94	02 O-	162	154	63/03/56

Table 4. N (10 µm) Photometry of Io (11) (continued).

OS'T	ES 1 75"! 621	c1 l	6'6E 9'6E	000-	SC:	1500	\$2/60/66 \$2/60/66
0s1	05 1 6* /	11 T 01 f	8.8E	10 0- 00 0	61Z 21Z	1101 856	\$2/60/66 \$2/60/66
Z9	0s1	t i f	1 6E	low	166	6 † L	52/69/66 52/69/56 52/69/66
SST E11 8 6E COO- SC1 P16 80/ZO/Z	SS ! Z9 t	91 T 61	0 O+	90 O-	86 T 96 T	5511 0+11	80/20/26 80/20/26
BS	SS 1 65 t 09 1	ET 1 91 I LT'T	3 6 6 40 6 7 1 7	90°0- 20°0-	338 338 338	716 6551 7551	\$0/20/26 25/05/01 25/05/01
19 t	85 T	\$1:1	9 0+	50 O-	SCc	0681	L0/20/26 L0/20/26 L0/20/26
18 25 99 0: v	19 t 59 1	17 17	£ 19	LO 0' o! 0 -	335 166	+151 1051	10/20/26 10/20/26 10/20/26
\$\$\frac{5\pi}{9\pi}1 \frac{9\pi}{9\pi}1 \frac{1}{2\pi} \frac{6\pi}{6\pi} \frac{6\pi}{6\pi} \frac{6\pi}{6\pi} \frac{6\pi}{6\pi} \frac{6\pi}{6\pi} \frac{1}{6\pi} \frac{1}{6\pi} \q	181	1.32	9.91	0: v	Otc	1536	L0/20/26 L0/20/26
28'1 CET 69"> 12'0" LBZ 156 L0/Z0/Z 8L'1 OE'1 65"> 81 v 28Z 816 L0/Z0/Z Z8 I cc1 69"> 12 v Otz 106 L0/Z0/Z	58 T 98'1	9s1 9£ °1	6 L+	52.0-	20€ 00c	Sell ♦211	L0/20/26 L0/20/26
	28.1 87.1	T 33	6 9 > 6 \$ >	12:0- 11 v	782 782	⊹\$6 \$16	10/20/26 10/20/26
72	17 l				Otz SLZ		10/20/26 \$0/20/26
Date time Phis Density Emittance	(101, Wm ⁻¹ _{pm} -1)	H:01)	(Kf)	ShuingaM	algeA (3ab)	(TU)	(TU)

Table 5. Q (20 µm) Photometry of io (11).

18 £ 59 Z +5£ 68 C- 121 1+5 88 £ 04 Z 095 16 C- 911 80£	
76 E EL'7 19E 26 E- 111 757	
50	
06 E GEZ PIC 92'E 72E 158	
00 50 70 900 FCC 150	LU/SU/LS
539 321 -5.76 314 2.35 3.90	
16 C 0+'T 0350 12'E- 61E LZS	
The ter ite str. The ite	
∯ 62.6 \$1°Z 982 W.6- 011 £+4	12/20/68
\$7.0 \$17. 987. 99°C- 601 \$50	83/01/21
\(\frac{11\cong 0}{2\cong 0} \)	17/10/68
11 E 40 Z 94Z 79 E- 26 9CL	
02 € £1.2 +82 20.6- u 305	
LbZ 861 592 45E- ZB 415	
£0 c 10 1 892 65 €- 14 951	
COC 107 897 65 E- SL EE	17/10/58
# EZC \$1Z 98Z 99'E- 1/2 9Z1	
E+ E 62.5 206 ET.E- OTS TOO	
356 2,44 3.65 M.C. 3.65 M.C.	
79 € 74 Z CZE 6L € 992 € CE 04 ° € 42 Z 268 Z 4 € 597 4 € 6 1 € 64 ° 5 Z 2 € 6UP 657 LM 2 € 65 Z 4 € 6UP 657 LM	83/01/20
01°E	83/01/50
ZL E 60: ZEE ZEE- 09Z SSI	
66 € 65 z SEE OUP 657 LM	
10 524 373 545 5 es	
69 C 94 7 6ZE 18 C C 700	
900 314 3 65 2.44 3 65	83/01/20
79 E 7 T 57 66 E- 967 209	
SO C SO Z ELZ 19 E- 19 SM	61/10/18
50 c 50 Z £LZ 19°E- £9 ZM	6 1/LO/E8
L67 66°1 992 25℃ 6s L 0 €	61/10/58
ZO'E 60 z 1LZ 09 E- LS 100	61/10/68
20 £ 50.5 CTZ 10.E. ER TEI	61 /LO/E8
80°E LOZ 9LZ 79E- 15 8Z	61/LO/E8
6L'E #9'Z LSE 06'E- 9s	\$2/\$0/E8
" €8 > ZL'Z Z9€ 26 €- Z01 015	1 11/60/68
88 > 1/2 99s C6 E 201 80s	1 11/50/58
76+ 89.2 72C 00.E- 000	90/50/58
98 1 0260 EZ! OLZ. OFE	83/02/08
16:1 9?60 9ZI LL'Z- OFC	83/05/18
28 68°T 252 25 E- LEE	81/20/68
98'E 06! PSZ ESE- EEE	81/20/68
	t 1/20/ET
to? 84.7 OLE +6.6- 646	95/03/16
914 68.2 LE 96.6- 246	61 /0/78
91'E 60'S \$7\$ 69 6- 26	21/10/08
(qcf.) (fy) (m ² ,m ⁴) Wm ² ,m ⁴)	
(101) **01) sbusingsaid algae.	un)
ime Phase Flux Density Emitlance	
innered image	_
tovisido (mm)	

Table 4. N (10 µm) Photometry of Io (11) (continued).

Smittance (10° Wm ³ µm ⁴)	Yasmad xi		ShuingsM	iaidiO Starfi SignA	CTU)	sseQ (TU)
	(,.ш _м г.mэW	(Ar)		(-gop)	ſ	
18.1	≯ € I	c ∠ ≯	2Z 0-	263	L+L	62/60/66
981	86.1	5.84	52 0-	10£	6£8	6Z/E0/E6
S8 T	LE 1	• 87	₽ Z <i>V</i>	Zos	ise	6Z/E0/E
28.1	7£.i	€ \$≯	≯ Z 0⁻	60€	626	62/60/6
28.1	86.1	+ 8+	●Zv	116	256	62/50/5
68.1	t ≠ f	\$ 60	LZ 0-	919	1020	62/60/6
981	₽C't	9 81	57 0-	372		62/60/6
ZB 1	35 E	S LP	22.0-	9ZC	19† I	62/60/6
64 1	1.33	8.54	12 v	328		62/E0/E
CL ?	82.1	Ŧ \$Þ	L1 0-	SEE	1 340	67/60/6
17.1	LZ I	9 **	21.0-	9SC	Z\$Z1	62/60/6
L9 I	● z 1	L ED	ET O-	338	SOE!	67/60/6
991	1.23	∌ c?	€10-	94€	61 €	62/60/6
19°T	1 33	6 2 >	11 0 -	345	EEET	62/00/6
Z9 1	12.1	• 7>	01"v	342	1325	6Z/E0/E
L9 '1	1.24	9 EÞ	E1 0-	LPE	1402	62/60/6

Table 5. Q (20 µm) Photometry of Io (11) (continued).

ES C	66Z	66E	£0.4	9s	6211	\$0/80/98
3 45	[6 Z	88£	66 €-	16	PSET	\$0/80/98
LZ C	2.78	OLC	₱6 E-	62	8661	\$0/80/98
9€€	\$8 z	648	79.E-	LZ	1353	\$0/80/98
3.40	58.2	986	66.E-	τz	1546	\$0/80/98
9s. €	98°Z	382	\$6'E-	02	1535	\$0/80/98
02.£	L6 z	L6S	ZO:1-	ĹĪ	\$121	\$0/80/98
91 ●	€§ c	049	4.20	951	1101	10/80/98
22.4	LS C	944	ZZ:\$-	>51	9S6	10/80/98
5€ ●	69 €	169	4.25	ZS 1	Z\$6	10/80/98
£\$.*	+8 €	215	62 ►	601	673	90/90/98
● L I	L¥ I	961	52.E-	99E	9561	£0/ \$ 0/98
07.€	3.12	910	£0°₽-	338	STET	£0/ 9 0/98
18 c	77 c	6Z >	01.1	EEC	1536	E0/80/98
€9 c	40 €	60+	\$0°₽	166	6121	E0/80/98
99′€	LO C	019	50 t-	61€	9s01	E0/80/98
\$+ €	76 'Z	160	00.1►	916	6001	£0/ 80/98
8 ₽ c	96Z	365	10'₽	313	8101	E0/90/98
16.4	10.6	100	€0.1	316	6 >> 1	20/10/98
1€●	10.€	TO+	€0.1≻	715	7641	ZO/LO/98
81 P	76.2	68E	00 ₺~	915	9711	20/10/98
75 ●	\$1.€	124	80.1►	60€	6EE1	ZO/LO/98
9++	lic	\$1+	LO ►	LOS	1356	20/ <i>L</i> 0/98
1919	80.E	110	90 ►	+0C	1305	20/10/98
45 *	20 €	√ û•	ço ►	505	1523	20/20/98
\$6 €	19.2	150	68 C	222	1350	\$2/90/98
96 €	79'7	255	68 C-	• IC	1524	\$7/90/98
6€.≯	£6 Z	160	00 ►	● IC	• ZC	! \$2/90/98
,, ,			•••	***	£121	\$2/90/91
<i>IL</i> ≯ 91?	≯1 c 9€ E	61) 277	90 F	CIC	718	21/60/\$1
16 C	91 c	727	\$1 1 - 20 1-	८8 6L	917	21/60/\$1
02.4	09°E	64	ZZ -	\$2		Z1/L0/S1
10?	EP:E	859	<i>ιι</i> +	78	1043	21/20/51
10:	Lr L	•>/	2.1 *	**	LFU	C1/20/34
10*	54.6	254	∠1 17►	и	1001	Z1/L0/S1
96.€	₱€ c	599	91 P	и	/ S6	21/40/51
22.6	SL.S	L9S	£6 E-	õζ	706	Z1/L0/S1
26.5	₹€	\$ 99	\$1' \$-	CRZ	6681	11/20/51
28.6	9.7.€	450	Z1:≯-	197	1258	11/20/51
1 22.5			•••			
≯ 1●	ES.E	1 L?	92.≯~	922	1424	11//.0/s8
cot	EÞ.E	857	∠1 ′⊁	ELZ	1431	11/20/51
1511	£8.€	91s	0€. ►	L9Z	1348	11/20/51
05.4	cc c	11s	62.1	Z9z	1310	11/20/51
08.5	60∵≯	SPS	9€:⊁-	258	1547	11/20/\$1
9 €●	OL"C	£6#	\$2.4-	722	1223	11/20/\$1
97.●	£9 É	727	£2.4~	727	1302	11/20/51
9Z Þ	€9.€	727	€Z.1-	741	1044	11/20/51
('mai'mW	Wcm² µm')	(Af)		(-gap)		
(10.	_{pt.} 01)	v-1/	Magnitude	olga∧	(TU)	(LL)
SOGETHERS.	Density	XUIT		Phase	этиТ	Date
Specital	<u> </u>		I	Orbital	_	-
	l	Observed				

Table 5. Q (20 µm) Photometry of Io (1) (continued).

90∵≱	LF C	79*	81 🕶	533	1033	11/10/5
901	LV C	29 *	81 🏲	, 167	£101	11/20/5
18 €	3.25	664	! 1 ● -	533	546	11/40/5
IL C		124		526	658	11/40/5
61' C	91 €		80° b- 10 v-	19 L9		01/20/5
	96'7	365				
61 C	5.96	\$6E	10 🗲	ç9	9561	01/20/5
cc £	2.83	LLξ	96C	54	6 t h	01/20/5
10 c	95.2	146	28 C-	• C	4101	01/20/5
10 c	9\$ z	1 ● C	28 €-	77	228	01/20/5
86 E	₹£	677	\$1 P-	16!	≯III	60/10/9
LO C	5 2 2	300	IT E.	69	1526	01/90/9
3.04	c: z	765	oγ €.	۷9	1244	01/90/9
LO f	\$2 z	90€	16.5-	99	1534	01/90/9
10 E	5 23	L6Z	0/ E-	63	1511	01/90/9
₩.E	£2.2	L6Z	OL E.	19	1021	01/90/5
EB z	LO"Z	942	79 €-	09	1125	01/90/9
ZB c	8L Z	370	₱6 f-	332	1521	60/90/9
*O:C	co -	324	68 f	*ZZ	£+21	60/90/9
99°€	\$9 z			t z:	9771	60/90/9
89.E	89 Z	72E	06 €		9021	60/90/s
85 C	2,60		£8 c-	518 514		
17.6	0/ Z	09s	[6 (1134	60/90/9
3 92	98.2	185	79 E-	212	cett	60/90,'s
35 €	>> z	976	0s c-	● C	1450	80/90/9
07.E	S9z	LSE	06 €-	0s	1322	80/90/5
91 c	62:	302	U c-	₽ Z	EIEI	80/90/9
18 Z	2.03	172	09 €-	20	1543	80/90/9
91 🖈	5'66	66€	₹0 ⊱	£91	0111	LO/W/S
69:♦	TE.E	699	\$1 % -	791	8501	L0/90/5
ĬĹ C	S9z	324	68 E-	333	1542	90/90/9
₹8.€	SL'T	L%	£6 E-	335	LEZI	90/90/
00.E	St z	982	99°E-	725	4511	90/90/9
₹8 z	10 z	S9z	65 E-	372	911	90/90/
9E.E	0≯°Z	330	8/ °€-	oz	6101	02/80/1
ye e z? ●	91 c	1 z?	80.4-	LPE	5+5	02/80/1
SO ?	16 Z	88E	66 E-	LPE 491	268	6 t /80/1
\$0 ? 10 ●	88.Z	*86	86 €-	991	578	61/80/1
10.4	88 Z	786	16 E-	\$91	918	61/80/1
		700	CO &	961	***	01.00:
4.20	70 c	405	€0 >-	851	732	61/80/1
₹6'€	2.83	LLE	96 E	123	LS9	61 /80/1
10*	LZ"C	964	71 •	>EE	028	61/20/1
LO*	16.6	299	CI 🏲	335	208	61 /L 0 /1
81.4	3.40	757	9! 🐦	LZC	6ZL	61 /LO/1
\$6.4	SS.€	CL?	17+	9ZC	817	61/10/1
4.20	₹ ¢	950	LT 🏲	353	959	61 /LO/1
09.E	7.51	332	E8 C-	0£1	£##1	80/50/1
TZ.E	2.49	ZEE	28 E-	128	1431	80/50/1
	("ma_moW		1			
("ma,"mW		(A)	Magnikude	(gop)]
	"o I)		Magnitude	Angle	(TU)	(Ln)
,.01)						
Spectral Spectral	Density	XU1-1		latichO Starff	omi⊺.	Date

٤٤ +	96'€	625	€€.⊁	912	LOL	LO/10/06
19 5	78 E	415	0E. > -	897	119	L0/10/06
15 *	64 E	sos	8Z Þ-	599	95s	LO/10/06
⊊ī e	6 † C	991	61.*	611	9621	90/10/06
SR c	3.25	£E>	tt'⊭-	£11	,	90/10/06
J., J	,,,,			•		707.0700
so t	10 €	+5+	LUP	83	! u	90/10/06
6L E	61 i	92>	60 1	SC	201	90/10/06
£L c	† € C	217	90'#~	и	LOL	90/10/06
LL C	81 C	*2*	60.4-	IL	CS9	90/10/06
98.€	3 25	>E >	21.1	9∳1	● co!	91/21/68
16 £	Occ	664	€1:19	SET	816	89/13/16
41 E	L9 Z	LSE	06 €-	611	627	91/21/6S
17.5	c1 €	L19	LO >-	9! 1	207	91/71/68
£6 C	16.6	199	£1'9-	1≠€	90s1	S1/21 /69
56 E	CC C	643	9T 9-	338	277	\$1/21/68
LO 🕈	3 43	LSP	LI 🗲	335	SO+1	\$1/21/68
Ltt	9L.C	105	TS. 4-	350	1340	\$1/21/68
\$6 P	9! 🕨	+55	\$C.7-	101	\$111	>1/71/68
18 +	10 P	6ES	SE.A	L6	1023	\$1/21/68
64.●	20 ●	966	#C.#-	56	9001	● <i>Inl/6</i> *
0s 🖈	87 €	¥05	82.4-	9s	sob	● IIz 1168
56 €	u f	799	91'9-	18	618	¥1/Z1/68
16 C	67 C	438	E! >-	и	6 U	\$1/21/68
16 €	6: £	167	€1:1=	69	907	₱1/Z1/68
3.83	0.0	LZ*	01 🗲	?8	27¢!	21/21/68
200	000	20,	0.7	.0	acr.	0.,0.,00
SL E	₱f c	217	20 📂	87	1961	21/21/68
69 €	60 E	415	90 >-	ZL	1529	21/21/68
78 C	>7 €	431	†1 > -	04	1544	21/21/68
95 c	L6 z	96€	20 🗲	45	4111	21/21/68
cc 🕻	6LZ	27.6	26.E-	69	1101	21/21/68
0€.€	9L Z	89 0	+6°€-	L*		21/21/68
3 00	1s'2	>EC	€8 €-	10	126	21/21/68
68 Z	1 ≠ z	333	6L E-	LΣ	?51	21/21/68
16.2	L S C	915	0€ №	60E	4601	81/10/68
82 S	58 c	+15	0E.≯-	40£	1023	\$1/10/68
2 1 s	87.E	105	87.5-	WC	1201	\$1/10/68
+9 €	90.E	204	20.₽-	Zoc	L1-6	\$1/10/68
28.E	EZ E	164	101	300	LE6	11/10/68
€6 C	OE.E	0++	€1:1	862	† 76	\$1/10/6\$
+0 *	9E.E	425	911	767	968	\$1/10/68
05.4	3.28	46>	ZI: }-	L \$Z	108	21/10/68
28 E	05c	494	4.20	08	cc!	91/11/88
≯L'E	BC.E	059	91:1	u	P111	91/11/88
06 E	SS.E	699	0Z:1-	69	0201	91/11/88
[0 +	Z9c	CRY	57.1	69	2101	91/11/88
99 E	IE.C	199	€1:⊁-	L9	5001	91/11/88
,,,		.,,	,	27	5501	
I md	Wcm²km²)	(41)		(- 2 5p)	Γ	Ι
('ma,'mW	,,OI)	(~[)	ppnyuZejų	olynA (Ase)	100	(10)
(10.1		vini é	'		(32)	(TU)
sonething	Densky	T.116	1	- Stadfi	SITING.	>MQ
Spectrai		ремонфо		ianich(O	l	
						L

€9 €	82 c	866	21 ≯ −	۲9	1001 91/11/81
₹8.€	8t E	191	611	. 59	6*6 91/11/81
O+.E	80.E	019	50°F-	LS	618 91/11/81
TE.E	50 c	901	Þ0 ·Þ-	, 98	81/11/81
₽ ₱°E	01 c	219	90'1-	05	008 91/11/81
99°E	01 E	*!*	90 +-	0s	6SL 91/11/81
18′€	99 Z	322	06.E-	w	2011 92/11/21
88 €	IL Z	195	Z6′€-	9€	0S6 9Z/11/L1
86 €	8L Z	ILE	≯6 €·	≯€	EE6 92/11/L1
101	78 z	9LC	96 E-	53	£12 9Z/11/L
. 16 €	5.11	69£	₱6'E-	50	854 9Z/11/41
20.8	ES C	149	02.4-	700	● ill \$2/11/L1
80 S	LS E	LL*	7 7.75	707	0 SO 1 SZ/11/L1
81'S	19°E	981	₽ ₹'\$~	200	11/1 1/52 1040
90.9	9: 🕶	895	IÞ 🏲	WI	\$\$9 \$2/11/L1
99 S	86 C	0s5	££.+-	w,	069 \$2/11/61
99 S	€8.€	015	62'1-	791	● 19 \$Z/\ 1/L!
10'5	εγ.ε εε ε	867	97 1-	L\$1	005 \$2/11/21
21.2	09 E	084	£5. }-		80\$ SZ/11/L1
21,5	09 €	Uav	16 7	ESI	BUS SCITTILE
Z8.2	\$1 b	ESS	86.4-	651	LZ11 EZ/11/LE
L#'S	68 c	615	16 🏲	†\$1	ESO1 EZ/1 1/LS
£2.23	3.72	L6*	97.1	501	PP6 EZ/11/LS
86.2	€8.€	115	6Z Þ-	561	11/11/23 832
04.2	3 85	ETS	9€ ►	cc!	17/11/23 820
5.23	CLc	L6#	92.4-	g _Z 1	9*L 62/11/4
44.5	3.90	250	1€ ≯−	121	L59 E7/1 I/if
20 +	L9 C		SZ +	60E	6051 82/60/15
60 >	7L E	86*	LZ >-	LOC	6S?1 \$7/60/L
7.42	90°9	662	5€.	10€	E 1+1 87/60/LE
ST'+	6L.E	906	82 🗠	262	2161 82/60/49
ES +	110	252	\$C'#-	L8Z	12/00/28
77.	99°€	589	*Z *-	6S2	8901 80/80/91
86 €	24.E	551	LT'D-	88	PEP! 40/80/91
\$6 €	6C C	254	91.4-	98	02? I LO/80/91
00 +	₩.€	659	2119-	18	8+C1 L0/90/91
#0°#	LP C	E9+	61.5-		
76.E	t≠c	+S+	91.)-	6L	1EE1 L0/80/91
ος <u>ε</u>	71.6		60°F	L9	6021 (0/80/9)
02 E	60C	453 415	90 ⊁-	55 23	110 1 L0/80/98 1 CO / L0/80/98
09 6	800	CIP	90 F	13	1 CO 1 20/80/98
89.E	91 c	12+	80 %-	61	6s6 <u>/0/80/98</u>
£7°7	87. E	505	82.2	741	ow/90/80/98
* 2.*	€9.€	683	€2 ►	722	0111 w/80/98
28.€	6ZC	6C?	E1 1-	LOZ	9601 90/80/98
Z8.E	UС	964	21.0~	504	\$101 W/so/%
09.€	wc	119	90'1∽	Zoz	1001 90/80/98
16.€	LO C	609	\$0.₽-	BE.	ETT Co/so/w
	T ind man		· · · · · · · · · · · · · · · · · · ·		,
(*'m4*@W	('ma'maw	(14)	1	(2 3p)	1
(10-1	w.O1)		Magnitude	οίξα∧	(un) (un)
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Table 5. Q (20 µm) Photometry of lo (11) (continued)

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ŧ	10.4	847	100	*n.		-01	17,00,00
ı	10 + 01 +	86 Z	L6C	20 p-	727	206	LZ/80/86
ı		₩.	909	-	522	698	72/60/66
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ł	\$E.Þ	3.23	430	117	243	SZL	75/60/66
1	ZÞ Þ	3 2 g	154	Z1 🏲	341	OIL	LZ/E0/E6
I	01 🖈	69 C	59+	61 🕶	238	ES9	LZ/60/66
1	105	3.72	567	9Z *	962	6C9	LZ/60/56
8	ILε	SLZ	L9C	£6 E-	100	1691	97/50/56
	8/ C	2.80	ELE	\$6 C	L6	91+1	92/00/06
H	EL C	LLZ	69s	≯6 ′	€6	9961	92/60/66
É	*LC	LL Z	OLc	*6 €-	L 8	1304	92/50/56
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ï	3.82		87.E	∠6 c	SB	1520	93/60/66
	b? C	5'29	345	78 C-	SL	1138	92/60/60
Ü	8t [25.5	94€	98 €-	w	1032	92/60/66
ı	LS C	59'7	323	68 € -	LS	626	92/60/60
Ĭ	191€	0L z	09s	16 E-	0s	198	92/50/56
4	19 c	2.68	LSE	06 c-	27	878	92/60/66
Ā	ĽC	z? z	eze.	ot 🗜	Z*	9 + L	92/50/56
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ŧ	3 3 8	15.5	SEE	EB.E.	6s	614	92/60/66
ı	_						
0	₱1 C	EE Z	116	27 E-	۷E	90L	92/50/56
Ħ	90 €	LT Z	WC	ST E-	32	559	9Z/E0/E6
į.	EL c	9L.2	69s	● 6 f	526	00)	\$2/E0/E6
į	80.€	uv	ÇQÇ	26 €	* \$Z	L191	\$2/60/66
	E8 c	2 83	BLE	L6 t-	248	9661	52/60/66
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Ň	6! >	o! £	+1+	90 ● -	238	1 3 5 e	SZ/EO/EI
	¢ 33	EI.E	21>	LOW	9CZ	0121	SZ/ED/EC
	71 9	20.E	L0+	\$0.4-	022	9101	SZ/CO/E
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	31 7	SU C	9V*	ru r	***	7001	SCHOOL
	68 C	2:88	?*C	86'E-	602	L S8	52/60/61
	69 €	2 88	+8€	\$6 €-	907	0+8	52/60/61
	84 c	2:80	PLE	\$6 €-	300	+ \$L	52/E0/Et
	*L E	<i>LL'</i> Z	0/E	n E-	861	194	52/60/61
	PO P	96°Z	56C	10.1	L 11	9+11	90/20/20
	78 c	7.80	ELE	\$6°E-	921	616	80/20/20
	341	6 > Z	332	€8 €-	t cc	1206	L0/20/21
	∠+ E	7.54 2.54	3 56	\$8 c-	LIC	9051 9ZEI	L0/20/21
	28 C	78 z	SLE	% E-	60€	1228	10/20/21
	3 8 C 10 ●	2.93	160	% t°	10s	6211	L0/20/21
	10 🛡	£0.6	105	W F	108	6611	LOILLIA
	69 €	OL'Z	09s	16.6-	282	9S6	10/20/21
	● 6 C	2.89	385	66 E-	787	906	10/20/21
	96 €	88.Z	786	16 C	714	ESET	50/20/24
	68 C	2.83	LLE	% c-	197	SIEI	\$0/20/21
	SL C	27.2	WC	26 € -	749		50/20/21
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	,01)				oignA.	(TU)	(TU)
	Ematence	Plux Densay		1 1	) निकार	>m₁)	Date
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	1		~	i			

Spectral		Observed		- Orbital		
namanga 101) 'ma'ranW	Plux Densky (10**		Magnitude	Phase Phase Assic (deg.)	miT (TU)	Date (TU)
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SL'+	26 €	62s	€€.⊁-	872	CZL	LO/10/06
89.4	<b>96</b> €	\$25	4.32	<b>784</b>	908	LO:10:06
£2.4	i g c	808	62.1←	286	820	40/10/06
28.4	80 +	C*S	9E+	293	too	LO/ 10/06
96 >	LT'+	9ss	86.4-	562	776	L0/10/06
65 🏲	98.€	<b>*</b> 1\$	9€ ⊁	wc	1023	LO/10/06
ts'+	6LC	90s	<b>2</b> Z <b>←</b>	310		10/10/06
67.7	ĽC	Cos	8Z.1	STE	2411	LO{10I06
07'7	OL C	€6≯	-4 52	353		LO/10/06
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Z+ +	14 C	564	92.4-	330	6261	LO/10/06
L+ +	9L C	105	72.A-	335		LO/ 10/06
ZB.€	3 S I	824	01.4-	115	019	80/10/06
¥8.€	3 23	430	11 ←	<b>*11</b>	<b>*</b> 79	80/10/06
16.€	9.29	664	£1:1 <del>-</del>	≯Z t	6CL	m!! 0/06
\$0.4	O+ E	+5+	91'1-	921	<b>+</b> \$Ł	80/10/06
11.4	3 42	09+	\$1° <del>5</del> -	134	848	CW10KI6
<b>*1 *</b>	84 C	191	61 1	9€1	106	80/10/06
42.4	LS C	94*	72 🕶	111	116	80/ 1 0/06
16.4	<b>29</b> c	683	€2.1≥	<b>L</b> >1		80/10/06
os Þ	87.€	105	82 <del>-</del>	ZS 1	1025	80/ t 0/06
4.12	9+ €	799	81.3-	9s1		80/10/06
64.49	81 c	<b>+0</b> 5	85 4-	851		80/1 0/06
89.4	\$6°E	252	ZE.+-	£91	1212	80/1 0/06
SL'+	66.€	262	<b>≯€</b> ′ <del>⊁</del> −	591	1558	80/10/06
21.2	0s?	*LS	ZP 9-	L91	8 1244	0/ 10/06
<b>*6</b> **	<b>†! †</b>	! \$\$	8E +-	*2E	7!L	60/10/06
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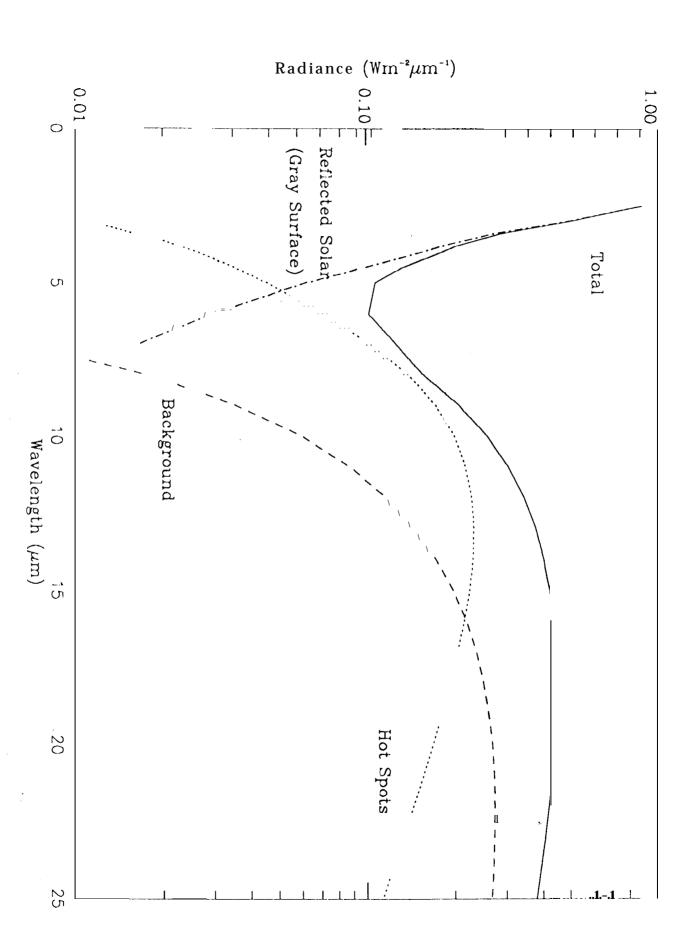
60/1 0/06 60/1 0/06

10L 506

6Z6

Table 5. Q (20 µm) Photometry of Io (II) (continued).

	<del></del>					
EE *	3 51	624	01 -	LPE	01+ 1	62/03/56
01 <b>C</b>	51.2	99s	£6 £-	916		52 /EO/E6
1 L C.	94 Z	L9C	EQ E-	344	2961	67/03/56
45 C	\$9.5	ESE	6S C-	!∳€	1353	63/03/56
						63/03/56
79 c	69Z	956	16 £-	336	1310	
88.E	88 Z	28€	86 €-	755	1521	62/60/66
10.0					A 17.	47,50,54
16 C	7 60	78£	66 €.	332	● *Z1	62/60/66
91 🗗	60 €	415	96*	6ZC	1200	62/60/66
<b>\$1</b> ●	80.€	017	90 Þ	LZC	7411	62/60/66
917	60 E	415	90?-	353	5111	62/20/26
17 ?	cl E	L10	LO+	350	9S0!	62/60/66
						1
16 >	Oz c	LZ#	01 <del>b</del> -	315	<b>LS6</b>	62/60/66
9€ ♦	61 C	974	01 🐦	310	596	62/03/56
87?	81 E	454	60 ►	303	LSB	62/03/56
62.7	61C	\$20	60 tr	106	448	62/60/66
<b>4</b> 23	+1′€	61+	80 <del>p</del> -	<b>≯6</b> Z	ZSL	62/69/66
CC +	., •	0.,	00 7	,,,,	CJL	00,00,00
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€.U	LOC	019	\$0 P-	882	807	62/60/66
£Z.>	<b>≯1.€</b>	619	80 P-	982	259	62/60/66
613						
	TO C	601	s o *	784	869	63/03/56
≱ò.£	OLz	09s	16 €-	0+1	1961	82/60/66
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29 €	£7.5	WC	C6 C	132	9061	82/60/66
26€	2 02	*55	68 €-	92 !	L!Z!	82/60/66
+9′€	OLz	09€	16 €-	<b>9</b> 7 1	1504	82/60/66
95°€	197	285	68 C-	111	1014	87/60/66
69 C	5.59	345	£8 €-	66!	2001	87/60/66
LS C	<b>\$9</b> z	324	68 <b>C</b> -	100	605	82/60/66
15 c	19 Z	348	<b>∠8 €</b> -	L6	9ss	82/60/66
ES c	79 z	6₽€	88 C-	\$6	823	87/60/66
z? c	5 24	Kc	₱ <b>8</b> c-	16	SSL	82/09/58
L+ c	LS Z	343	98 €-	68	247	82/60/66
EP c	2.55	940	28 E-	● 8	OIL	82/60/66
2+°E	<b>8</b> \$ Z	345	₹8 €-	28	589	82/60/66
SC C	Is Z	332	C8 C-	08	019	82/60/66
90'*	70 c					
90 F	LU C	704	c o *	WC	9691	LZ/E0/E6
00°C	00.7	700	06 F.	700	1741	17/50/51
38.E	2.86	382	86 €-	302	1451	63/03/51
98 €	98.2	186	86 c-	oos	80+1	72\co\ce
18 E	28.2	080	L6 €-	967	1340	72\E0\E9
ig c	28 Z	Щ	96 €	<b>767</b>	1327	72/60/69
69 €	2.74	59€	£6 E-	162	1 305	72/60/66
41.6			•-	*07		
67.€	2.8.2	STE	96 c-	*8Z	1213	LZ/E0/E6
9L C	6L Z	372	S6 €	282	1500	72/E0/E9
66 E	96 <b>Z</b>	365	10*	172	1901	LZ/E0/E6
3.94	7.92	390	00 ≯	697	LZO1	LZ/ED/E6
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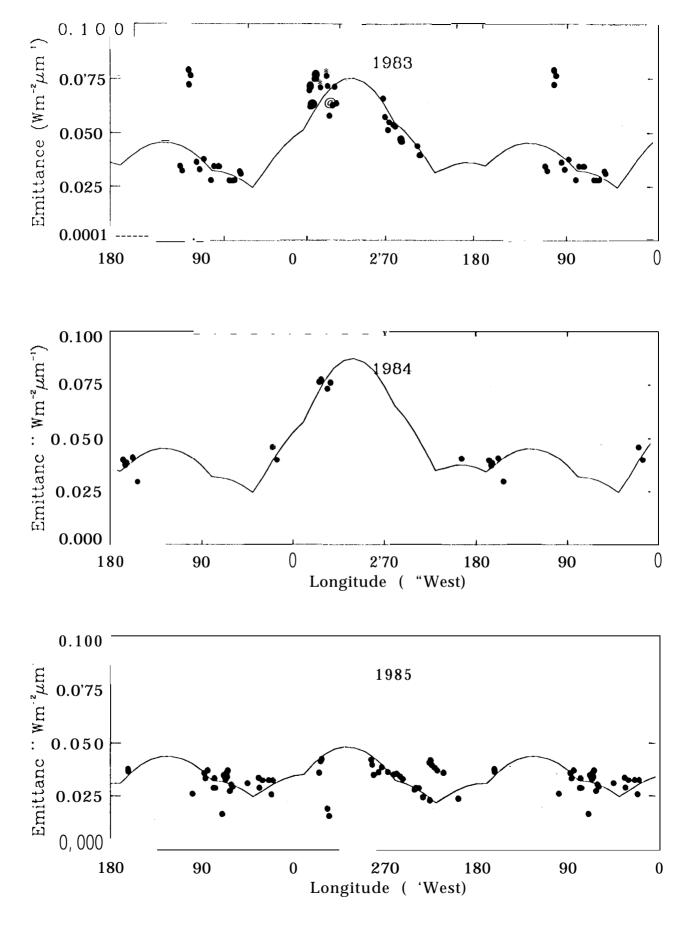
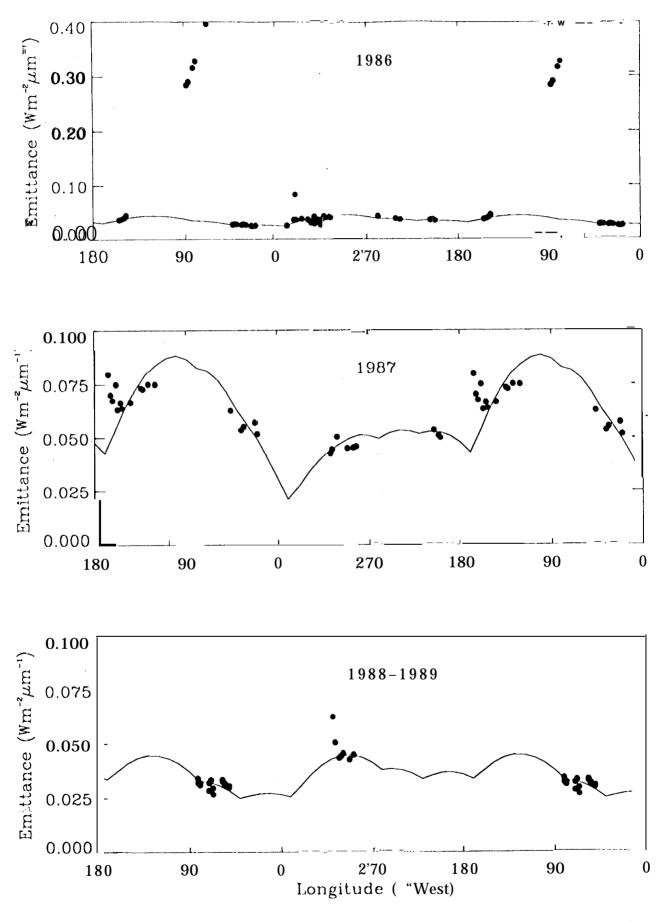


Fig. 2a



F1926

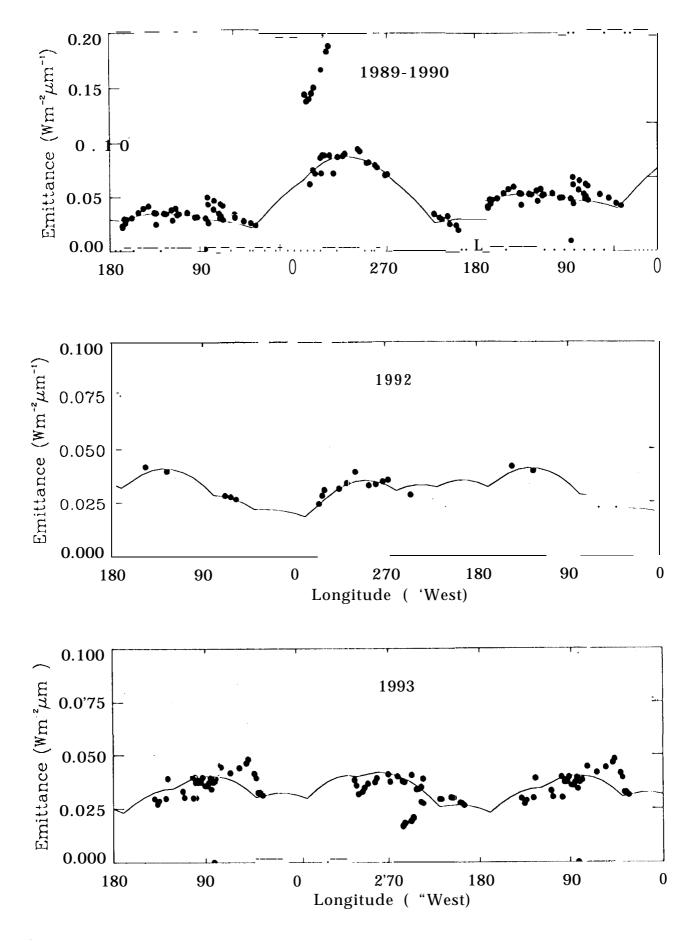


Fig.2e

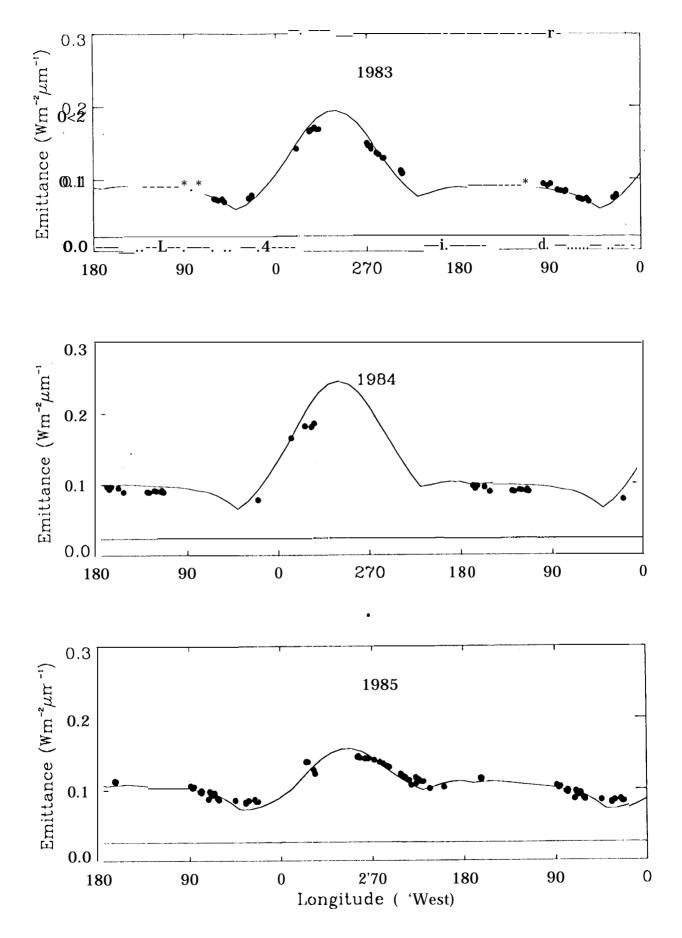
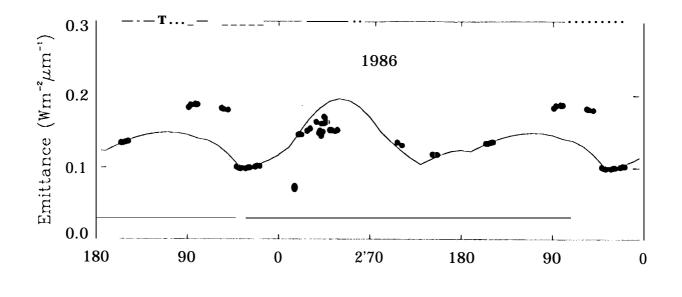
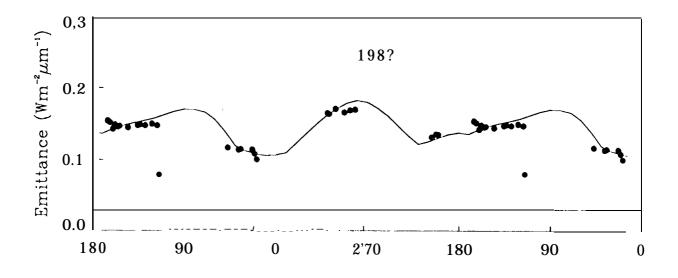


Fig.3a





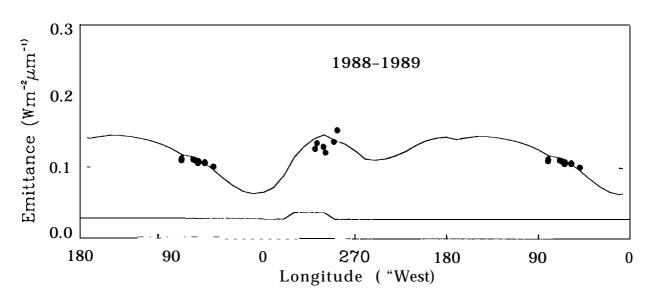
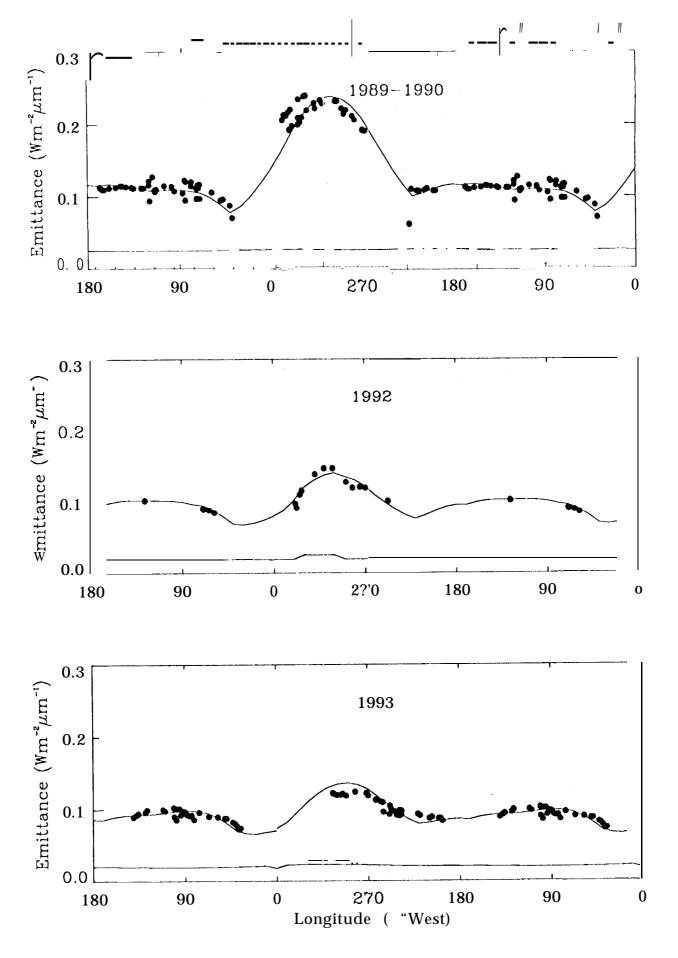
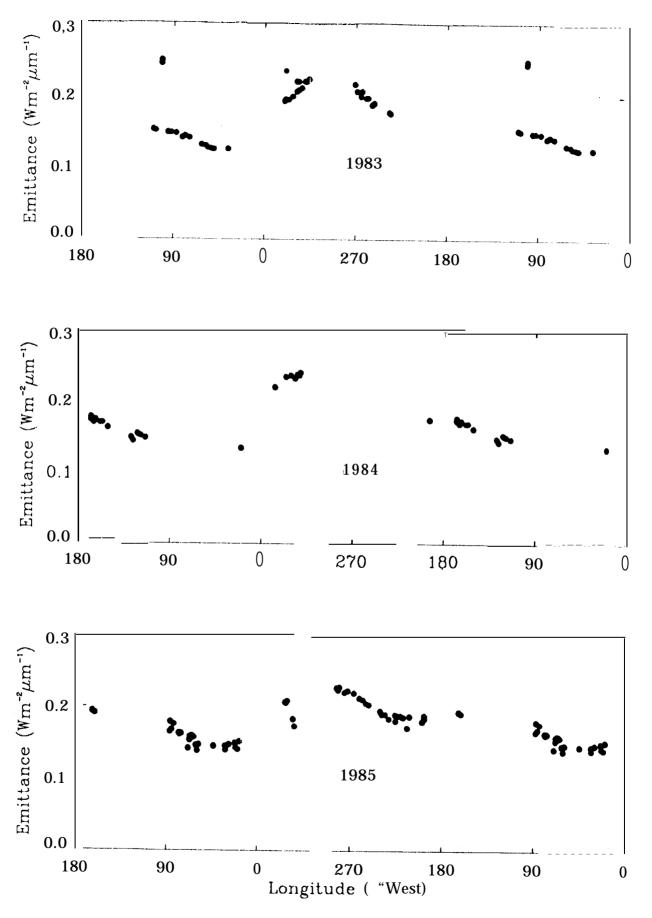


Fig. 3b



Figh



F18.4a

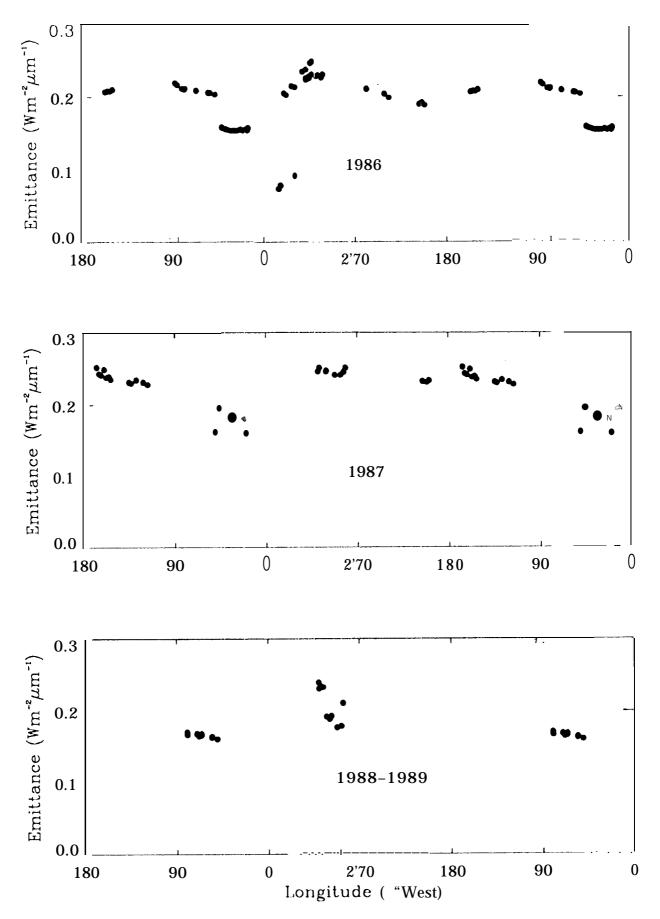


Fig4b

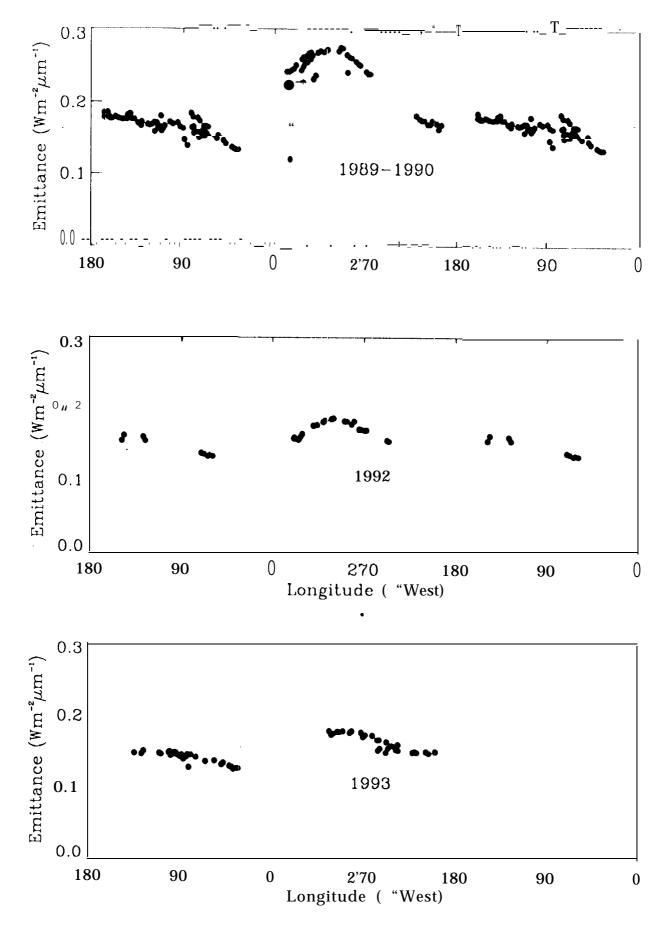
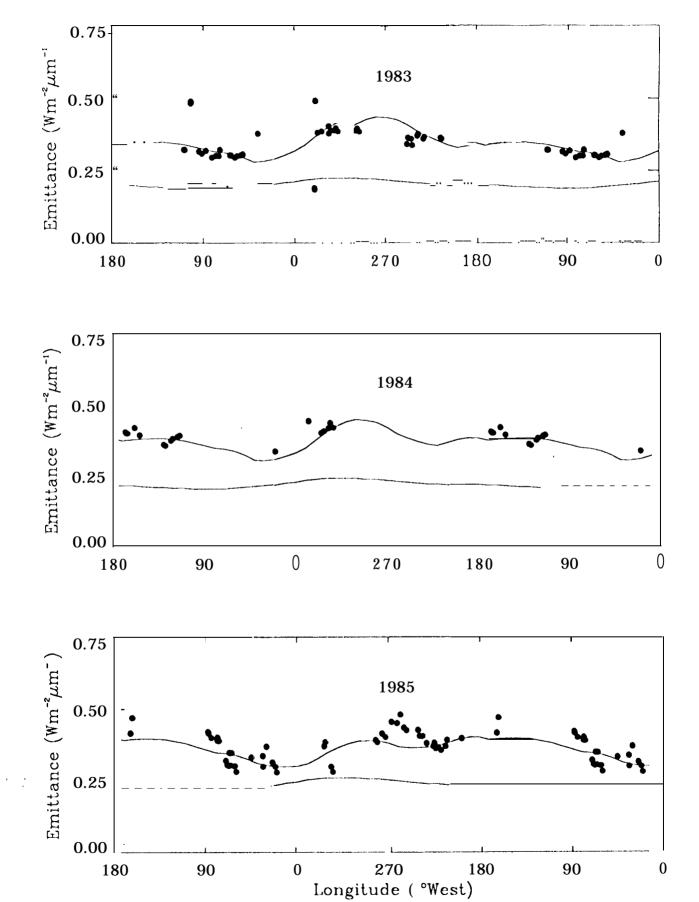


Fig. 40



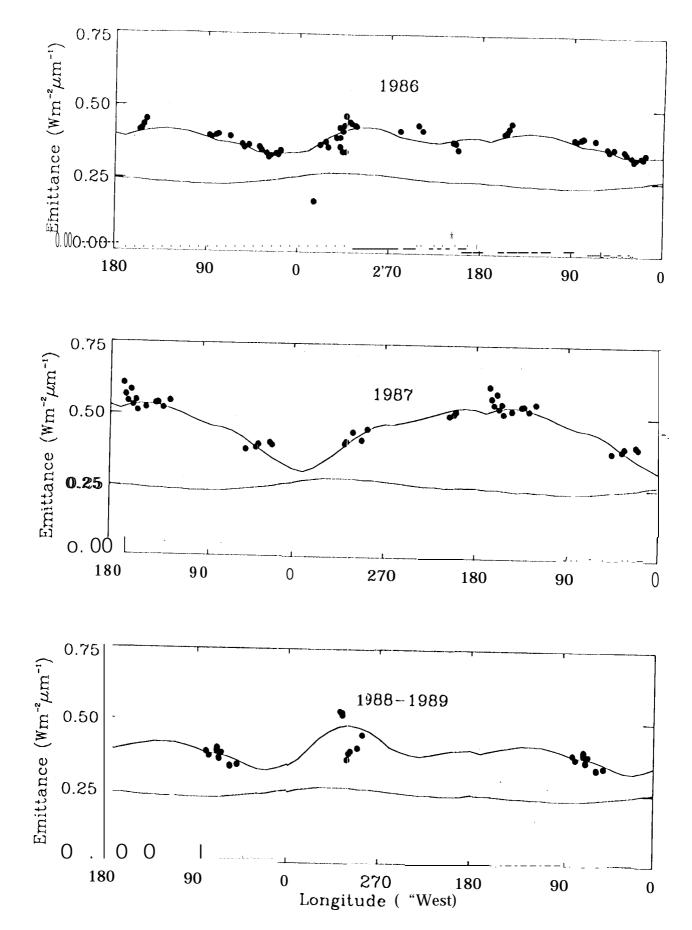


Fig.Sk

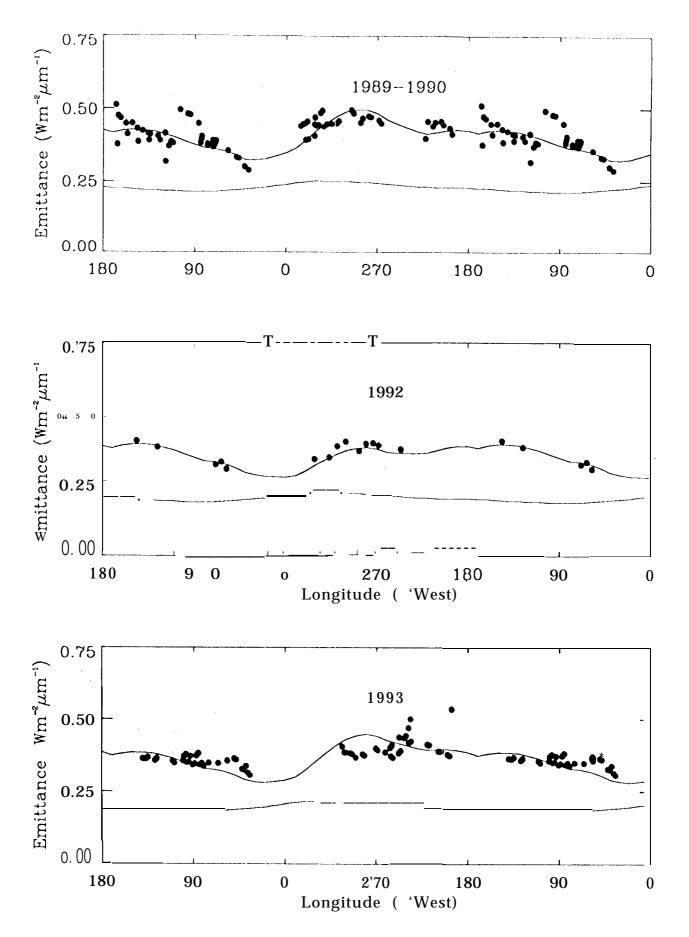


Fig. 5e

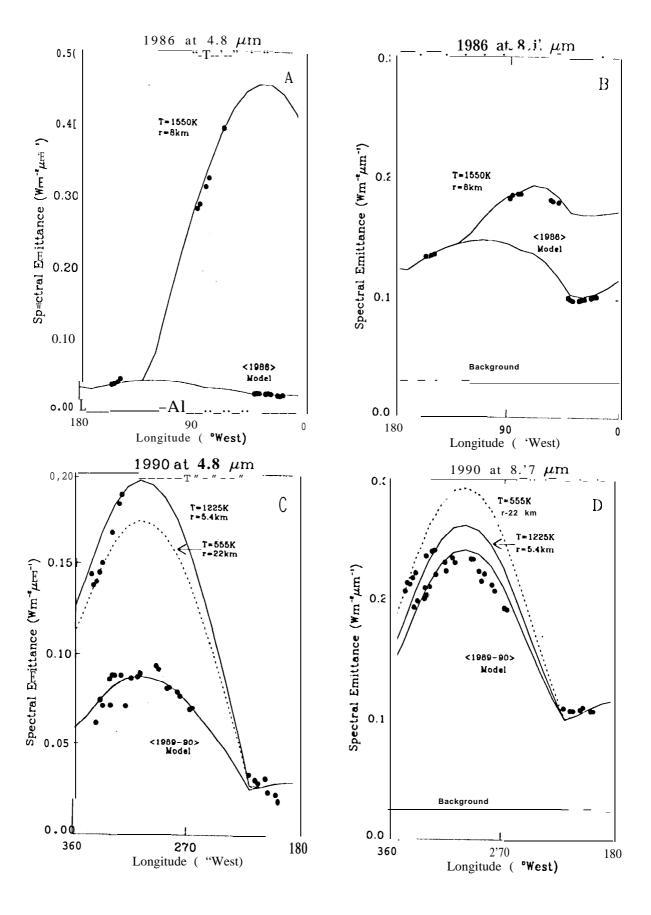


Fig.9

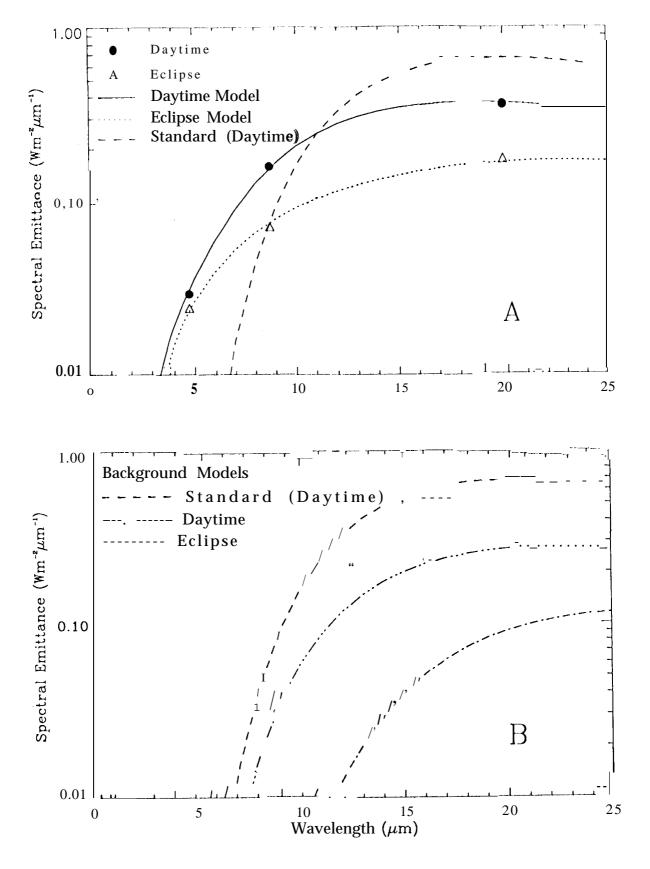


Fig.6

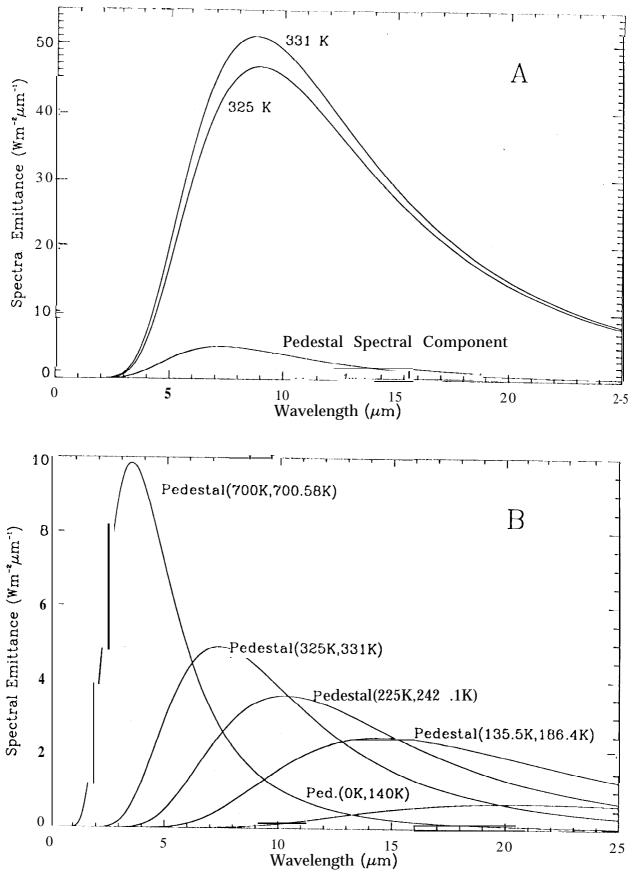


Fig. 7

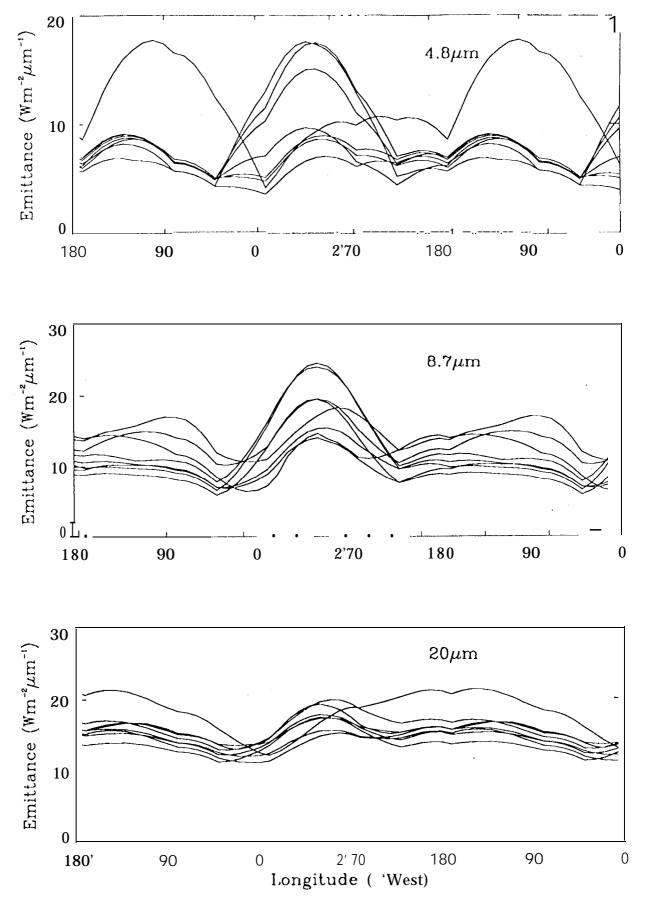
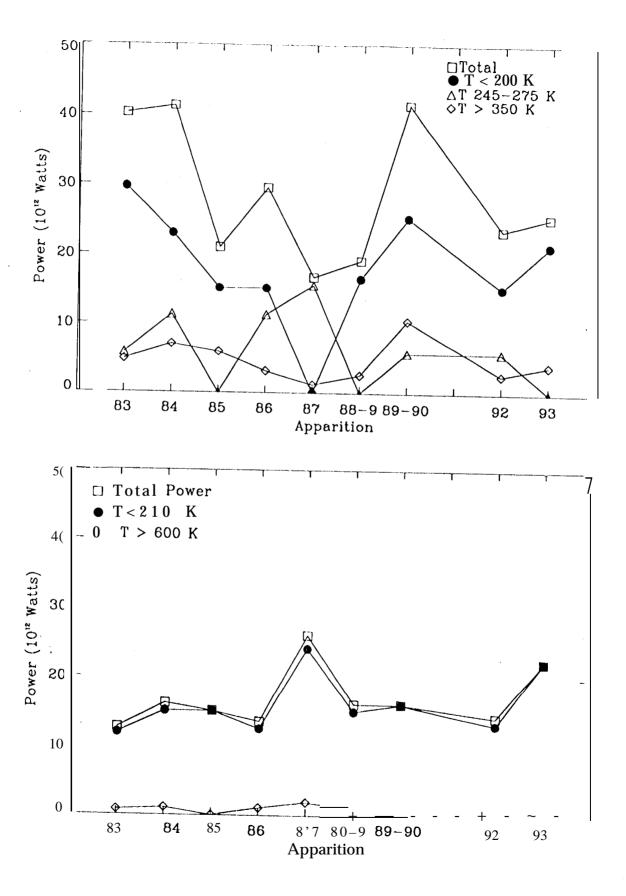


Fig.8



F18.10a

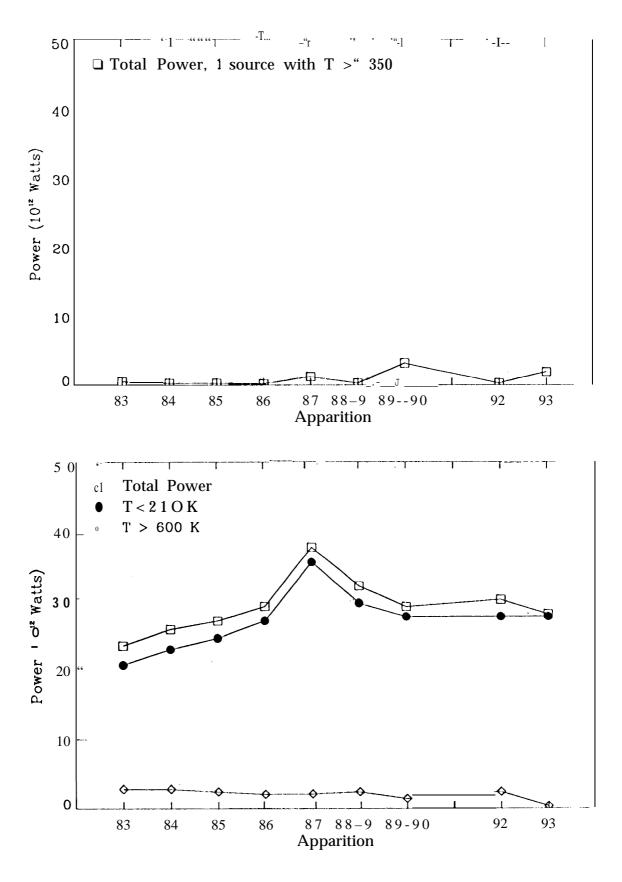


Fig.12

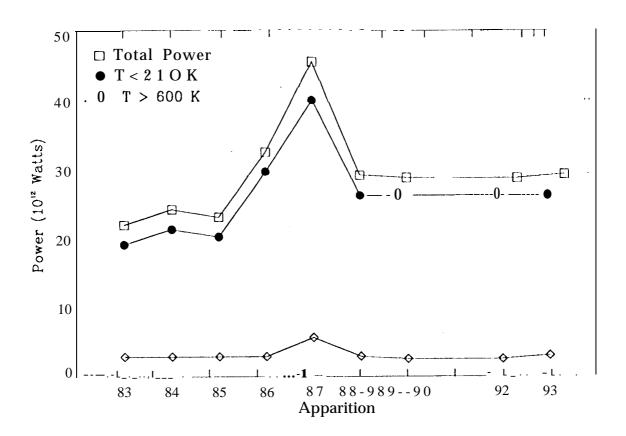
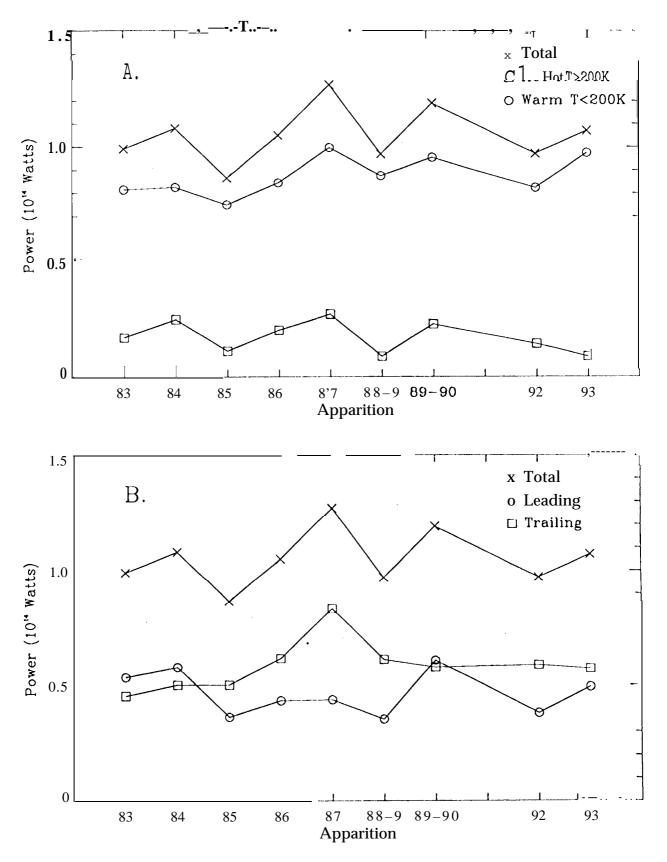
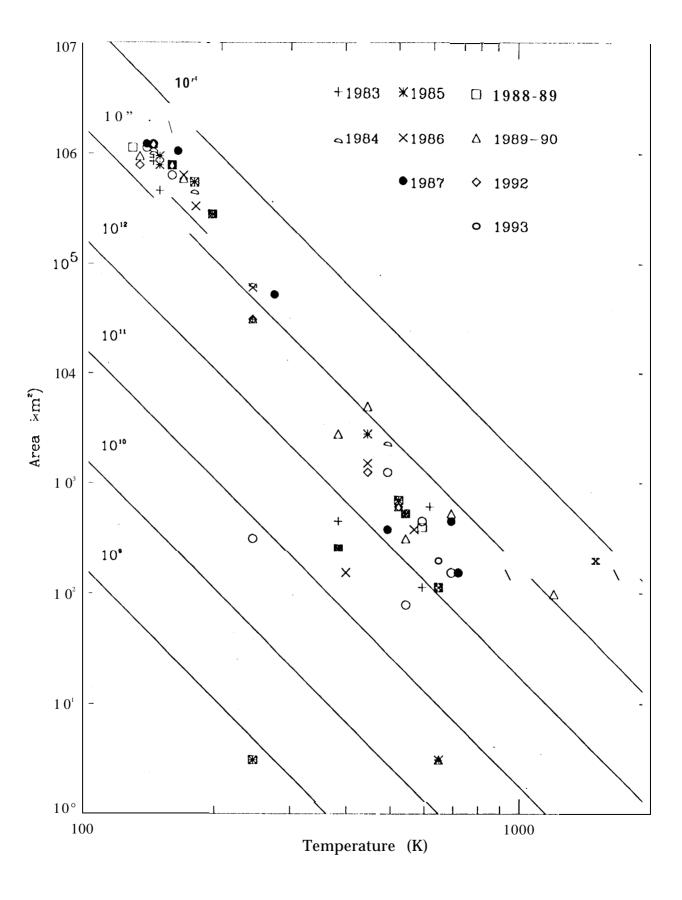


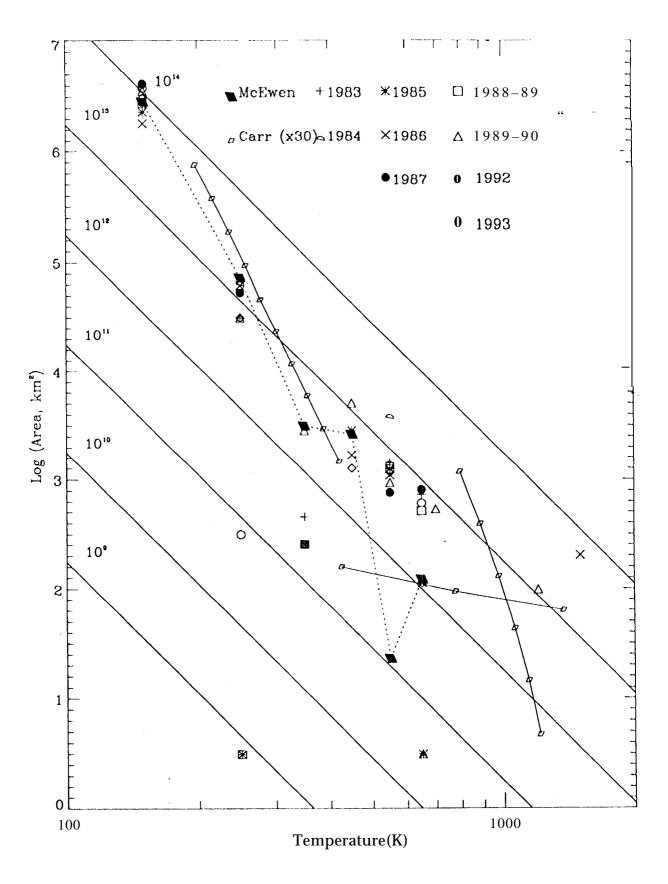
Fig 10.



F18. 11



F18.12a



Figla

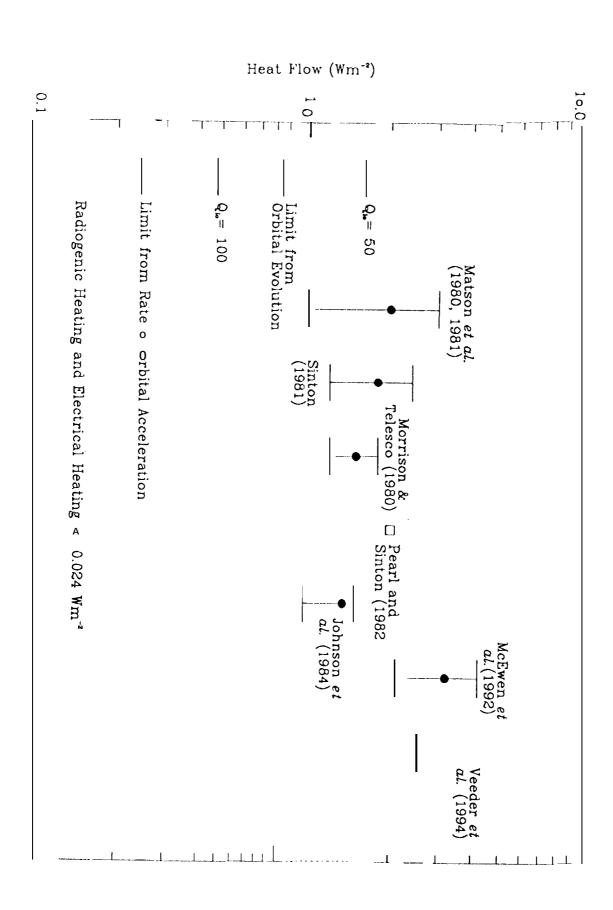


Fig. 13